

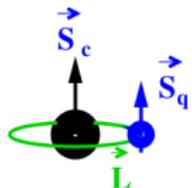


## Study of $\bar{B}^0 \rightarrow D^{*+} \omega \pi$ at Belle

Dmitry Matvienko

Novosibirsk State University, Novosibirsk, Russia  
and Budker Institute of Nuclear Physics, SB RAS

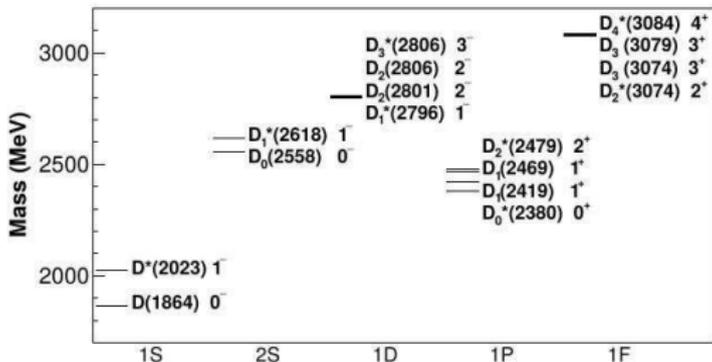
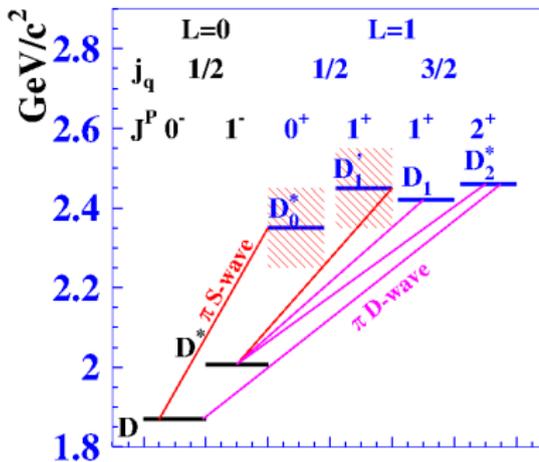
**Mini-workshop on  $D^{(*)} \tau \bar{\nu}$  and related topics  
Nagoya University, March 27-28, 2017**

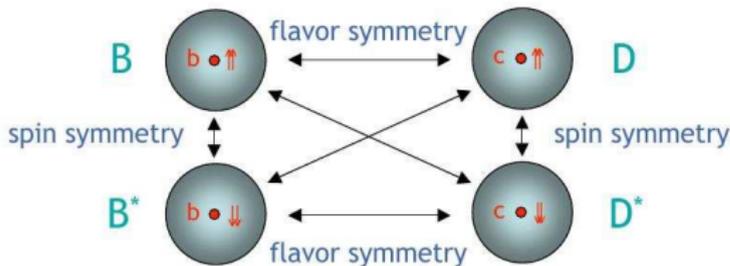
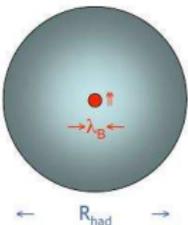


$$\vec{J}_q = \vec{L} + \vec{S}_q$$

$$\vec{J} = \vec{J}_q + \vec{S}_c$$

- $D_0^* \rightarrow D\pi$  in  $S$ -wave
- $D_1^* \rightarrow D^*\pi$  in  $S$ -wave
- $D_1 \rightarrow D^*\pi$  in  $D$ -wave
- $D_2^* \rightarrow D\pi$   $D^*\pi$  in  $D$ -wave



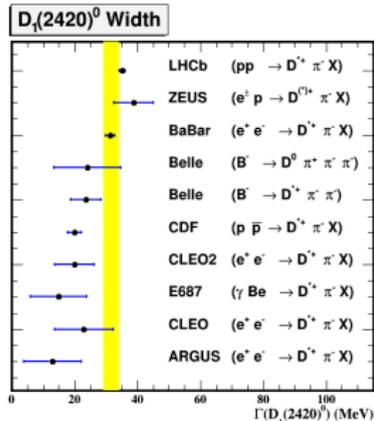
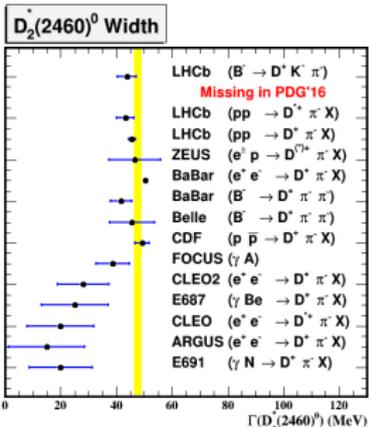
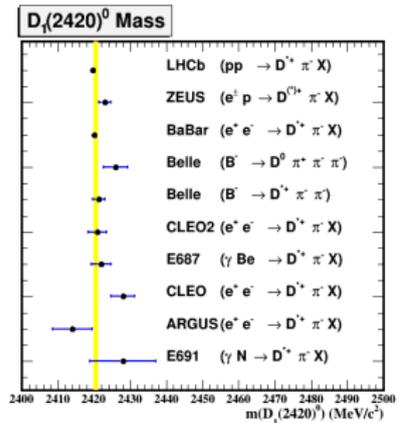
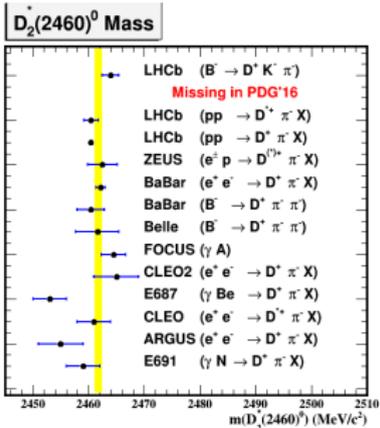


