Classification of Simple W' Models

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What is W'?

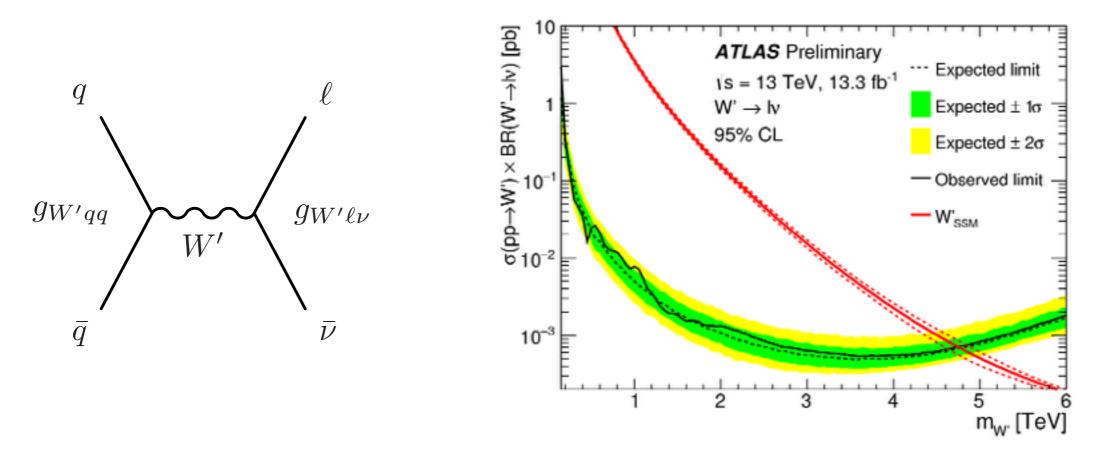
Hypothetical particle

- spin 1
- heavier version of W boson

Motivated by various models beyond the standard model

- dynamical electroweak symmetry breaking
- Extra dimension models (KK gauge boson)
- Unification models (W_R, \cdots)
- • • •

current bound W'

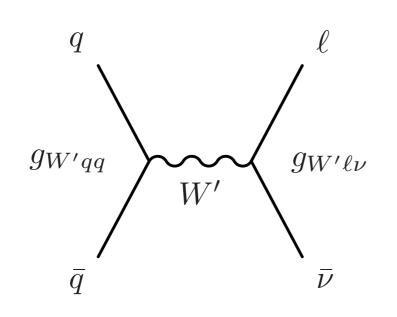


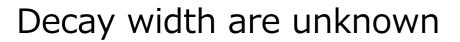
 $m_{W'} > 4.7 \text{ TeV}$ with the assumiptions, $g_{W'ff} = g_{Wff}^{SM}$

This bound depends on W' couplings.

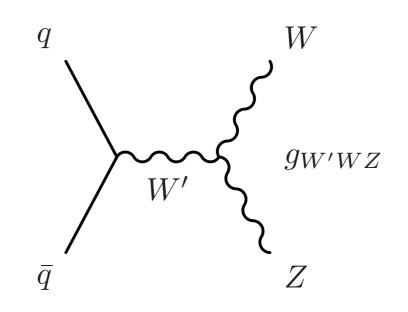
However, couplings are unknown

W' couplings are unknown parameters





$$\begin{split} \Gamma(W' \to WZ) &\simeq \frac{1}{192\pi} \frac{m_{W'}^5}{m_W^2 m_Z^2} g_{WW'Z}^2, \\ \Gamma(W' \to u_i d_j) &\simeq \frac{1}{16\pi} |V_{CKM}^{ij}|^2 m_{W'} g_{W'}^2, \\ \Gamma(W' \to \ell \nu) &\simeq \frac{1}{48\pi} m_{W'} g_{W'}^2, \end{split}$$



However, the ratio could be predicted, if we know relations among couplings

$$\frac{\Gamma(W' \to WZ)}{\Gamma(W' \to \ell\nu)} \simeq \frac{m_{W'}^4}{4m_W^2 m_Z^2} \frac{g_{WW'Z}^2}{g_{W'}^2}$$

