### Muon g-2: a new data-based analysis

## **D. Nomura (KEK)**

### talk at 2018 WPI-next Mini-workshop "Hints for New Physics in Heavy Flavors" @ Nagoya U.

November 16, 2018

Ref: A. Keshavarzi, DN and T. Teubner (KNT) Phys. Rev. D97 (2018) 114025 [arXiv:1802.02995]

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SM prediction for muon g-2

# Muon g-2: introduction

Lepton magnetic moment  $\vec{\mu}$ :

$$\vec{\mu} = -g \frac{e}{2m} \vec{s}$$
,  $(\vec{s} = \frac{1}{2} \vec{\sigma}$  (spin),  $g = 2 + 2F_2(0)$ )

where

$$\overline{u}(p+q)\Gamma^{\mu}u(p) = \overline{u}(p+q)\left(\gamma^{\mu}F_{1}(q^{2}) + \frac{i\sigma^{\mu\nu}q_{\nu}}{2m}F_{2}(q^{2})\right)u(p)$$

Anomalous magnetic moment:  $a \equiv (g-2)/2 \ (=F_2(0))$ 

Historically,

★ g = 2 (tree level, Dirac) ★  $a = \alpha/(2\pi)$  (1-loop QED, Schwinger)

Today, still important, since...

★ One of the most precisely measured quantities:

 $a_{\mu}^{\exp} = 11\ 659\ 208.9(6.3) \times 10^{-10}$  [0.5ppm]

#### ★ Extremely useful in probing/constraining physics beyond the SM

SM prediction for muon g-2



(Bennett et al)



# Muon g-2: previous exp. (after 1960)

Experiment	Years	Polarity	$a_{\mu} \times 10^{10}$	Precision [	ppm]	Sensitivity
CERN I	1961	$\mu^+$	11450000(220000)	4300	2-lo	op QED contrib. (3600 ppm)
CERN II	1962-1968	$\mu^+$	11661600(3100)	270	3-lo	op QED contrib. (260 ppm)
CERN III	1974-1976	$\mu^+$	11659100(110)	10	had	ronic vacuum polarization
CERN III	1975-1976	$\mu^{-}$	11659360(120)	10	Com	(60 ppm)
BNL	1997	$\mu^+$	11659251(150)	13		
BNL	1998	$\mu^+$	11659191(59)	5	4-lo	op QED contrib. (3.3 ppm)
BNL	1999	$\mu^+$	11659202(15)	1.3	elec	troweak contrib. (1.3 ppm)
BNL	2000	$\mu^+$	11659204(9)	0.73	had	ronic light-by-light contrib.
BNL	2001	$\mu^{-}$	11659214(9)	0.72	had	ronic NLO vacuum pol.
Average			11659208.0(6.3)	0.54	contrib. (-0.85 ppm)	

Table from BNL-E821 final report, Phys. Rev. D 73 (2006) 072003

#### Testing the SM for more than 50 years! 2 more exp. to come: those from Fermilab and J-PARC

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## Breakdown of SM prediction for muon g-2

	<u>2011</u>		<u>2018</u>		
QED	11658471.81 (0.02)	$\longrightarrow$	11658471.90(0.01) [arXiv:1712.06060]		
EW	15.40 (0.20)	$\longrightarrow$	15.36 (0.10) [Phys. Rev. D 88 (2013) 053005]		
LO HLbL	10.50 (2.60)	$\longrightarrow$	9.80 (2.60) [EPJ Web Conf. 118 (2016) 01016]		
NLO HLbL			0.30 (0.20) [Phys. Lett. B 735 (2014) 90]		
	HLMNT11		<u>KNT18</u>		
LO HVP	694.91 (4.27)	$\longrightarrow$	693.27 (2.46) this work		
NLO HVP	-9.84 (0.07)	$\longrightarrow$	-9.82 (0.04) this work		
NNLO HVP			1.24 (0.01) [Phys. Lett. B 734 (2014) 144]		
Theory total	11659182.80 <mark>(4.94)</mark>	$\longrightarrow$	11659182.05 (3.56) this work		
Experiment			11659209.10 (6.33) world avg		
Exp - Theory	26.1 (8.0)	$\longrightarrow$	27.1 (7.3) this work		
$\Delta a_{\mu}$	3.3σ	$\rightarrow$	<b>3.7</b> $\sigma$ this work		
(HVP: Hadronic Vacuum Polarization) (HLbL: Hadronic Light-by-Light) Slide by A. Keshavarzi (Liverpool) at 'Muon $g - 2$ Workshop' at Mainz, June 18-22, 2018					
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# **Hadronic Contributions**

