

Status and prospect of R measurements at Belle II

Yuki Sue

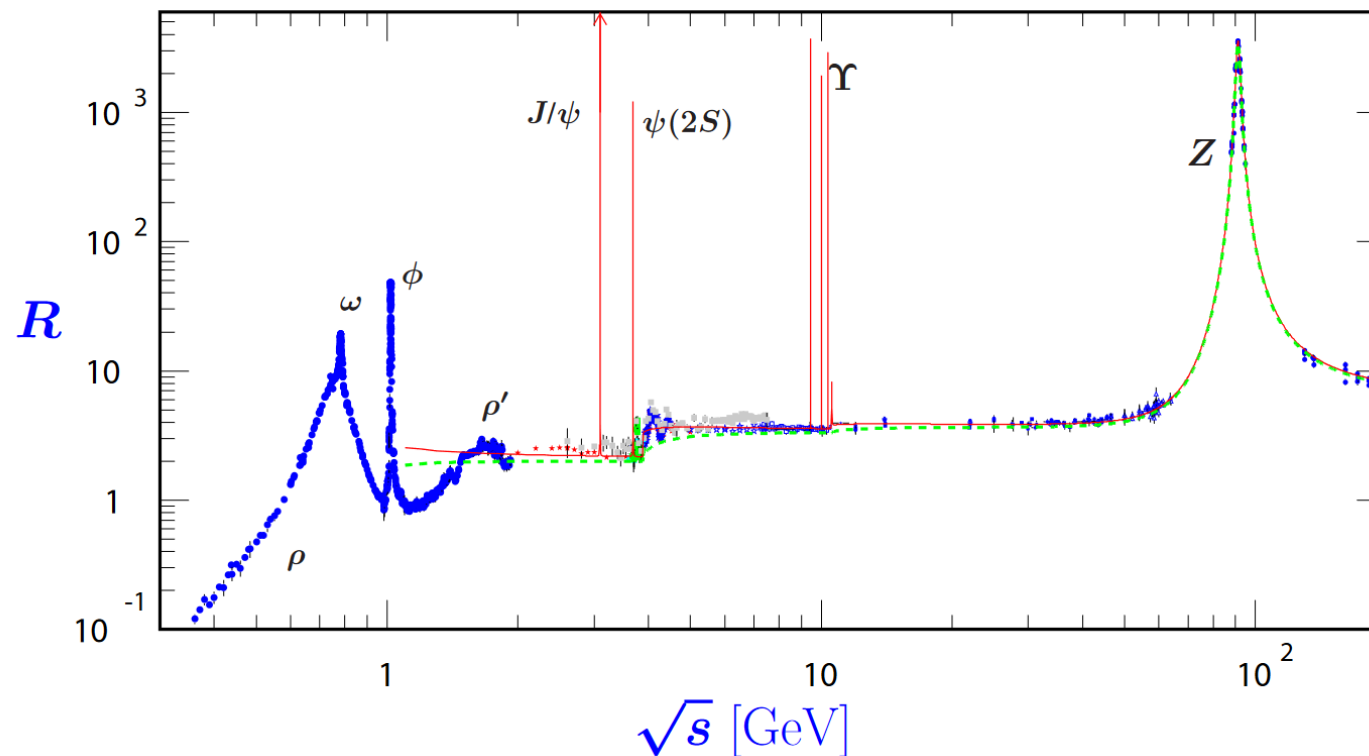
Nagoya University

B1 Heavy Flavor and Dark Matter Joint Unit Symposium, March 29th, 2023

Hadronic cross section measurement

$$R(s) = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$$