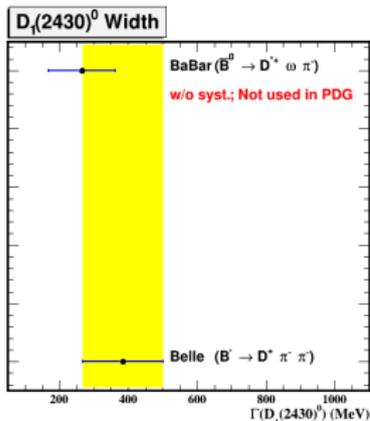
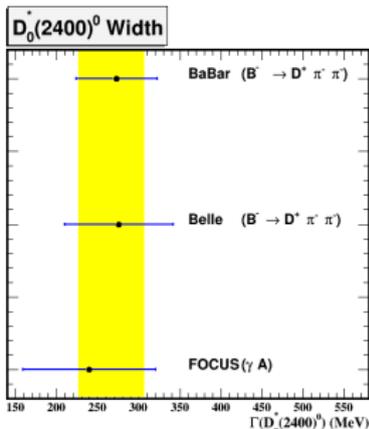
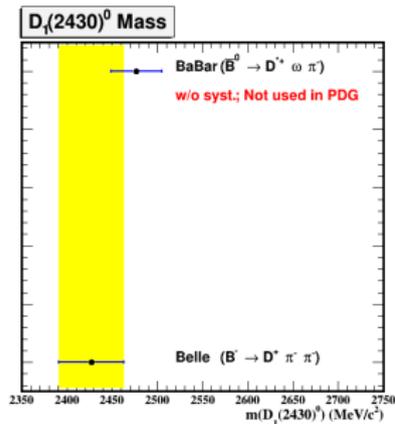
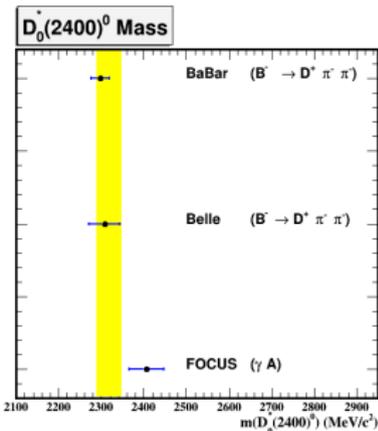
- Binding energy

$$M_H = m_Q + \bar{\Lambda} + \frac{1}{m_Q} \cdot (\lambda_1 + c_s \lambda_2) + O(1/m_Q^2)$$

- Power corrections to heavy-quark limit

- $\lambda_1$ : spin-independent
- $\lambda_2$ : spin-dependent



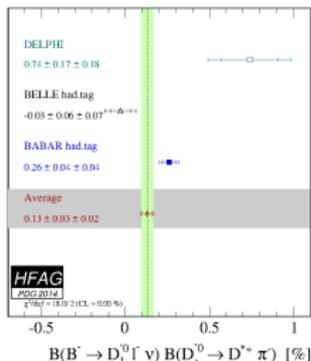
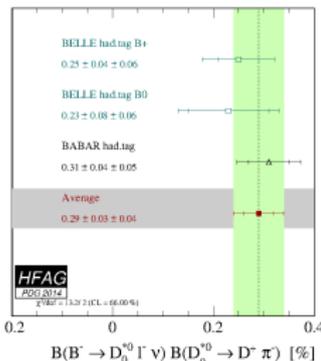
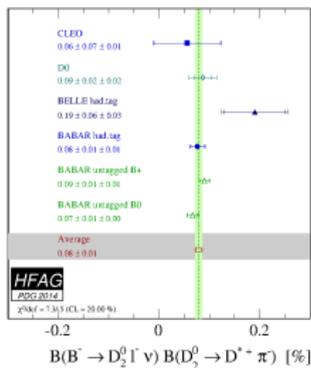
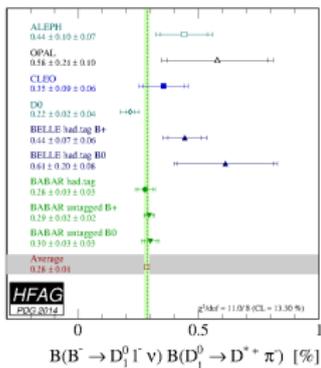