There are several hadronic contributions:



LO: Leading Order (or Vacuum Polarization) Hadronic Contribution NLO: Next-to-Leading Order Hadronic Contribution I-by-I: Hadronic light-by-light Contribution



The diagram to be evaluated:



pQCD not useful. Use the dispersion relation and the optical theorem.



$$a_{\mu}^{\rm had,LO} = \frac{m_{\mu}^2}{12\pi^3} \int_{s_{\rm th}}^{\infty} ds \ \frac{1}{s} \hat{K}(s) \sigma_{\rm had}(s)$$



• Weight function  $\hat{K}(s)/s = \mathcal{O}(1)/s$   $\implies$  Lower energies more important  $\implies \pi^{+}\pi^{-}$  channel: 73% of total  $a_{\mu}^{\text{had,LO}}$ 

- Lots of new input  $\sigma(e^+e^- 
  ightarrow$  hadrons) data
- Improvements in the estimates of uncertainties due to radiative corrections (Vacuum Polarization Radiative Corrections & Final State Radiations)
- Improvements in data-combination method

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Channel	Energy range [GeV]	$a_{\mu}^{\rm had,LOVP} \times 10^{10}$	$\Delta \alpha_{\rm had}^{(5)}(M_Z^2) \times 10^4$	New data	
	Chiral perturbation the	eory (ChPT) threshold contr	ibutions		Breakdown of contributions
$\pi^0 \gamma$	$m_x \le \sqrt{s} \le 0.600$	$0.12 \pm 0.01$	$0.00 \pm 0.00$		to a (had $IO VD$ ) from
$\pi^{+}\pi^{-}$	$2m_{\pi} \le \sqrt{s} \le 0.305$	$0.87 \pm 0.02$	$0.01 \pm 0.00$		$u_{\mu}(hau, LOVP)$ from
$\pi^{+}\pi^{-}\pi^{0}$	$3m_{\pi} \le \sqrt{s} \le 0.660$	$0.01 \pm 0.00$	$0.00 \pm 0.00$		various hadronic final states
117	$m_{\eta} \le \sqrt{s} \le 0.660$	$0.00 \pm 0.00$	$0.00 \pm 0.00$		various nauronic final states
	Data based c	hannels ( $\sqrt{s} \le 1.937$ GeV)			
$\pi^{0}\gamma$	$0.600 \le \sqrt{s} \le 1.350$	$4.46 \pm 0.10$	$0.36 \pm 0.01$	[65]	
<i>π</i> <sup>-</sup> <i>π</i> <sup>-</sup>	$0.305 \le \sqrt{s} \le 1.937$	$502.97 \pm 1.97$	$34.26 \pm 0.12$	[34,35]	
<i>π</i> <sup>-</sup> <i>π</i> <sup>-</sup> <i>π</i> <sup>0</sup>	$0.660 \le \sqrt{s} \le 1.937$	$47.79 \pm 0.89$	$4.77 \pm 0.08$	[36]	
<i>π π π π</i>	$0.613 \le \sqrt{s} \le 1.937$	14.87 ± 0.20	4.02 ± 0.05	[40,42]	
$\pi^{+}\pi^{-}\pi^{-}\pi^{-}\pi^{-}$	$0.850 \le \sqrt{s} \le 1.937$	19.39 ± 0.78 0.00 ± 0.00	$5.00 \pm 0.20$ $0.33 \pm 0.03$	[44]	
$(2\pi^{+}2\pi^{-}\pi^{0})_{noy}$	$1.013 \le \sqrt{s} \le 1.937$	0.99 ± 0.09	$0.33 \pm 0.03$		We have included new data sets
$3\pi^{-}3\pi^{-}$	$1.313 \le \sqrt{s} \le 1.937$	$0.23 \pm 0.01$	$0.09 \pm 0.01$	[00]	
$(2\pi^{+}2\pi^{-}2\pi^{''})_{nayo}$	$1.322 \le \sqrt{s} \le 1.937$	1.35 ± 0.17	$0.51 \pm 0.06$		from $\sim 30$ papers,
K * K *	$0.988 \le \sqrt{s} \le 1.937$	$23.03 \pm 0.22$	$3.37 \pm 0.03$	[45,46,49]	
K <sup>o</sup> <sub>S</sub> K <sup>o</sup> <sub>L</sub>	$1.004 \le \sqrt{s} \le 1.937$	$13.04 \pm 0.19$	$1.77 \pm 0.03$	[50,51]	in addition to those included