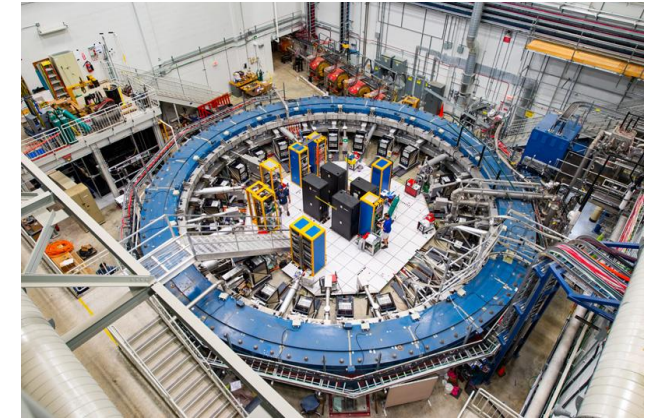
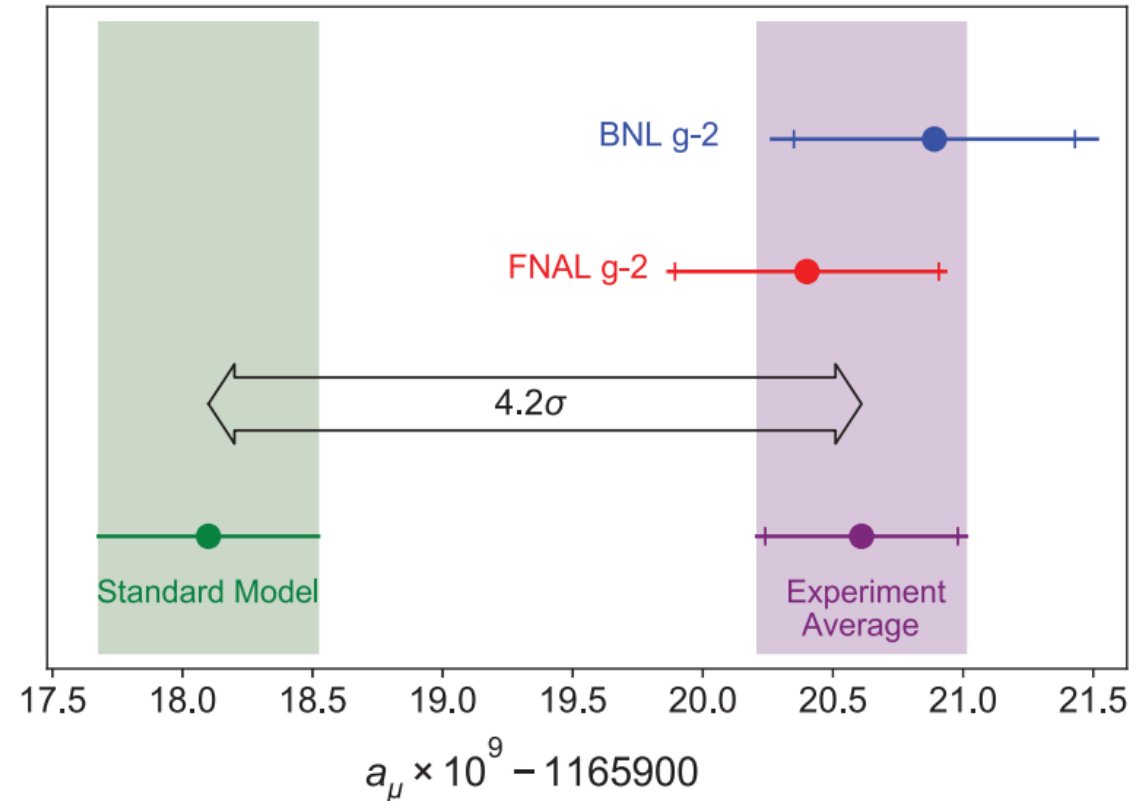
- The hadronic cross section data, R-ratio, is crucial inputs to the QCD physics.
- Tests of perturbative QCD, light vector mesons parameters, hadronic corrections to $\alpha(M_Z^2)$, muon g-2, ...