Mixing is suppressed in the heavy quark limit

$$|D_1(2420)^0 \rangle = \sin \omega |j_q = 1/2 \rangle + \cos \omega e^{-i\varphi} |j_q = 3/2 \rangle$$

$$|D_1^*(2430)^0 \rangle = \cos \omega |j_q = 1/2 \rangle - \sin \omega e^{i\varphi} |j_q = 3/2 \rangle$$

$$\left. \begin{aligned} \omega &= (-0.10 \pm 0.03 \pm 0.02 \pm 0.02) \text{ rad} \\ \varphi &= (+0.05 \pm 0.20 \pm 0.04 \pm 0.06) \text{ rad} \end{aligned} \right\} \text{Belle } (B^- \rightarrow D^{*+} \pi^- \pi^-)$$



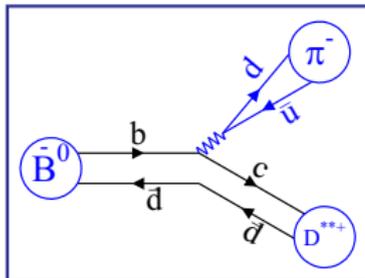
HQET vs. experiment discrepancy

Belle II data could clarify situation

Hadronic dynamics is welcome

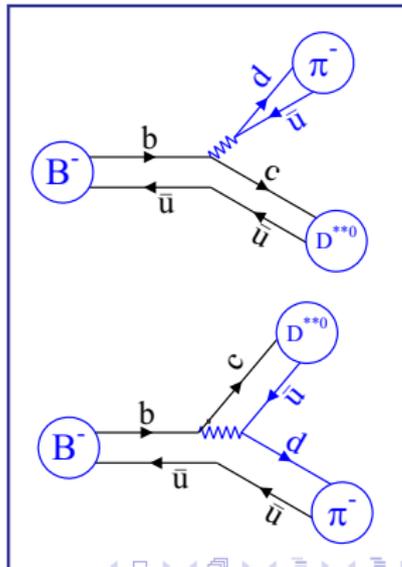
## Class I decays

Pion emission: " $B \rightarrow D^{**}$ "  $\times$   $f_\pi$   
 Isgur-Wise functions:  $\tau_{1/2} \ll \tau_{3/2}$   
 Narrow  $D^{**}$ 's dominate



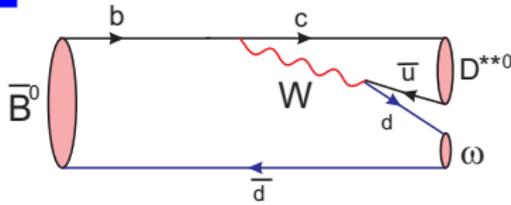
## Class III decays

Pion +  $D^{**}$  emission:  
 (" $B \rightarrow D^{**}$ "  $\times$   $f_\pi$ ) + (" $B \rightarrow \pi$ "  $\times$   $f_{D^{**}}$ )  
 Constructive interference  
 Isgur-Wise functions:  $\tau_{1/2} \ll \tau_{3/2}$   
 $D^{**}$  weak constants:  $f_{1/2} > f_{3/2}$   
 The rates of all  $D^{**}$ 's are of the same order

