

# Muon $g - 2$ theory: beyond the SM

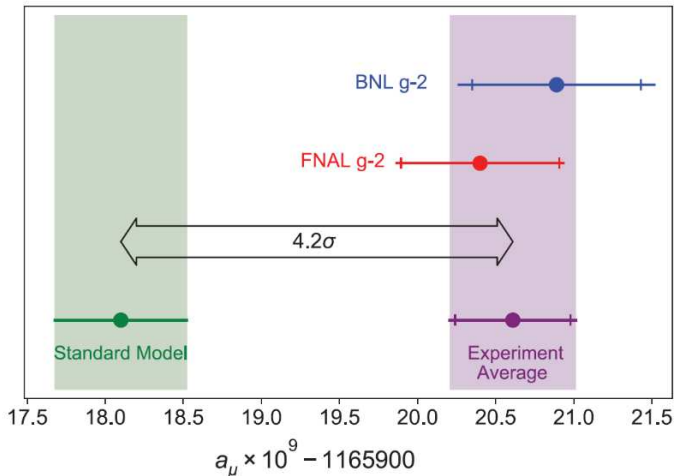
Dominik Stöckinger, TU Dresden

Theory Initiative Workshop KEK (virtual), 28th June 2021

$(g_\mu - 2)/2 = a_\mu$  is among the most precise observables  
sensitive to all known (and unknown?) interactions

# Outline

- 1 Overview and SM theory
- 2  $g = 2$  and BSM — important general remarks
- 3 Examples of concrete models and constraints
- 4 General lessons and conclusions



**Questions: Which models can(not) explain it?**

- Here: general remarks and examples from survey [2104.03691](#)

[Peter Athron, Csaba Balasz, Douglas Jacob, Wojciech Kotlarski, DS, Hyejung Stöckinger-Kim]

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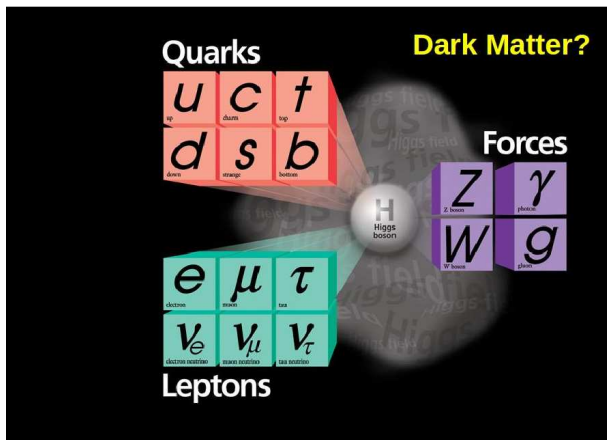
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  - Simple, distinctive properties
  - Overview of contributions
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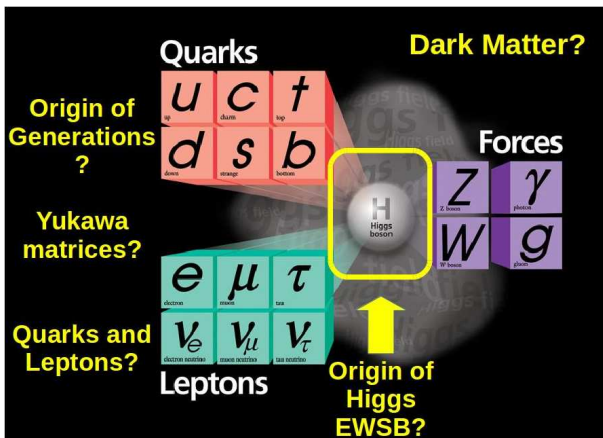
# Open questions require Beyond the Standard Model (BSM) physics



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- experimental clues needed!  $\rightsquigarrow g - 2!$   
not easy to explain!
- relevant and deep questions may be related to  $g - 2$

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# Discrepancy

SM prediction too low by  $\approx (25 \pm 6) \times 10^{-10}$

## Two important general points

discrepancy  $\approx 2 \times a_\mu^{\text{SM,weak}}$

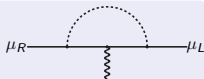
but: expect  $a_\mu^{\text{NP}} \sim a_\mu^{\text{SM,weak}} \times \left(\frac{M_W}{M_{\text{NP}}}\right)^2 \times \text{couplings}$

# Two important general points

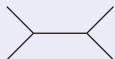
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loop-induced, CP- and Flavor-conserving, chirality-flipping



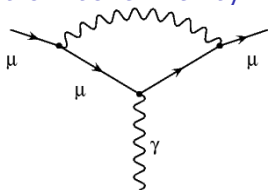
compare:



$b \rightarrow s\gamma$   
EDMs,  $B \rightarrow \tau\nu$   
 $\mu \rightarrow e\gamma$

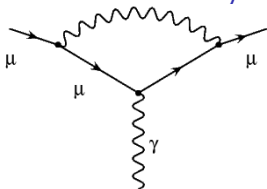
EWPO

# Connection to chirality flip, and structure of BSM



$$\mathcal{L}_{\text{eff}} = - \frac{Qe}{4m_\mu} \mathbf{a}_\mu \times \bar{\psi}_L \sigma_{\mu\nu} \psi_R F^{\mu\nu}$$

# Connection to chirality flip, and structure of BSM



$$\mathcal{L}_{\text{eff}} = - \frac{Qe}{4m_\mu} a_\mu \times \bar{\psi}_L \sigma_{\mu\nu} \psi_R F^{\mu\nu}$$

In the SM QED and QCD: not crucial, trivial chirality flip via gauge invariant muon mass term appearing in propagator numerators

$$a_\mu^{\text{QED, } \tau\text{-loop}} = \frac{\alpha^2}{45\pi^2} \frac{m_\mu^2}{m_\tau^2}$$

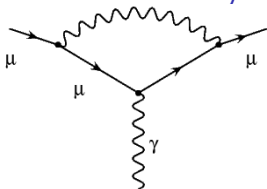
$$a_\mu^{\text{EFT-LbL (Knecht, Nyffeler, Perrotet, de Rafael)}} \sim \left(\frac{\alpha}{\pi}\right)^3 \frac{m_\mu^2}{F_\pi^2}$$

Even in the EW SM contributions, things don't change much, e.g.

$$a_\mu^{\text{W,Z-boson loops}} \sim \frac{\alpha}{4\pi} \frac{m_\mu^2}{M_{W,Z}^2} \approx 15 \times 10^{-10}$$

But the required chirality flip is key to understand BSM contributions!

# Connection to chirality flip, and structure of BSM



$$\mathcal{L}_{\text{eff}} = - \frac{Qe}{4m_\mu} \mathbf{a}_\mu \times \bar{\psi}_L \sigma_{\mu\nu} \psi_R F^{\mu\nu}$$

But:

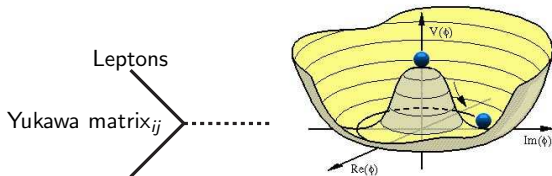
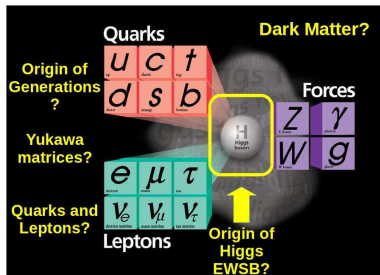
EW gauge invariant  $\mathbf{a}_\mu$ -operator:

$$\bar{L} \sigma_{\mu\nu} \mu_R F^{\mu\nu} \langle H \rangle$$

$$\mathbf{a}_\mu \sim m_\mu \times \underbrace{(\text{some VEV}) \times (\mu_{L \leftrightarrow R}\text{-flipping param.})}_{\text{related to muon mass generation, potential enhancement!}} \times \frac{(\text{other couplings})}{M_{\text{typical}}^2}$$

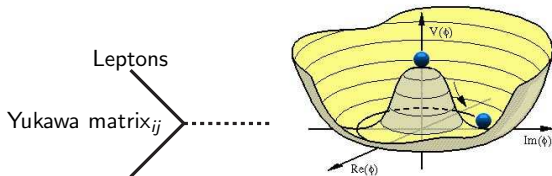
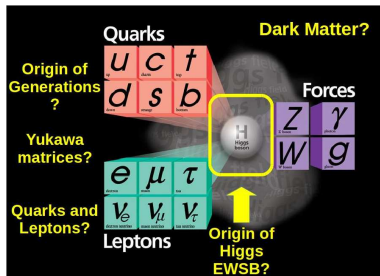
$$m_\mu(\text{SM}) \sim (\text{SM Higgs-VEV}) \times (\text{muon Yukawa coupling})$$

# Window to the muon mass generation mechanism (Higgs/Yukawa sectors)



(changed by new physics?)

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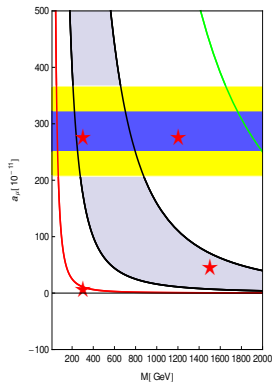
Second connection of  $a_\mu$ : Dark Matter, (light?) dark sectors? Hard to see in detectors but could couple to muon  $\rightsquigarrow$  large effects possible!