We can predict the main decay mode of W'

This work

Goal

 \bullet find the main decay mode of W' in a model independent manner

Strategy

- perturbative unitarity
- find relations among various couplings
- compare widths $\Gamma(W' \rightarrow ff)$ and $\Gamma(W' \rightarrow WZ)$
- find the main decay mode of W'

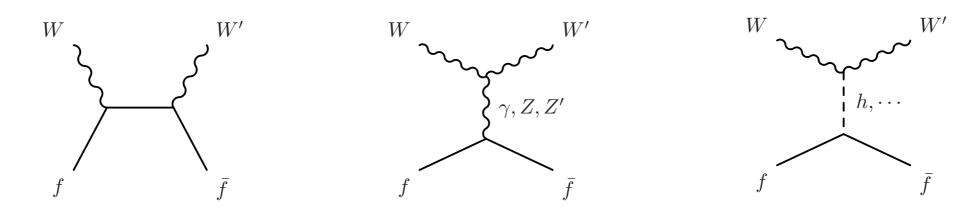
I will discuss Two setups

(1) SM + W' + Z' + CP-even scalars

- minimal flavor violation is assumed
- arbitrary number of CP-even scalars

$$\begin{split} g_{\bar{u}_{i}d_{j}V^{+}}^{L} &= V_{CKM}^{ij}g_{V}, \quad g_{\bar{\nu}_{i}\ell_{j}V^{+}}^{L} = \delta^{ij}g_{V}^{\ell}, \\ g_{\bar{u}_{i}u_{j}V^{0}}^{L,R} &= \delta^{ij}g_{\bar{u}uV^{0}}^{L,R}, \quad g_{\bar{d}_{i}d_{j}V^{0}}^{L,R} = \delta^{ij}g_{\bar{d}dV^{0}}^{L,R}, \quad g_{\bar{\ell}_{i}\ell_{j}V^{0}}^{L,R} = \delta^{ij}g_{\bar{\ell}\ell V^{0}}^{L,R}, \\ g_{\bar{u}_{i}u_{j}h} &= \delta^{ij}g_{\bar{u}uh}, \quad g_{\bar{d}_{i}d_{j}h} = \delta^{ij}g_{\bar{d}dh}, \quad g_{\bar{\ell}_{i}\ell_{j}h} = \delta^{ij}g_{\bar{\ell}\ell h}. \end{split}$$

(2) SM + W' + Z' + CP-even scalars + CP-odd scalars



amplitude **M**_{ij} at high energy limit (i, j are the twice of the helicity of f and f-bar)

$$\mathcal{M}_{-+} \simeq \frac{s}{2m_W m_{W'}} \mathcal{A} \sin \theta,$$

$$\mathcal{M}_{+-} \simeq \frac{s}{2m_W m_{W'}} \mathcal{B} \sin \theta,$$

$$\mathcal{M}_{++} \simeq \frac{\sqrt{s}}{2m_W m_{W'}} \left(\mathcal{C}^{(0)} + \mathcal{C}^{(1)} \cos \theta \right),$$

$$\mathcal{M}_{--} \simeq \frac{\sqrt{s}}{2m_W m_{W'}} \left(\mathcal{D}^{(0)} + \mathcal{D}^{(1)} \cos \theta \right),$$

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We impose A = 0, B = 0, $C^{(0)} = 0$, $D^{(0)} = 0$, $C^{(1)} = 0$, $D^{(1)} = 0$

After taking linear combination, we find four independent relations

$$\frac{1}{2}g_W g_{W'} = \sum_{V=Z,Z'} g_{WW'V} g_{\bar{u}uV}^L,$$

$$0 = \sum_{V=Z,Z'} g_{WW'V} g_{\bar{u}uV}^R,$$

$$g_W g_{W'} = 2 \sum_h g_{WW'h} \frac{g_{\bar{u}uh}}{m_u},$$

$$0 = \sum_{V=Z,Z'} \frac{m_W^2 - m_{W'}^2}{m_V^2} g_{WW'V} (g_{\bar{u}uV}^L - g_{\bar{u}uV}^R)$$