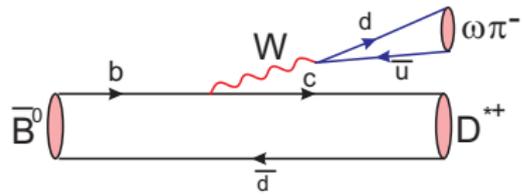


Class I & Class III BF products (PDG'16)	Value
$\mathcal{B}(\bar{B}^0 \rightarrow D_0^*(2400)^+ \pi^-) \times \mathcal{B}(D_0^*(2400)^+ \rightarrow D^0 \pi^+)$	$(6.0 \pm 3.0) \times 10^{-5}$
$\mathcal{B}(\bar{B}^0 \rightarrow D_1(2430)^+ \pi^-) \times \mathcal{B}(D_1(2430)^+ \rightarrow D^{*0} \pi^+)$	$< 7 \times 10^{-5}$
$\mathcal{B}(\bar{B}^0 \rightarrow D_1(2420)^+ \pi^-) \times \mathcal{B}(D_1(2420)^+ \rightarrow D^{*0} \pi^+)$	$(3.7 \pm 0.9) \times 10^{-4}$
$\mathcal{B}(\bar{B}^0 \rightarrow D_2^*(2460)^+ \pi^-) \times \mathcal{B}(D_2^*(2460)^+ \rightarrow D^0 \pi^+)$	$(2.2 \pm 0.4) \times 10^{-4}$
$\mathcal{B}(\bar{B}^0 \rightarrow D_2^*(2460)^+ \pi^-) \times \mathcal{B}(D_2^*(2460)^+ \rightarrow D^{*0} \pi^+)$	$(2.5 \pm 0.6) \times 10^{-4}$
$\mathcal{B}(\bar{B}^- \rightarrow D_0^*(2400)^0 \pi^-) \times \mathcal{B}(D_0^*(2400)^0 \rightarrow D^- \pi^+)$	$(6.4 \pm 1.4) \times 10^{-4}$
$\mathcal{B}(\bar{B}^- \rightarrow D_1(2430)^0 \pi^-) \times \mathcal{B}(D_1(2430)^0 \rightarrow D^{*-} \pi^+)$	$(5.0 \pm 1.3) \times 10^{-4}$
$\mathcal{B}(\bar{B}^- \rightarrow D_1(2420)^0 \pi^-) \times \mathcal{B}(D_1(2420)^0 \rightarrow D^{*-} \pi^+)$	$(6.8 \pm 1.5) \times 10^{-4}$
$\mathcal{B}(\bar{B}^- \rightarrow D_2^*(2460)^0 \pi^-) \times \mathcal{B}(D_2^*(2460)^0 \rightarrow D^- \pi^+)$	$(3.5 \pm 0.4) \times 10^{-4}$
$\mathcal{B}(\bar{B}^- \rightarrow D_2^*(2460)^0 \pi^-) \times \mathcal{B}(D_2^*(2460)^0 \rightarrow D^{*-} \pi^+)$	$(2.2 \pm 1.1) \times 10^{-4}$

**Class I & III data confirm the theoretical expectations**  
**Class II data are needed**



**Class II decays**



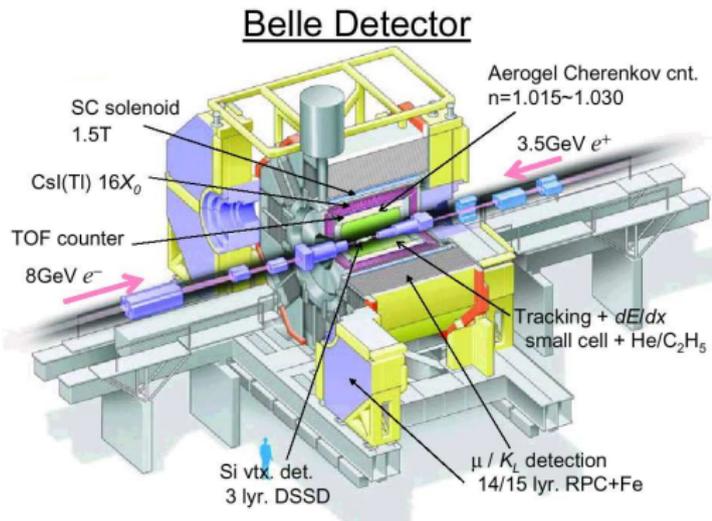
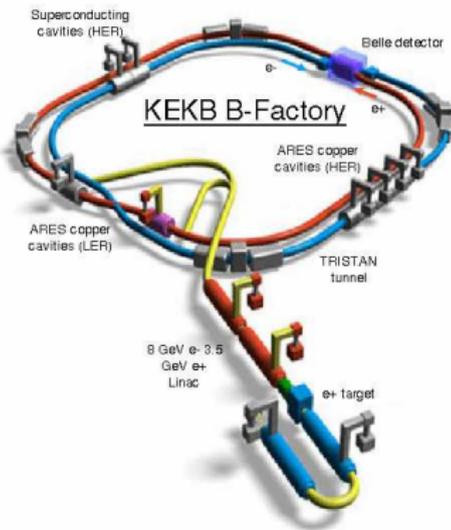
- 1. What are the  $j = 1/2$  and  $j = 3/2$  rates in class II decays?
- 2. What is the tensor  $j = 3/2$  rate in class II decays?

Soft Collinear Effective Theory (SCET) (Phys. Rev. D **70**, 114006 (2004)) predicts equal rates and strong phases for  $D_1(2420)$  and  $D_2^*(2460)$  ( $j = 3/2$ ) in the  $\bar{B}^0 \rightarrow D^{*0} M$  decays,  $M = \pi, \rho, K$  or  $M = K^*$  with long. polarization.