Introduction for muon g-2

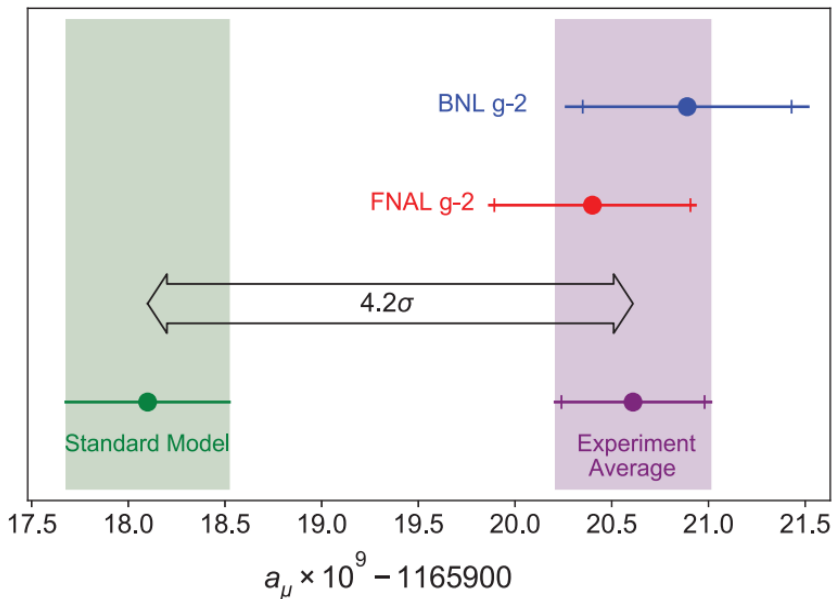
B. Abi *et al.*, PRL126, 141801 (2021).

T. Aoyama *et al.*, Phys. Rept. 887 (2020).



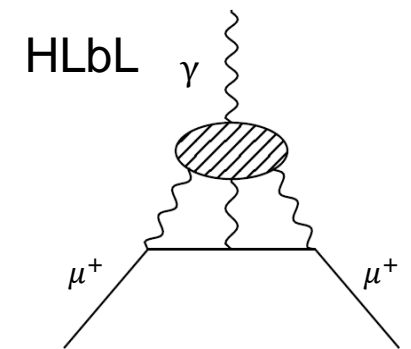
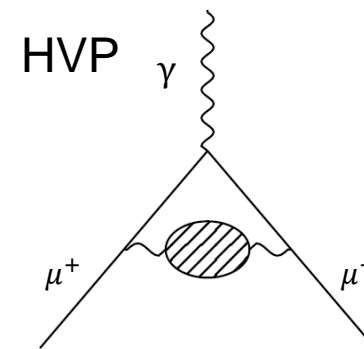
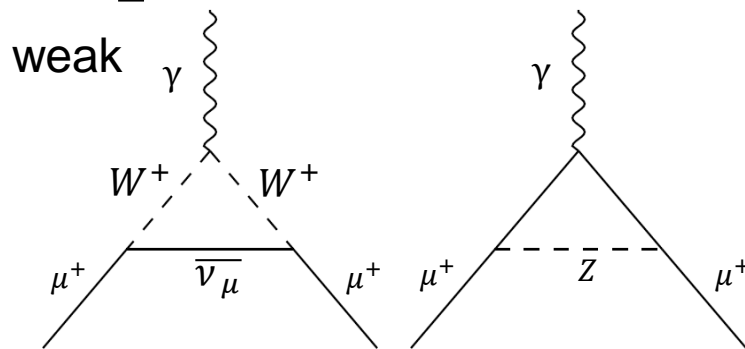
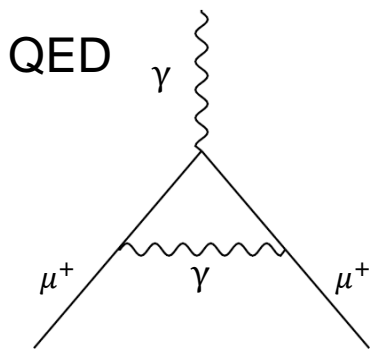
Muon anomalous magnetic moment $a_\mu = \frac{g_\mu - 2}{2}$

Introduction for muon g-2



Contribution	Value $a_\mu \times 10^{11}$	Error $\delta a_\mu \times 10^{11}$
QED	116 584 718.931	0.104
HVP LO (Leading-Order)	6931	40
HVP HO (Higher-Order)	-85.9	1.2
HLbL (Light-by-Light)	92	19
EW (Electroweak)	153.6	1
SM total (Dispersive)	116591810	43
Experiment (BNL+FNAL)	116592061	41
Experiment – SM	251	59