КК <b>л</b>	$1.260 \le \sqrt{s} \le 1.937$	$2.71 \pm 0.12$ 1.02 ± 0.00	$0.89 \pm 0.04$	[53,54]	
KK2π	$1.350 \le \sqrt{s} \le 1.957$	$1.93 \pm 0.08$ 0.70 ± 0.02	$0.75 \pm 0.03$ $0.00 \pm 0.00$	[50,53,55]	in the HLMNT11 analysis
ηγ w= <sup>+</sup> = <sup>-</sup>	$1.001 \le \sqrt{3} \le 1.700$	$1.20 \pm 0.02$	$0.09 \pm 0.00$ $0.30 \pm 0.02$	[67]	
$(n\pi^{+}\pi^{-}\pi^{0})$	$1.091 \le \sqrt{3} \le 1.937$ $1.333 \le \sqrt{5} \le 1.937$	$1.29 \pm 0.00$ $0.60 \pm 0.15$	$0.39 \pm 0.02$ $0.21 \pm 0.05$	[08,09]	
"2"+2"-	$1.338 \le \sqrt{s} \le 1.937$	$0.08 \pm 0.01$	$0.03 \pm 0.00$	[70]	
100	$1.333 \le \sqrt{s} \le 1.937$	$0.31 \pm 0.03$	$0.00 \pm 0.00$ $0.10 \pm 0.01$	[70.71]	We have included $\sim$ 30 hadronic
$\alpha(\rightarrow \pi^0 \gamma) \pi^0$	$0.920 \le \sqrt{s} \le 1.937$	$0.88 \pm 0.02$	$0.19 \pm 0.00$	[72,73]	C
nd	$1.569 \le \sqrt{s} \le 1.937$	$0.42 \pm 0.03$	$0.15 \pm 0.01$		final states
$\phi \rightarrow$ unaccounted	$0.988 \le \sqrt{s} \le 1.029$	$0.04 \pm 0.04$	$0.01 \pm 0.01$		
$\eta \omega \pi^0$	$1.550 \le \sqrt{s} \le 1.937$	$0.35 \pm 0.09$	$0.14 \pm 0.04$	[74]	
$\eta \rightarrow npp K \bar{K}_{roden K\bar{K}}$	$1.569 \le \sqrt{s} \le 1.937$	$0.01 \pm 0.02$	$0.00 \pm 0.01$	[53,75]	
pp	$1.890 \le \sqrt{s} \le 1.937$	$0.03 \pm 0.00$	$0.01 \pm 0.00$	[76]	At $2 \lesssim \sqrt{s} \lesssim 11$ GeV,
nñ	$1.912 \le \sqrt{s} \le 1.937$	$0.03 \pm 0.01$	$0.01 \pm 0.00$	[77]	we use inclusively measured data
	Estimated cont	ributions ( $\sqrt{s} \le 1.937$ GeV)			we use inclusively measured data
$(\pi^{+}\pi^{-}3\pi^{0})_{ma}$	$1.013 \le \sqrt{s} \le 1.937$	$0.50 \pm 0.04$	$0.16 \pm 0.01$		
$(\pi^+\pi^-4\pi^0)_{max}$	$1.313 \le \sqrt{s} \le 1.937$	$0.21 \pm 0.21$	$0.08 \pm 0.08$		
ККЗл	$1.569 \le \sqrt{s} \le 1.937$	$0.03 \pm 0.02$	$0.02 \pm 0.01$		At higher energies > 11 GeV
$\omega(\rightarrow npp)2\pi$	$1.285 \le \sqrt{s} \le 1.937$	$0.10 \pm 0.02$	$0.03 \pm 0.01$		At higher energies $\gtrsim$ 11 GeV,
$\omega(\rightarrow npp)3\pi$	$1.322 \le \sqrt{s} \le 1.937$	$0.17 \pm 0.03$	$0.06 \pm 0.01$		
$\omega(\rightarrow npp)KK$	$1.569 \le \sqrt{s} \le 1.937$	$0.00 \pm 0.00$	$0.00 \pm 0.00$		we use pool
$\eta \pi^{+} \pi^{-} 2 \pi^{0}$	$1.338 \le \sqrt{s} \le 1.937$	$0.08 \pm 0.04$	$0.03 \pm 0.02$		
	Other contril	putions ( $\sqrt{s} > 1.937$ GeV)			
Inclusive channel	$1.937 \le \sqrt{s} \le 11.199$	$43.67 \pm 0.67$	$82.82 \pm 1.05$	[56,62,63]	
$J/\psi$		$6.26 \pm 0.19$	$7.07 \pm 0.22$		
$\psi'$		$1.58 \pm 0.04$	$2.51 \pm 0.06$		
$\Upsilon(1S - 4S)$		$0.09 \pm 0.00$	$1.06 \pm 0.02$		
pQCD	$11.199 \le \sqrt{s} \le \infty$	$2.07 \pm 0.00$	$124.79 \pm 0.10$		
		(02.26 ) 2.46	076111111		