### SCET vs naive factorization (?)

### Advantage of $\bar{B}^0 \rightarrow D^{*0} \omega$ study

- Lower fractions of  $q\bar{q}$  continuum and combinatorial  $B\bar{B}$  backgrounds than in  $B \rightarrow D^{*0} \pi^0$ .
- Possibility to measure the polarizations and partial-wave fractions of  $D^{*0}$ -states.
- Possibility to perform the coherent study of  $\rho(770)$ ,  $\rho(1450)$  and  $\rho(1700)$  in the  $\omega\pi$  final state.



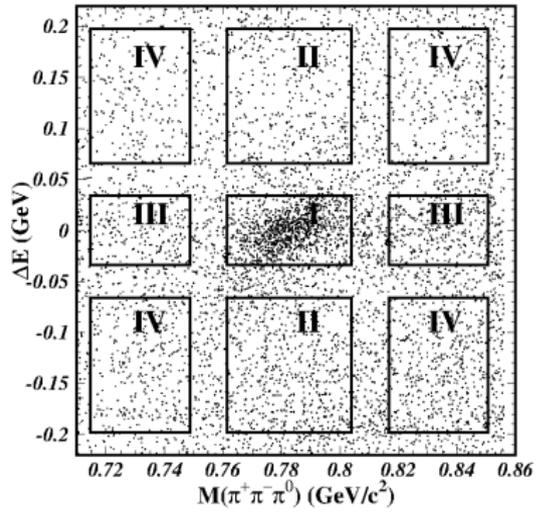
- Asymmetric beam energies:  $E_{e^-} = 8 \text{ GeV}$ ,  $E_{e^+} = 3.5 \text{ GeV}$
- Peak luminosity:  $\mathcal{L} = 2.1 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$
- $B\bar{B}$  luminosity:  $\int \mathcal{L} dt \approx 711 \text{ fb}^{-1}$  corresponds to  $(772 \pm 11) \times 10^6 B\bar{B}$  pairs produced.
- Reported analysis uses the full  $B\bar{B}$  data sample collected at Belle.

$$M_{bc} = \sqrt{(E_{\text{beam}}^{\text{CMS}})^2/c^4 - \left| \sum_i \mathbf{p}_i^{\text{CMS}} \right|^2/c^2}$$

$$\Delta E = \sqrt{|\mathbf{p}_i^{\text{CMS}}|^2 c^2 + m_i^2 c^4} - E_{\text{beam}}^{\text{CMS}}$$

$$5.2725 \text{ GeV}/c^2 < M_{bc} < 5.2845 \text{ GeV}/c^2$$

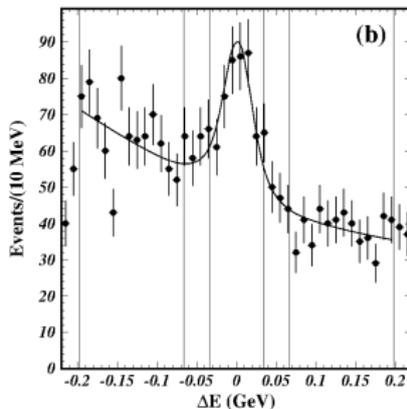
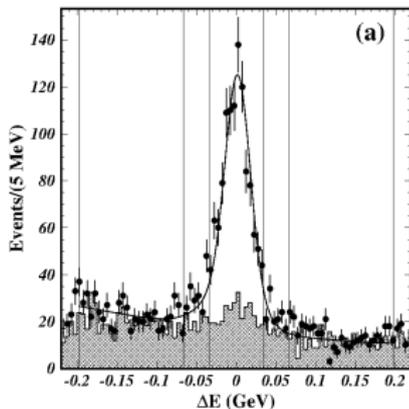
Region I — signal  
 Regions II, III, IV — sidebands



## Multidimensional amplitude analysis

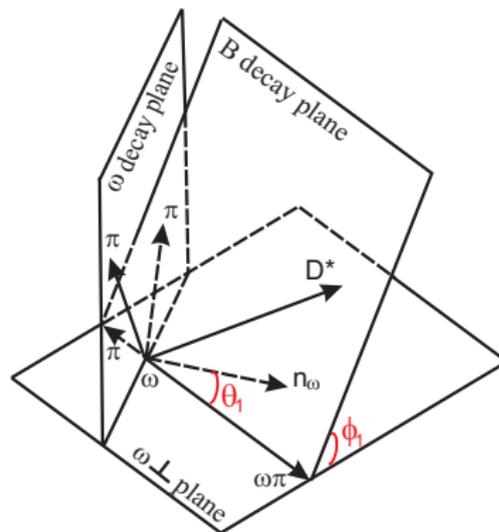
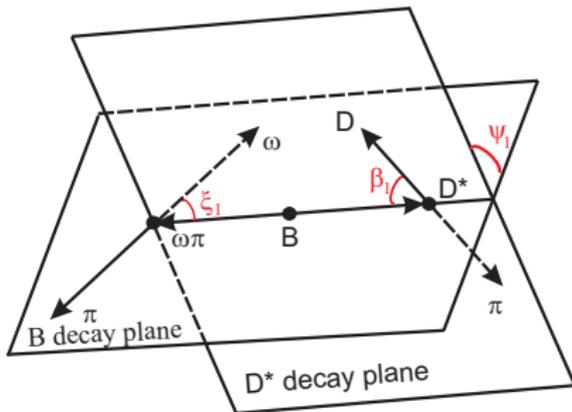
- Amplitude analysis as an efficient tool for distinguishing the contributions of narrow & broad  $D^{**}$  states.
- Unbinned likelihood fit in the decay kinematic phase space.
- Likelihood function constructed from the background and signal PDF functions.