By combining sum rules, we find

$$g_{WW'Z} \simeq -\frac{m_W m_Z}{m_{Z'}^2} g_{W'}$$

$$(m_{W'} \simeq m_{Z'} \gg m_W, m_Z)$$

W' decay width

$$\Gamma(W' \to WZ) \simeq \frac{1}{192\pi} \frac{m_{W'}^5}{m_W^2 m_Z^2} g_{WW'Z}^2,$$

$$\Gamma(W' \to u_i d_j) \simeq \frac{1}{16\pi} |V_{CKM}^{ij}|^2 m_{W'} g_{W'}^2,$$

$$\Gamma(W' \to \ell\nu) \simeq \frac{1}{48\pi} m_{W'} g_{W'}^2,$$

ratio

$$\frac{\Gamma(W' \to WZ)}{\Gamma(W' \to f_i f_j)} \simeq \frac{1}{4c_{ij}}$$
$$c_{ij} = \begin{cases} N_c |V_{CKM}^{ij}|^2 & \text{(for quarks)}\\ \delta^{ij} & \text{(for leptons)} \end{cases}$$

We find

$$\frac{\Gamma(W' \to WZ)}{\sum_{f} \Gamma(W' \to f_i f_j)} \simeq \frac{1}{4 \times ((N_c + 1) \times 3)} = \frac{1}{48},$$

and thus

$$Br(W' \to WZ) = \frac{\Gamma(W' \to WZ)}{\Gamma(W' \to WZ) + \sum_{f} \Gamma(W' \to ff) + \sum_{X} \Gamma(W' \to X)}$$
$$\leq \frac{\Gamma(W' \to WZ)}{\Gamma(W' \to WZ) + \sum_{f} \Gamma(W' \to ff)}$$
$$\simeq 2\%$$

 $\Gamma(W' \to X)$ is the sum of the other partial decay widths of W' such as $\Gamma(W' \to Wh)$

short summary and next step

SM + W' + Z' + CP-even scalars

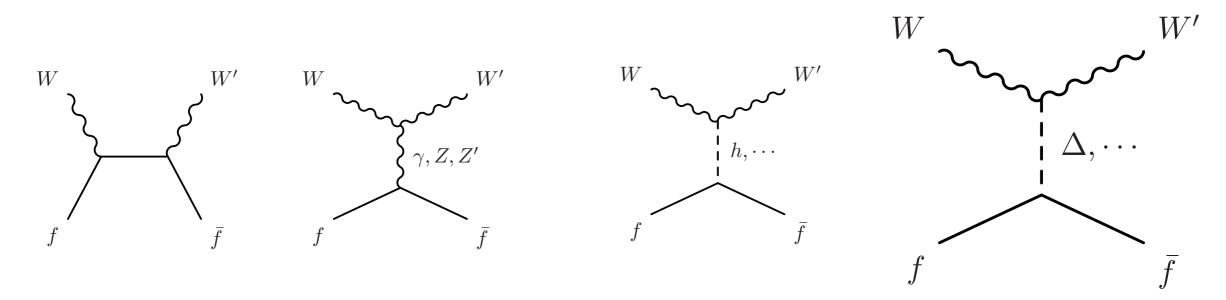
- unitarity analysis
- $Br(W' \rightarrow WZ) < 2\%$
- W' \rightarrow fermions is the main decay mode
- independent of the number of CP-even scalars

If W' is discovered in future through W' \rightarrow WZ, then it implies Br(W' \rightarrow WZ) >> 2%. Is it possible?