#### Table from KNT18, Phys. Rev. D97 (2018) 114025

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## $\sigma^0_{\mathrm{had},\gamma}:$ vacuum polarisation corrections



 $\Rightarrow$  Photon VP corresponds to higher order contributions to  $a_{\mu}^{\rm had,\,VP}$ 



- $\Rightarrow$  Fully updated, self-consistent VP routine: [vp\_knt\_v3\_0], available for distribution
  - $\rightarrow$  Cross sections undressed with full photon propagator (must include imaginary part),  $\sigma_{\rm had}^0(s) = \sigma_{\rm had}(s) |1 \Pi(s)|^2$
- $\Rightarrow \text{ If correcting data, apply corresponding radiative correction uncertainty} \\ \rightarrow \text{Take } \frac{1}{3} \text{ of total correction per channel as conservative extra uncertainty}$

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### $\sigma_{\rm had \ \gamma}^0$ : final state radiation corrections



 $\Rightarrow$  Photon FSR formally higher order corrections to  $a_u^{\rm had,\,VP}$ 



- $\Rightarrow$  Cannot be unambiguously separated, not accounted for in HO contributions
  - $\rightarrow$  Must be included as part of 1PI hadronic blobs
- $\Rightarrow$  Experiment may cut/miss photon FSR  $\rightarrow$  Must be added back
- $\Rightarrow$  For  $\pi^+\pi^-$ , sQED approximation [Eur. Phys. J. C 24 (2002) 51, Eur. Phys. J. C 28 (2003) 261]
- $\Rightarrow$  For higher multiplicity states, Need new, more developed tools to increase difficult to estimate correction precision here
  - (e.g. CARLOMAT 3.1 [Eur.Phys.J. C77 (2017) no.4, 254 ]?) . Apply conservative uncertainty

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# **Data Combination**

To evaluate the vacuum polarization contribution, we have to combine lots of experimental data.

To do so, we usually construct a  $\chi^2$  function and find the value of R(s) at each bin which minimizes  $\chi^2$ .

Naively, the  $\chi^2$  function defined as

$$\chi^2(\{\overline{R}_i\}) \equiv \sum_{n=1}^{N_{ ext{exp}}} \sum_{i=1}^{N_{ ext{bin}}} \sum_{j=1}^{N_{ ext{bin}}} (R_i^{(n)} - \overline{R}_i) (V_n^{-1})_{ij} (R_j^{(n)} - \overline{R}_j) \;,$$

where  $V_n$  is the cov. matrix of the *n*-th exp.,

$$V_{n,ij} = \begin{cases} (\delta R_{i,\text{stat}}^{(n)})^2 + (\delta R_{i,\text{sys}}^{(n)})^2 & (\text{for } i = j) \\ (\delta R_{i,\text{sys}}^{(n)})(\delta R_{j,\text{sys}}^{(n)}) & (\text{for } i \neq j) \end{cases}$$

may seem OK, but when there are non-negligible normalization uncertainties in the data, we have to be more careful.