Typical behaviour:  $\sim$  chirality flip ( $\rightsquigarrow$  Higgs!) and masses

$$a_\mu \sim \frac{m_\mu \times (\text{some VEV}) \times (\mu_{L \leftrightarrow R}\text{-flipping parameter})}{M_{\text{typical}}^2} \left[ \lesssim \frac{m_\mu^2}{M_{\text{typical}}^2} \text{ (no finetuning)} \right]$$

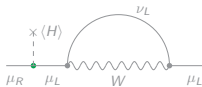
$$\Delta m_\mu \sim (\text{some VEV}) \times (\mu_{L \leftrightarrow R}\text{-flipping parameter})$$



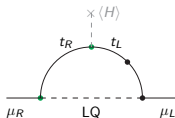
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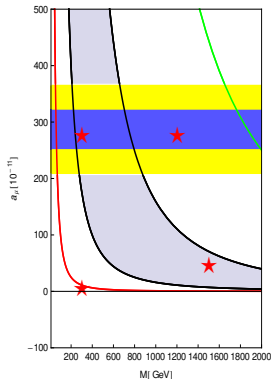
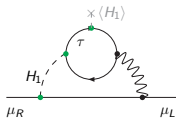
- EWSM:  $\alpha \frac{m_\mu^2}{M_W^2}$   
Similar in  $Z'$ , Dark  $Z_d$  models



- LQ:  $g_L g_R \frac{m_\mu m_t}{M_{LQ}^2}$



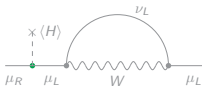
- 2HDM:  $\alpha^2 \tan^2 \beta \frac{m_\mu^2}{M_H^2}$



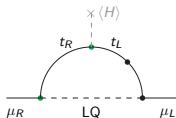
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• EWSM:  $\alpha \frac{m_\mu^2}{M_W^2}$



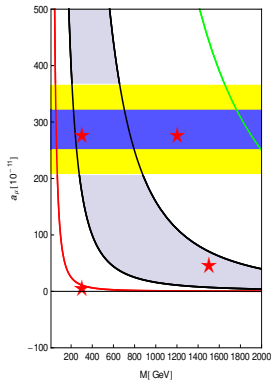
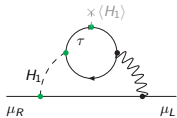
• LQ:  $g_{LGR} \frac{m_\mu m_t}{M_{LQ}^2}$



Can also involve Higgs couplings to  $b$ ,  $c$  or new particles.

Beware:  $\Delta m_\mu / m_\mu \sim g_{LGR} m_t / m_\mu$  restricts couplings

• 2HDM:  $\alpha^2 \tan^2 \beta \frac{m_\mu^2}{M_H^2}$



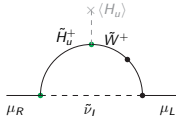
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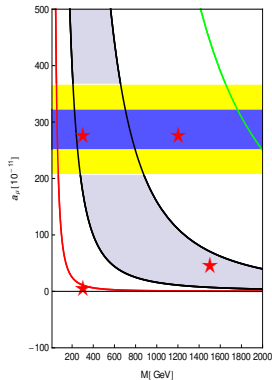
- EWSM:  $\propto \frac{m_\mu^2}{M_W^2}$



- SUSY:  $\propto \frac{m_\mu^2 \tan \beta}{M_{\text{SUSY}}^2} \frac{\mu}{M_{\text{SUSY}}}$



- rad.  $m_\mu \sim \frac{m_\mu^2}{M_{\text{NP}}^2}$



# Typical behaviour: $\sim$ chirality flip ( $\rightsquigarrow$ Higgs!) and masses

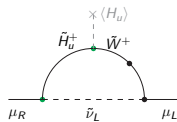
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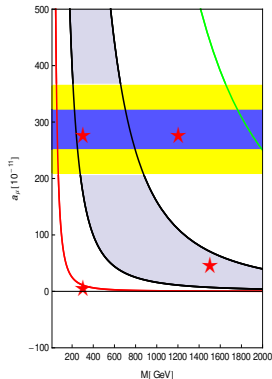
- SUSY:  $\alpha \frac{m_\mu^2 \tan \beta}{M_{\text{SUSY}}^2} \frac{\mu}{M_{\text{SUSY}}}$

Well-motivated theory. Many other advantages



- rad.  $m_\mu \sim \frac{m_\mu^2}{M_{\text{NP}}^2}$

E.g. MSSM for  $\tan \beta \rightarrow \infty$  [Bach,Park,DS,Stöckinger-Kim'15]



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# There are many more examples. . .