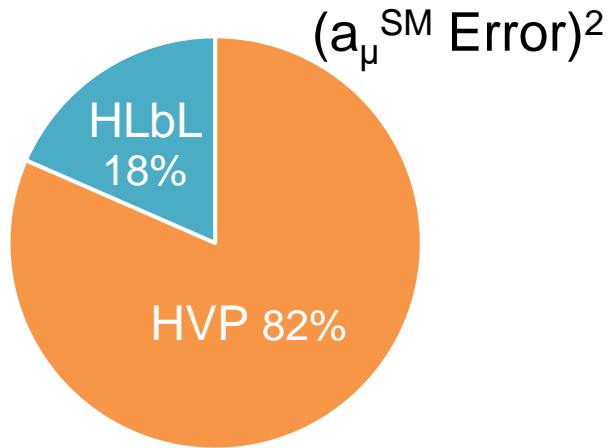
$$a_\mu^{\text{SM}} = \frac{g-2}{2} = a_\mu^{\text{QED}} + a_\mu^{\text{weak}} + a_\mu^{\text{HVP}} + a_\mu^{\text{HLbL}}$$



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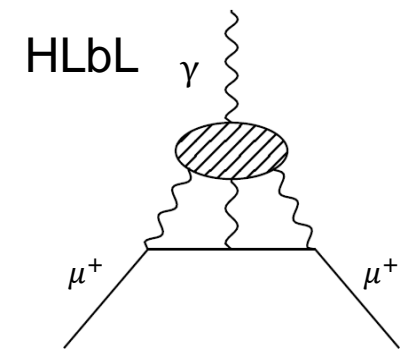
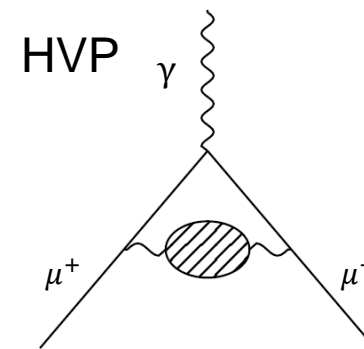
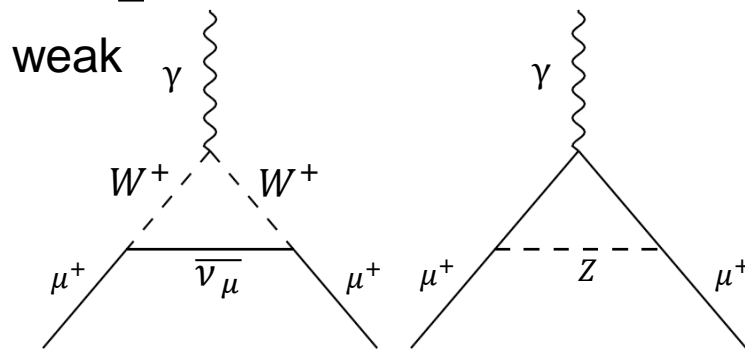
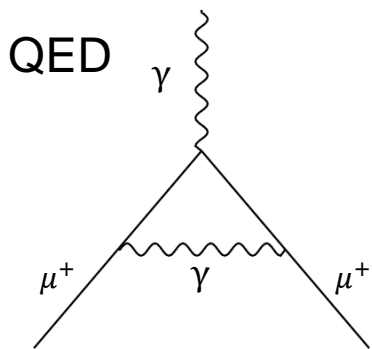
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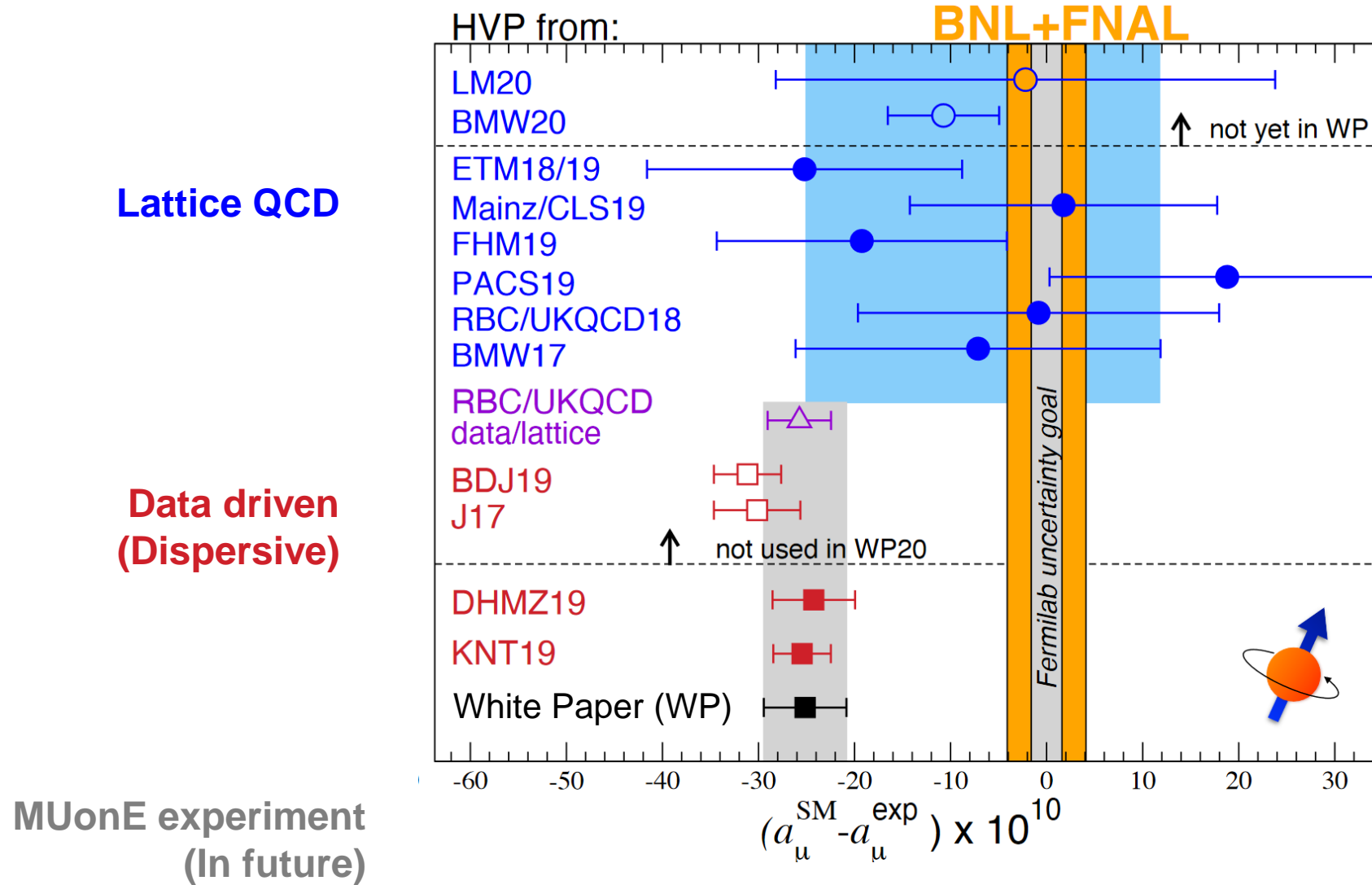