A possible way to make $Br(W' \rightarrow WZ) > 2\%$

- add other interactions and/or particles
- (or unitarity violation)

Adding CP-odd scalars Δ



coupling relations

$$\frac{1}{2}g_{W}g_{W'} = \sum_{V=Z,Z'} g_{WW'V}g_{\bar{u}uV}^{L},$$

$$0 = \sum_{V=Z,Z'} g_{WW'V}g_{\bar{u}uV}^{R},$$

$$g_{W}g_{W'} = 2\sum_{h} g_{WW'h}\frac{g_{\bar{u}uh}}{m_{u}}$$

$$\sum_{\Delta^{0}} \frac{g_{\bar{u}u\Delta^{0}}}{m_{u}}g_{WW'\Delta} = \sum_{V=Z,Z'} 2g_{WW'V}\frac{m_{W}^{2} - m_{W'}^{2}}{m_{V}^{2}}(g_{\bar{u}uV}^{L} - g_{\bar{u}uV}^{R})$$

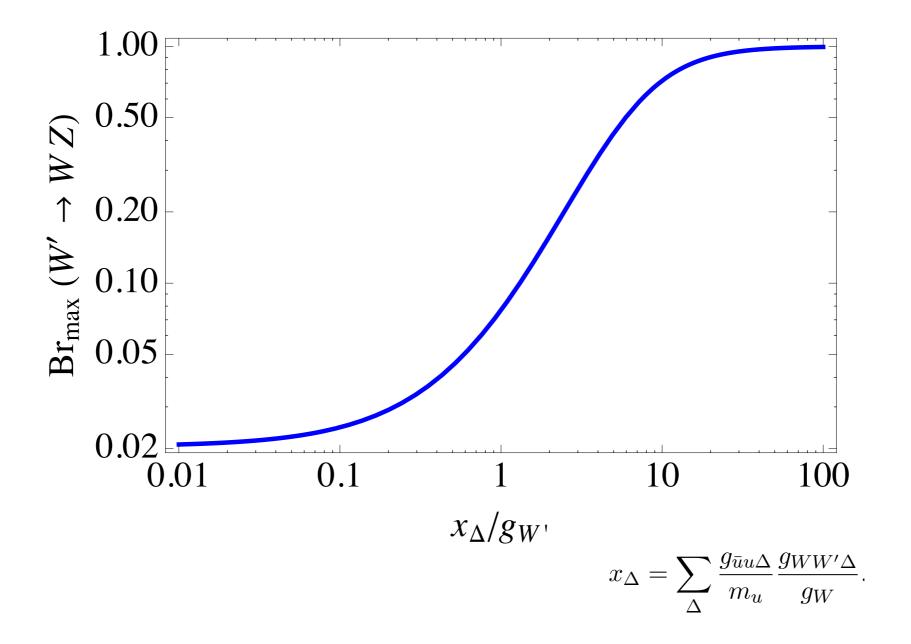
Adding CP-odd scalars Δ

$$g_{WW'Z} \simeq -\frac{m_W m_Z}{m_{Z'}^2} \left(g_{W'} + x_\Delta\right)$$
$$x_\Delta = \sum_\Delta \frac{g_{\bar{u}u\Delta}}{m_u} \frac{g_{WW'\Delta}}{g_W}$$

We can estimate the branching ratio

$$Br(W' \to WZ) = \frac{\Gamma(W' \to WZ)}{\Gamma(W' \to WZ) + \sum_{f} \Gamma(W' \to ff) + \sum_{X} \Gamma(W' \to X)}$$
$$\leq \frac{\Gamma(W' \to WZ)}{\Gamma(W' \to WZ) + \sum_{f} \Gamma(W' \to ff)}$$

Br(W'→WZ)>2% thanks to CP-odd scalars



Summary

W' is a popular particle in physics beyond the standard model

- dynamical electroweak symmetry breaking
- Extra dimension models (KK gauge boson)
- Unification models (W_R, \cdots)

We have studied W' in a model independent manner

- ★ perturbative unitarity
- ★ coupling relations

Setup 1: SM + W' + Z' + CP-even scalars

★ Br(W' → WZ) < 2%

Setup 2: SM + W' + Z' + CP-even scalars + CP-odd scalars

★ Br(W' \rightarrow WZ) can be >> 2%