## SUSY: **MSSM**, **MRSSM**

- **MSugra**. . . many other generic scenarios
- **Bino-dark matter**+some coannihil.+mass splittings
- **Wino-LSP**+specific mass patterns

## Two-Higgs doublet model

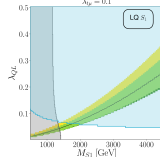
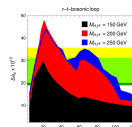
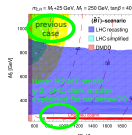
- **Type I, II, Y, Type X**(lepton-specific), flavour-aligned

## Lepto-quarks, vector-like leptons

- scenarios with muon-specific couplings to  $\mu_L$  and  $\mu_R$

## Simple models (one or two new fields)

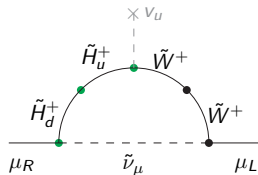
- **Mostly excluded**
- light N.P. (**ALPs**, **Dark Photon**, **Light  $L_\mu - L_\tau$** )



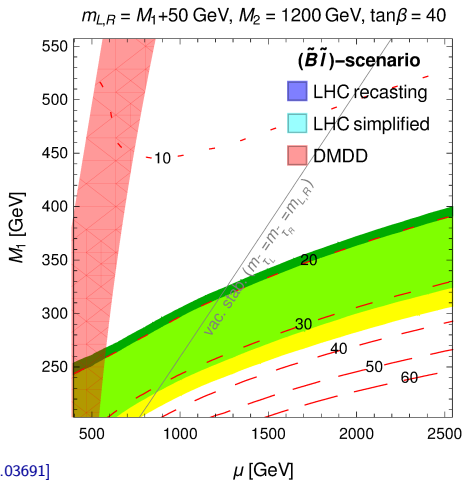
Model	Mass	Spin	CP	Parity	Charge	Color	Notes
1	1.0	0	+	+	0	1	Standard Model (SM)
2	1.0	0	+	-	0	1	Standard Model (SM)
3	1.0	0	-	+	0	1	Standard Model (SM)
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[Athron,Balazs,Jacob,Kotlarski,DS,Stöckinger-Kim, 2104.03691]

# SUSY (MSSM): can explain $g - 2$ and dark matter



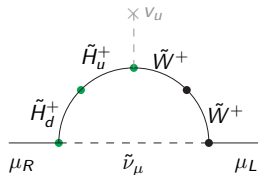
$$a_\mu^{\text{SUSY}} \approx 25 \times 10^{-10} \frac{\tan \beta}{50} \frac{\mu}{M_{\text{SUSY}}} \left( \frac{500 \text{ GeV}}{M_{\text{SUSY}}} \right)^2$$



[2104.03691]

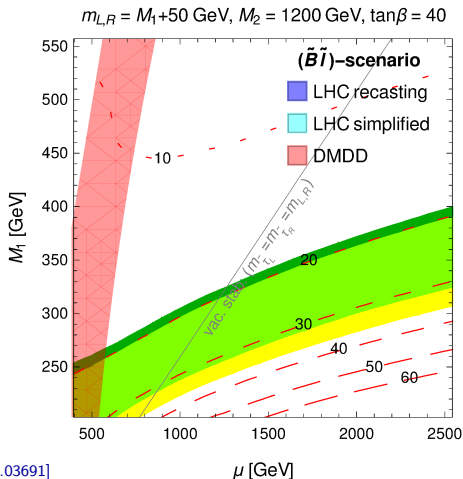


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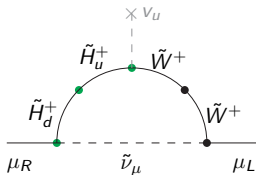
$$a_{\mu}^{\text{SUSY}} \approx 25 \times 10^{-10} \frac{\tan \beta}{50} \frac{\mu}{M_{\text{SUSY}}} \left( \frac{500 \text{ GeV}}{M_{\text{SUSY}}} \right)^2$$

- “Dark matter mass” versus  $\mu$
- explains  $g - 2$  in large region (expands for  $\tan \beta \neq 40$ )
- DM explained by stau/slepton-coannihilation



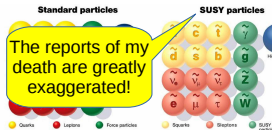
[2104.03691]