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### $\chi^2$ vs normalization error: d'Agostini bias

G. D'Agostini, Nucl. Instrum. Meth. A346 (1994) 306 We first consider an observable x whose true value is 1. Suppose that there is an experiment which measures xand whose normalization uncertainty is 10%. Now, assume that this experiment measured x twice:

$$\begin{array}{ll} \mbox{1st result:} & 0.9\pm 0.1_{\rm stat}\pm 10\%_{\rm syst} \;, \\ \mbox{2nd result:} & 1.1\pm 0.1_{\rm stat}\pm 10\%_{\rm syst} \;. \end{array}$$

Taking the systematic errors 0.09 and 0.11, respectively, the covariance matrix and the  $\chi^2$  function are

$$egin{aligned} (\mathsf{cov.}) &= egin{pmatrix} 0.1^2 + 0.09^2 & 0.09 \cdot 0.11 \ 0.09 \cdot 0.11 & 0.1^2 + 0.11^2 \end{pmatrix} \ , \ \chi^2 &= egin{pmatrix} x - 0.9 & x - 1.1 \end{pmatrix} (\mathsf{cov.})^{-1} egin{pmatrix} x - 0.9 \ x - 1.1 \end{pmatrix} \ . \end{aligned}$$

 $\chi^2$  takes its minimum at x=0.98: Biased downwards!

### d'Agostini bias (2): improvement by iterations

What was wrong? In the previous page,

$$\begin{array}{ll} \mbox{1st result:} & 0.9\pm 0.1_{\rm stat}\pm 10\%_{\rm syst} \ , \\ \mbox{2nd result:} & 1.1\pm 0.1_{\rm stat}\pm 10\%_{\rm syst} \ . \end{array}$$

we took the syst. errors 0.09 and 0.11, respectively, which made the downward bias. Instead, we should take 10% of some estimator  $\bar{x}$  as the syst. errors. Then,

$$( ext{cov.}) = egin{pmatrix} 0.1^2 + (0.1ar{x})^2 & (0.1ar{x})^2 \ (0.1ar{x})^2 & 0.1^2 + (0.1ar{x})^2 \end{pmatrix} \,, \ \chi^2 = egin{pmatrix} x - 0.9 & x - 1.1 \end{pmatrix} ( ext{cov.})^{-1} egin{pmatrix} x - 0.9 \ x - 1.1 \end{pmatrix} \,.$$

 $\chi^2$  takes its minimum at x = 1.00: Unbiased! In more general cases, we use iterations: we find an estimator for the next round of iteration by  $\chi^2$ -minimization. R.D.Ball et al, JHEP 1005 (2010) 075.

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 $\sigma(e^+e^- 
ightarrow \pi^+\pi^-)$  data





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## $\sigma(e^+e^- ightarrow \pi^+\pi^-)$ : relative differences



## Contribution to $(g-2)_{\mu}$ from $\pi^+\pi^-$ channel



Fig. from KNT18, Phys. Rev. D97 (2018) 114025

#### Other notable exclusive channels [KNT18: arXiv:1802.02995, PRD (in press)]



Slide by A. Keshavarzi (Liverpool) at 'Muon g - 2 Workshop' at Mainz, June 18-22, 2018

## KNT18 $a_{\mu}^{had, VP}$ update



Slide by A. Keshavarzi (Liverpool) at 'Muon g - 2 HVP Workshop' at KEK, Feb. 12-14, 2018

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/ <mark>P: Hadronic Vacuum Polarization)</mark> . <mark>bL: Hadronic Light-by-Light)</mark> de by A. Keshavarzi (Liverpool) at 'Muon <i>g</i> — 2 Workshop' at Mainz, June 18-22, 2018						
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## Exp. value of muon g-2 vs SM prediction



KNT18 update

### Comparison with HLMNT11

Channel	This work (KNT18)	HLMNT11	Difference
$\pi^+\pi^-$	$502.99 \pm 1.97$	$505.77 \pm 3.09$	$-2.78\pm3.66$
$\pi^{+}\pi^{-}\pi^{0}$	$47.82 \pm 0.89$	$47.51 \pm 0.99$	$0.31 \pm 1.33$
$\pi^+\pi^-\pi^+\pi^-$	$15.17\pm0.21$	$14.65\pm0.47$	$0.52\pm0.51$
$\pi^+\pi^-\pi^0\pi^0$	$19.80\pm0.79$	$20.37 \pm 1.26$	$-0.57\pm1.49$
$K^+K^-$	$23.05\pm0.22$	$22.15\pm0.46$	$0.90 \pm 0.51$
$K_{S}^{0}K_{L}^{0}$	$13.05\pm0.19$	$13.33\pm0.16$	$-0.28\pm0.25$
Inclusive channel	$41.27\pm0.62$	$41.40\pm0.87$	$-0.13\pm1.07$
Total	$693.27 \pm 2.46$	$694.91 \pm 4.27$	$-1.64 \pm 4.93$