Peaking bkg from  $\bar{B}^0 \rightarrow D^{*+} \pi^+ \pi^- \pi^0 \pi^-$

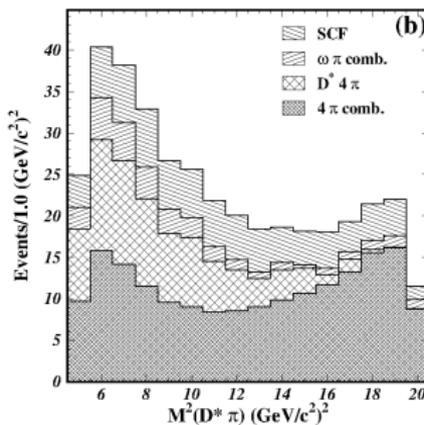
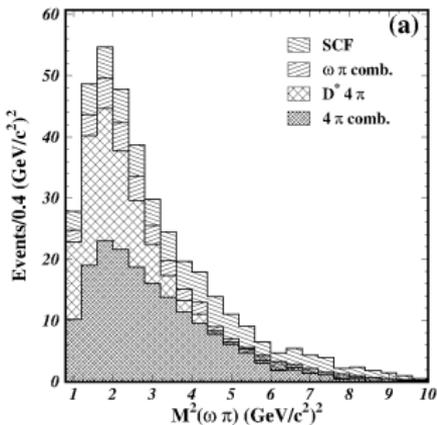


$$\mathcal{B}(\bar{B}^0 \rightarrow D^{*+} \omega \pi^-) = (2.31 \pm 0.11 \text{ (stat.)} \pm 0.14 \text{ (syst.)}) \times 10^{-3}$$

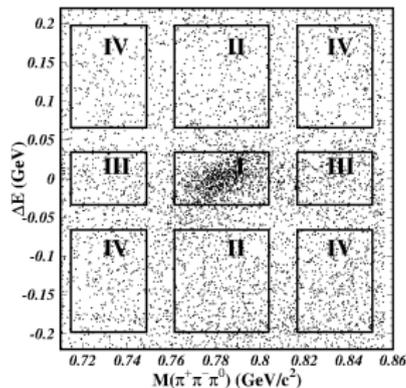
BF consistent with CLEO and BaBar but has higher precision.



- 5 angles +  $q^2 = 6$  variables
- Two different bases for  $\omega\pi$  and  $D^{**}$
- Full angular analysis for all intermediate states



- SCF (Self Cross-Feed) (regions I, II, III and IV) — misreconstruction of true signal
- $\omega\pi$  comb. (regions I and III) — combinatorial  $B\bar{B}$  bkg w/ correctly reconstructed  $\omega$
- $D^*4\pi$  (regions I and II) — peaking bkg from  $\bar{B}^0 \rightarrow D^{*+}\pi^+\pi^-\pi^0\pi^-$  w/o  $\omega$  contribution (basically  $\bar{B}^0 \rightarrow D^{*+}a_1(1260)\pi^-$ )
- $4\pi$  comb. (regions I, II, III and IV) — combinatorial  $B\bar{B}$  bkg w/ misreconstructed  $\omega$



- Sum of quasi-two-body resonant amplitudes:

$$\mathcal{M} = \sum_{i=R} a_R e^{i\phi_R} \mathcal{M}_R$$

- Partial-wave formalism for each amplitude:

$$\mathcal{M}_R \sim \frac{1}{\mathcal{D}_R(q^2)} \sum_{L_1} f_{L_1} \sum_{L_2} \mathcal{A}_{L_1, L_2}; \quad L_1 = L(D^{**}\omega), \quad L_2 = L(D^*\pi)$$

- Relativistic Breit-Wigner function with  $q^2$ -dependent width for all resonances:

$$\mathcal{D}_R(q^2) = q^2 - m_R^2 + im_R \Gamma_R(q^2)$$

- Angular distributions in terms of defined variables:

$$\mathcal{A}_{L_1=S, L_2=P} = -s_\theta s_\phi c_\beta s_\xi + s_\theta c_\phi s_\beta s_\psi - s_\theta s_\phi s_\beta c_\psi c_\xi$$