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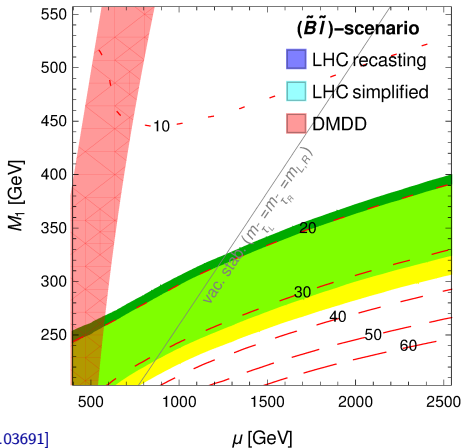


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- this automatically evades (current) LHC limits

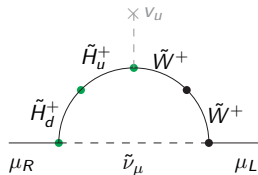


$m_{L,R} = M_1 + 50 \text{ GeV}$ ,  $M_2 = 1200 \text{ GeV}$ ,  $\tan \beta = 40$



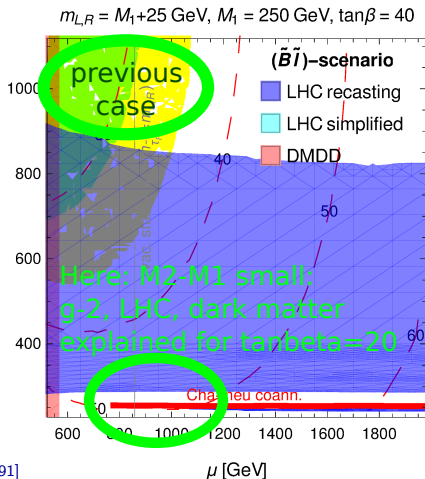
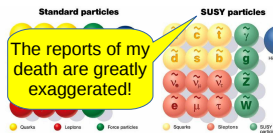
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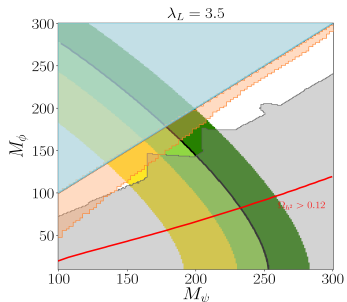
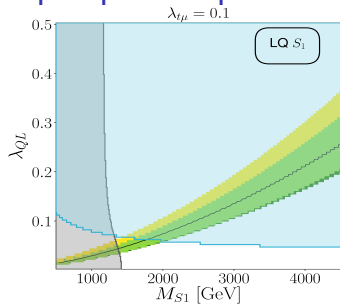
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- Strong LHC limits on  $M_2$
- DM also explained by Wino-coannihilation
- again evades (current) LHC limits



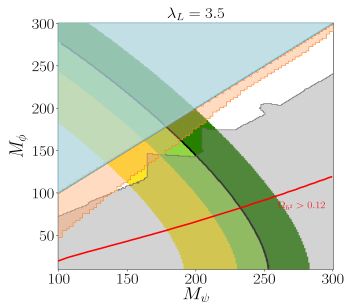
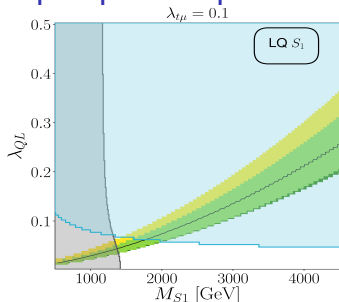
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# Leptoquarks: promising; 2-field models: less



[Athron,Balazs,Jacob,Kotlarski,DS,Stöckinger-Kim, 2104.03691 ]

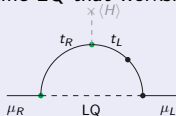
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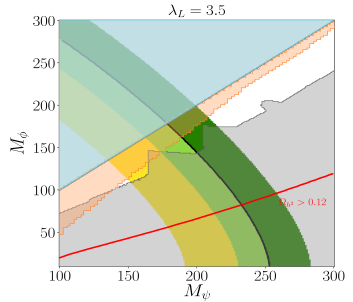
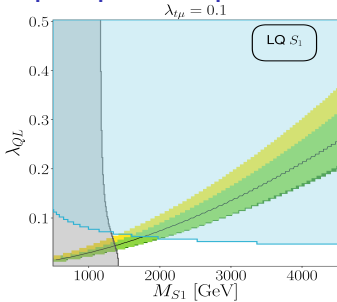
$$a_\mu \text{ from LQ (or VLL)} \quad \mathcal{L}_{S_1} = - (\lambda_{QL} Q_3 \cdot L_2 S_1 + \lambda_{t\mu} t_\mu S_1^*)$$

Specific LQ that works:



- Chiral enhancement  $\sim y_{\text{top}}, y_{\text{VLL}}$  versus  $y_\mu$
- LHC: lower mass limits
- Flavour constraints  $\rightsquigarrow$  assume **only couplings to muons**
- Viable window above LHC (without  $m_\mu$ -finetuning)

# Leptoquarks: promising; 2-field models: less

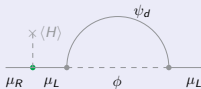


[Athron,Balazs,Jacob,Kotlarski,DS,Stöckinger-Kim, 2104.03691 ]

## $a_\mu$ from 2-field model L

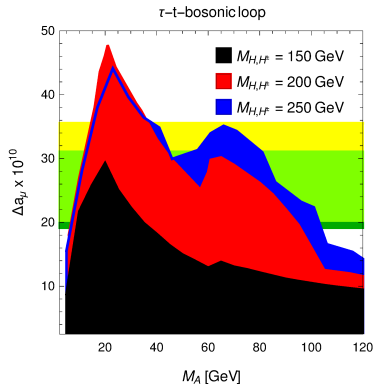
- No chiral enhancement, need very large couplings
- LHC: lower mass limits
- Dark matter candidate, but incompatible with large  $a_\mu$

General result:  $a_\mu$  and DM require at least three new fields! see also: [Arcadi,Calibbi,Fedele,Mescia] on B-physics



# Two-Higgs doublet model: promising in specific versions

- Aligned 2-Higgs doublet model, rich new Higgs/Yukawa sectors



[2104.03691]

Details on Yukawa couplings:

Type X/lepton-specific:  $Y_\ell \propto \tan \beta$

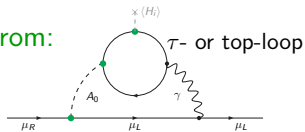
Type II:  $Y_{\ell,d} \propto \tan \beta$

Aligned:  $Y_\ell \propto \zeta_\ell$

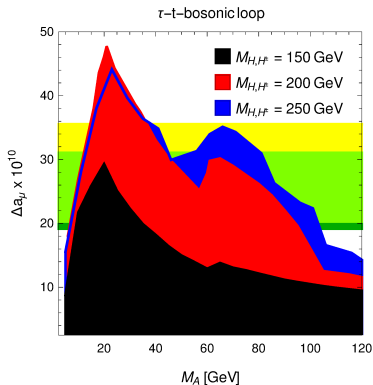
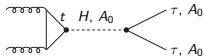
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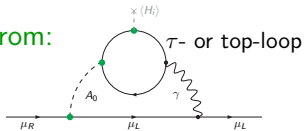
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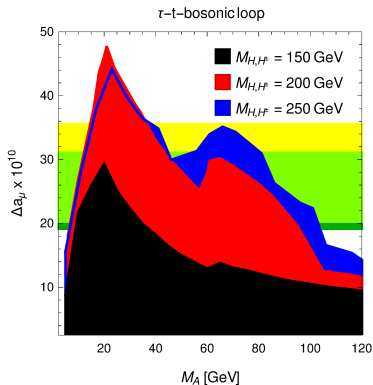
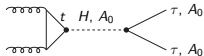
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$a_\mu$  from:



constraints (LHC, B-,  $\tau$ -physics):



[2104.03691]

- can explain  $g - 2$
- need large new Yukawa couplings
- under pressure, testable at LHC, lepton colliders, B-physics

# Light/dark sectors — compatible with large $a_\mu$ ?

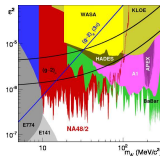
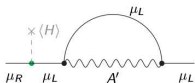
Very light, weakly interacting new particles

# Light/dark sectors — compatible with large $a_\mu$ ?

Very light, weakly interacting new particles

- “dark photon” **NO**

$$\mathcal{L} = -\frac{\epsilon}{2 \cos \theta_W} F^{\mu\nu} B_{\mu\nu} \quad a_\mu \sim \frac{\alpha}{2\pi} \epsilon^2$$



[NA48: 1504.00607]  
excludes minimal dark photon for  $a_\mu$

- “dark  $Z_d$ ” **Better**

$$a_\mu \sim \frac{\alpha}{2\pi} (\epsilon + \delta' m_{Z_d} / m_Z)^2$$

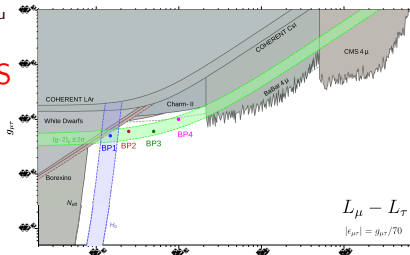
Additional mass mixing  $\delta$ , may assume invisible decays into dark sector, can evade limits (still nontrivial)