- $\Rightarrow \text{Biggest difference in } 2\pi \text{ channel} \\ \rightarrow \text{ large reduction in mean} \\ \text{ and uncertainty} \end{cases}$
- ⇒ Tensions with HLMNT11 analysis for both two-kaon channels
- $\Rightarrow$  Overall agreement with HLMNT11
- ⇒ Notable improvement of about one third in uncertainty



Slide by A. Keshavarzi (Liverpool) at 'Muon g = 2 HVP Workshop' at KEK, Feb. 12-14, 2018

### Comparison with other similar works

Channel	This work (KNT18)	DHMZ17	Difference
$\pi^+\pi^-$	$503.74 \pm 1.96$	$507.14 \pm 2.58$	$-3.40\pm3.24$
$\pi^{+}\pi^{-}\pi^{0}$	$47.70\pm0.89$	$46.20 \pm 1.45$	$1.50 \pm 1.70$
$\pi^+\pi^-\pi^+\pi^-$	$13.99\pm0.19$	$13.68\pm0.31$	$0.31 \pm 0.36$
$\pi^+\pi^-\pi^0\pi^0$	$18.15\pm0.74$	$18.03 \pm 0.54$	$0.12 \pm 0.92$
$K^+K^-$	$23.00\pm0.22$	$22.81 \pm 0.41$	$0.19\pm0.47$
$K_{S}^{0}K_{L}^{0}$	$13.04\pm0.19$	$12.82\pm0.24$	$0.22 \pm 0.31$
$1.8 \leq \sqrt{s} \leq 3.7~{ m GeV}$	$34.54\pm0.56~\mathrm{(data)}$	$33.45 \pm 0.65 \text{ (pQCD)}$	$1.09\pm0.86$
Total	$693.3 \pm 2.5$	$693.1 \pm 3.4$	$0.2 \pm 4.2$

- $\Rightarrow$  Total estimates from two analyses in very good agreement
- $\Rightarrow$  Masks much larger differences in the estimates from individual channels
- $\Rightarrow$  Unexpected tension for  $2\pi$  considering the data input likely to be similar
  - $\rightarrow$  Points to marked differences in way data are combined
  - $\rightarrow$  From  $2\pi$  discussion:  $a_{\mu}^{\pi^+\pi^-}$  (Weighted average) = 509.1  $\pm$  2.9
- $\Rightarrow$  Compensated by lower estimates in other channels

 $\rightarrow$  For example, the choice to use pQCD instead of data above 1.8 GeV

 $\Rightarrow$  FJ17:  $a_{\mu, \text{FJ17}}^{\text{had, LO VP}} = 688.07 \pm 41.4$ 

 $\rightarrow$  Much lower mean value, but in agreement within errors

Slide by A. Keshavarzi (Liverpool) at 'Muon g = 2 HVP Workshop' at KEK, Feb. 12-14, 2018

# **Comparison with Lattice Results**



• Lattice errors  $\sim 2\%$  vs phenomenology errors  $\sim 0.4\%$ 

• Some lattice results suggest new physics others not but all compatible with phenomenology

 Slide by L. Lellouch (Marseille) at 'Muon g-2 Workshop' at Mainz, June 18-22, 2018

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## Summary

- Standard Model prediction for  $(g-2)_{\mu}$ :  $\gtrsim 3.5\sigma$ deviation from measured value  $\implies$  New Physics?
- Recent data-driven evaluations of hadronic vacuum polarization contributions seem convergent (Similar mean values from KNT18 and Davier et al with slightly smaller uncertainty from KNT18.)
- To better establish the g-2 anomaly, better data for  $e^+e^- \rightarrow \pi^+\pi^-$  welcome (from Belle II, CMD-3, SND, . . .)
- Lattice calculations still suffer from large uncertainties (but a hybrid approach gives a slight improvement)
- New exp. at Fermilab and J-PARC expected to reduce the uncertainty of  $(g-2)_{\mu}$  by a factor of 4