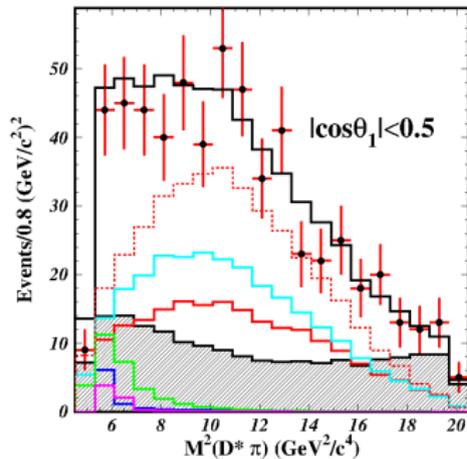
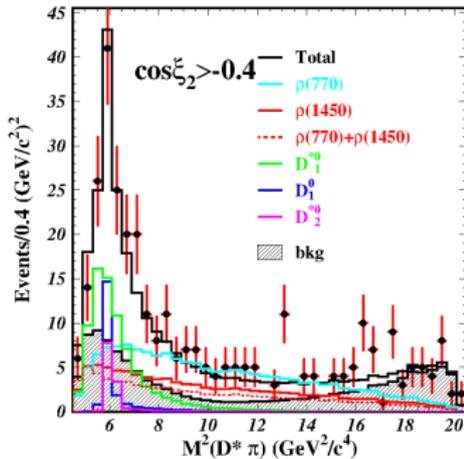
$$\mathcal{A}_{L_1, L_2} = \dots$$

- Blatt-Weisskopf factor for each partial wave:  $f_{L_1}(q^2) \sim B_{L_1}(q^2)$
- Mixing between  $j_q = 1/2$  and  $j_q = 3/2$

No significant mixing observed:

$$\omega = -0.03 \pm 0.02 \text{ (stat.)},$$

$$\varphi = -0.27 \pm 0.75 \text{ (stat.)}$$

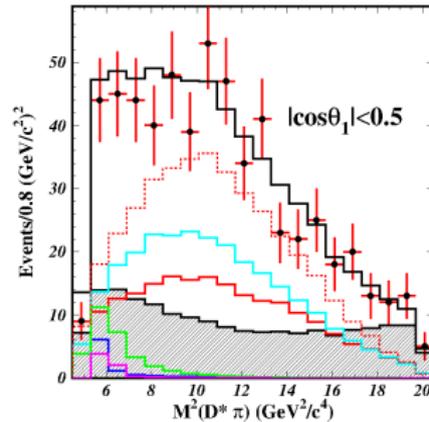
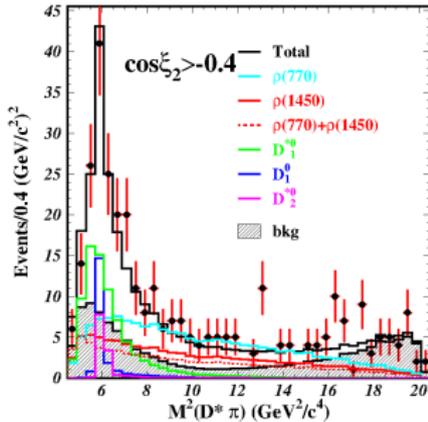


$$\mathcal{B}_{D_1(2430)^0} = (2.5 \pm 0.4 \text{ (stat.)}^{+0.7}_{-0.2} \text{ (syst.)}^{+0.4}_{-0.1} \text{ (model)}) \times 10^{-4}$$

$$\delta = 8.6\sigma$$

• Consistent with HQET  $\Rightarrow \mathcal{B}_{D_1(2430)^0} \simeq 3.5 \times 10^{-4}$

# $D_1(2420)^0$ & $D_2^*(2460)^0$ resonances



$$\mathcal{B}_{D_1(2420)^0} = (0.7 \pm 0.2 \text{ (stat.)}^{+0.1}_{-0.0} \text{ (syst.)} \pm 0.1 \text{ (model)}) \times 10^{-4}$$

$$\delta = 5.5\sigma$$

$$\mathcal{B}_{D_2^*(2460)^0} = (0.4 \pm 0.1 \text{ (stat.)}^{+0.0}_{-0.1} \text{ (syst.)} \pm 0.1 \text{ (model)}) \times 10^{-4}$$

$$\delta = 5.0\sigma$$

- Narrow  $D^{**}$  are suppressed according to HQET
- SCET prediction could be reasonable  
 $\mathcal{B}(\bar{B}^0 \rightarrow D_1(2420)^0 \omega_{\parallel}) = \mathcal{B}(\bar{B}^0 \rightarrow D_2^*(2460)^0 \omega_{\parallel})$   
 (but only for the longitudinal component)