Davoudiasl, Lee, Marciano . . . Cadeddu, Cargioli, Dordei, Giunti, Picciao

- $Z'$  with quantum number  $L_\mu - L_\tau$  **YES**

[Ma, Roy, Roy'01, Heeck, Rodejohann'11. . .]  
(plot from [Amaral, Cerdeno, Cheek, Foldenauer'21])

Evades collider constraints,  
subject to low-E constraints,  
viable window 10 . . . 100 MeV



$L_\mu - L_\tau$   
 $|c_{\mu\tau}| = g_{\mu\tau} / 70$

# Outline

- 1 Overview and SM theory
- 2  $g = 2$  and BSM — important general remarks
- 3 Examples of concrete models and constraints
- 4 General lessons and conclusions

# There are many more examples. . .

## SUSY: **MSSM**, **MRSSM**

- **MSugra**. . . many other generic scenarios
- **Bino-dark matter**+some coannihil.+mass splittings
- **Wino-LSP**+specific mass patterns

## Two-Higgs doublet model

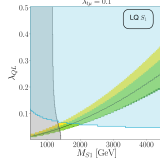
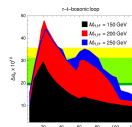
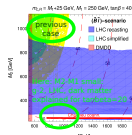
- **Type I, II, Y, Type X**(lepton-specific), flavour-aligned

## Lepto-quarks, vector-like leptons

- scenarios with muon-specific couplings to  $\mu_L$  and  $\mu_R$

## Simple models (one or two new fields)

- **Mostly excluded**
- light N.P. (**ALPs**, **Dark Photon**, **Light  $L_\mu - L_\tau$** )



MSSM	Non-CP	CP	W	Z	U(1)	CP	W	Z	U(1)	CP	W	Z	U(1)
1	1	1	1	1	1	1	1	1	1	1	1	1	1
2	1	1	1	1	1	1	1	1	1	1	1	1	1
3	1	1	1	1	1	1	1	1	1	1	1	1	1
4	1	1	1	1	1	1	1	1	1	1	1	1	1
5	1	1	1	1	1	1	1	1	1	1	1	1	1
6	1	1	1	1	1	1	1	1	1	1	1	1	1
7	1	1	1	1	1	1	1	1	1	1	1	1	1
8	1	1	1	1	1	1	1	1	1	1	1	1	1
9	1	1	1	1	1	1	1	1	1	1	1	1	1
10	1	1	1	1	1	1	1	1	1	1	1	1	1
11	1	1	1	1	1	1	1	1	1	1	1	1	1
12	1	1	1	1	1	1	1	1	1	1	1	1	1
13	1	1	1	1	1	1	1	1	1	1	1	1	1
14	1	1	1	1	1	1	1	1	1	1	1	1	1
15	1	1	1	1	1	1	1	1	1	1	1	1	1
16	1	1	1	1	1	1	1	1	1	1	1	1	1
17	1	1	1	1	1	1	1	1	1	1	1	1	1
18	1	1	1	1	1	1	1	1	1	1	1	1	1
19	1	1	1	1	1	1	1	1	1	1	1	1	1
20	1	1	1	1	1	1	1	1	1	1	1	1	1
21	1	1	1	1	1	1	1	1	1	1	1	1	1
22	1	1	1	1	1	1	1	1	1	1	1	1	1
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48	1	1	1	1	1	1	1	1	1	1	1	1	1
49	1	1	1	1	1	1	1	1	1	1	1	1	1
50	1	1	1	1	1	1	1	1	1	1	1	1	1

[Athron,Balazs,Jacob,Kotlarski\_DS,Stöckinger-Kim, 2104.03691]

# Summary of main points

discrepancy  $\approx 2 \times a_\mu^{\text{SM,weak}}$

but: expect  $a_\mu^{\text{NP}} \sim a_\mu^{\text{SM,weak}} \times \left(\frac{M_W}{M_{\text{NP}}}\right)^2 \times \text{couplings}$

$a_\mu$  is loop-induced, CP- and flavor-conserving and chirality-flipping

rather light, neutral (?) particles  $\rightsquigarrow$  Connection to dark matter?

Chirality flip enhancement  $\rightsquigarrow$  Window to muon mass generation? EWSB/generations?