## $D^{**}$ longitudinal polarizations

Long-distance QCD effects could be essential

$$\mathcal{P}_{D_1(2430)^0} = (63.0 \pm 9.1 \text{ (stat.)} \pm 4.6 \text{ (syst.)} {}^{+4.6}_{-3.9} \text{ (model)})\%$$

$$\mathcal{P}_{D_1(2420)^0} = (67.1 \pm 11.7 \text{ (stat.)} {}^{+0.0}_{-4.2} \text{ (syst.)} {}^{+2.3}_{-2.8} \text{ (model)})\%$$

$$\mathcal{P}_{D_2^*(2460)^0} = (76.0 {}^{+18.3}_{-8.5} \text{ (stat.)} \pm 2.0 \text{ (syst.)} {}^{+2.9}_{-2.0} \text{ (model)})\%$$

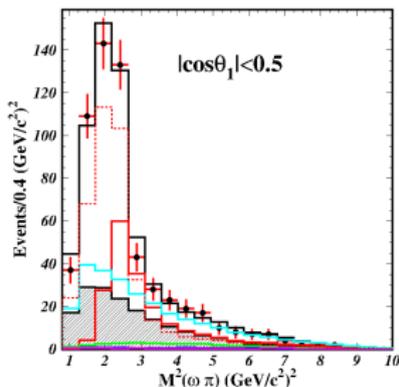
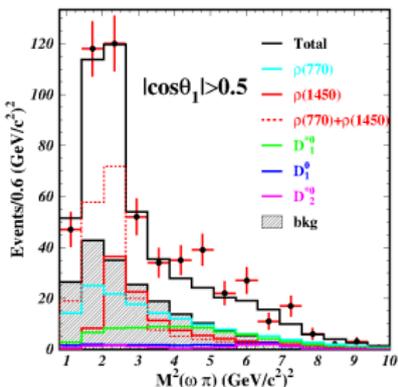
$$\mathcal{P}_{D^{*0}} = (66.5 \pm 4.7 \text{ (stat.)} \pm 1.5 \text{ (syst.)})\% \quad \leftarrow \text{BaBar } (\bar{B}^0 \rightarrow D^{*0} \omega)$$

## $S$ -, $P$ -, $D$ -wave rates for $D_1(2430)^0$

$$f_S = (38.9 \pm 10.8 \text{ (stat.)} {}^{+4.3}_{-0.7} \text{ (syst.)} {}^{+1.2}_{-1.1} \text{ (model)})\%$$

$$f_P = (33.1 \pm 9.5 \text{ (stat.)} {}^{+2.4}_{-5.5} \text{ (syst.)} {}^{+3.0}_{-4.0} \text{ (model)})\%$$

$$f_D = (28.3 \pm 8.9 \text{ (stat.)} {}^{+3.0}_{-0.8} \text{ (syst.)} {}^{+3.9}_{-2.9} \text{ (model)})\%$$



$$\mathcal{B}_{\rho(770)^-} = (1.48 \pm 0.27 \text{ (stat.)}^{+0.15}_{-0.09} \text{ (syst.)}^{+0.21}_{-0.56} \text{ (model)}) \times 10^{-3}$$

$$\delta = 10.5\sigma$$

$$\mathcal{B}_{\rho(1450)^-} = (1.07^{+0.15}_{-0.31} \text{ (stat.)}^{+0.06}_{-0.13} \text{ (syst.)}^{+0.40}_{-0.02} \text{ (model)}) \times 10^{-3}$$

$$\delta = 15.0\sigma$$

$$\mathcal{B}_{\Sigma\rho} = (1.90 \pm 0.11 \text{ (stat.)}^{+0.11}_{-0.13} \text{ (syst.)}^{+0.02}_{-0.06} \text{ (model)}) \times 10^{-3}$$

$$\delta = 29.8\sigma$$

$$\mathcal{B}_{b_1(1235)} < 0.7 \times 10^{-4} \text{ (90\% C.L.)}$$

$$\mathcal{P}_{\rho(1450)} = (66.5 \pm 0.6 \text{ (stat.)}^{+0.1}_{-0.3} \text{ (syst.)}^{+1.2}_{-0.8} \text{ (model)})\%$$

HQET & factorization prediction based on SL data  $\implies (68.4 \pm 0.9)\%$

- $\bar{B}^0 \rightarrow D^{*+} \omega \pi^-$  decay rate is consistent with the CLEO and BaBar measurements but it has higher precision (7.7%).
- Broad  $D_1(2430)^0$  production is consistent with HQET.
- Narrow  $D_1(2420)$  and  $D_2^*(2460)$  productions are suppressed as predicted in HQET
- Significant tensor  $D_2^*(2460)$  production is observed.
- $D^{**}$  longitudinal polarizations and partial-wave rates are measured.
- The consistent study of the  $\rho$ -meson-like states is performed.

**These results could be useful for SL studies**