## Which models can still accommodate large deviation?

Many (but not all) models!

but always: **experimental constraints!**

## Outlook:

- $g - 2 + \text{LHC, DM} \rightsquigarrow$  **constraints on BSM physics, great potential for future**
- often chirality flips/new flavor structures/light particles  $\rightsquigarrow$  tests: Higgs couplings,  $B$ -physics, CLFV, EDM, light-particle searches,  $e^+e^-$ /muon collider

**20 years after BNL... deviation confirmed ... very promising future!**

# Dark matter and LHC constraints (rules of thumb)

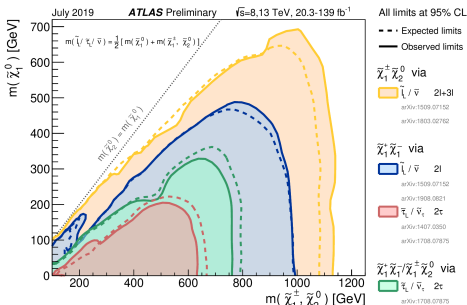
Dark matter (Wimps):

- Need to obtain  $\Omega_{\text{DM}} = \Omega_{\text{obs}}$  or at least  $\Omega_{\text{DM}} < \Omega_{\text{obs}}$ , requires specific annihilation cross sections
- for Binos often too high, for scalars often too low
- Higgsinos or general fermion doublets: need masses around 1 TeV
- DM direct detection: limits on interactions with quarks: sometimes correlated with annihilation interactions

LHC:

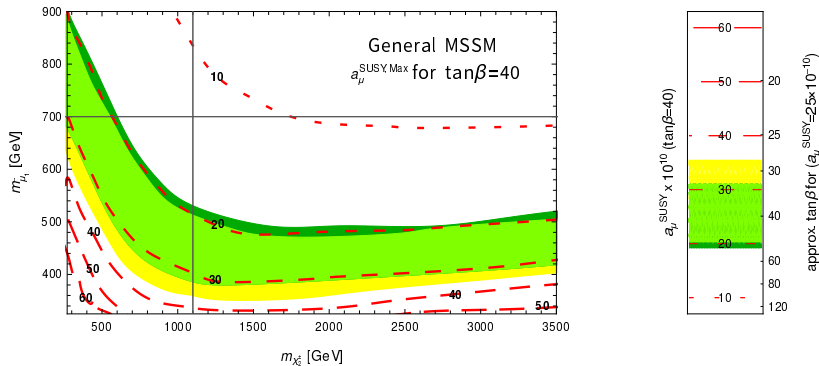
- “triangular limits”, depending on assumptions/decay modes
- either  $\text{Cha} > 1100$ ,  $\text{smuons} > 700$  GeV
- or somewhat small mass splittings  $\lesssim 100$  to LSP
- or more complicated decay modes

E.g.: limits on Higgsinos are weak, limits on Winos weaker if sleptons  $>$  Winos



# Full MSSM overview in 7 plots

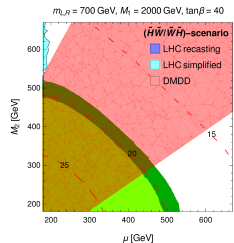
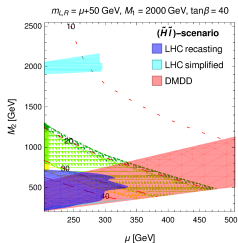
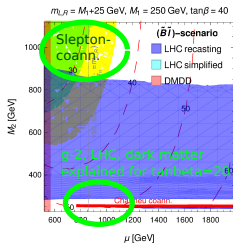
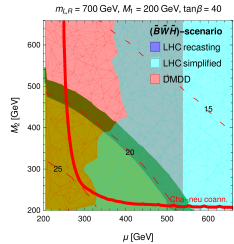
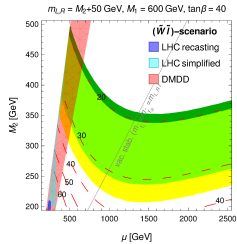
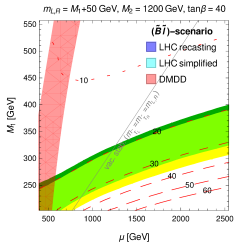
[Peter Athron, Csaba Balasz, Douglas Jacob, Wojciech Kotlarski, DS, Hyejung Stöckinger-Kim, 2104.03691]





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[Peter Athron, Csaba Balasz, Douglas Jacob, Wojciech Kotlarski, DS, Hyejung Stöckinger-Kim, 2104.03691]



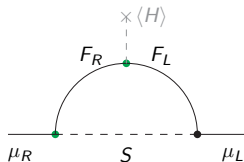
Summary: Bino-LSP:  $a_\mu$  and DM. Wino-/Higgsino-LSP:  $a_\mu$ . Both  $\chi \langle \text{slepton} \rangle \approx$  disfavoured.

DM+LHC  $\Rightarrow$  mass patterns! Coannihilation regions help! Specific cases excluded, e.g. Constrained MSSM





# Three-field models



- many models: viable, large chirality enhancements
- can explain  $a_\mu^{\text{BNL}}$  and LHC and dark matter