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Hadron vacuum polarization (HVP) contribution

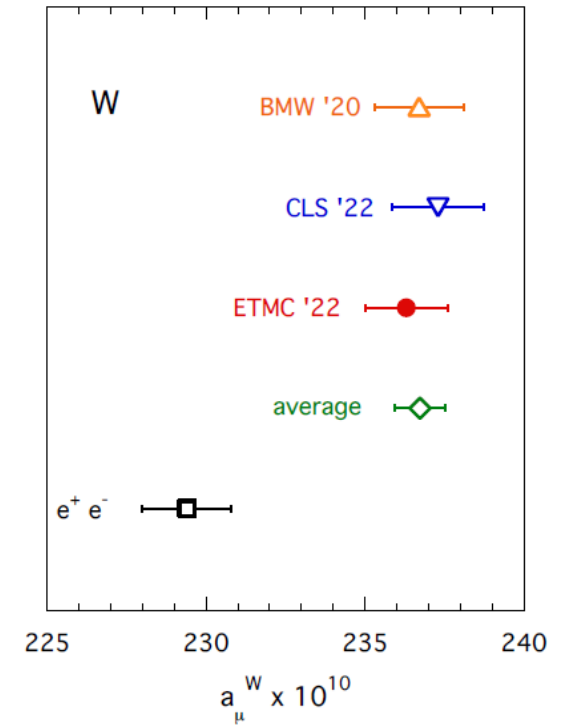
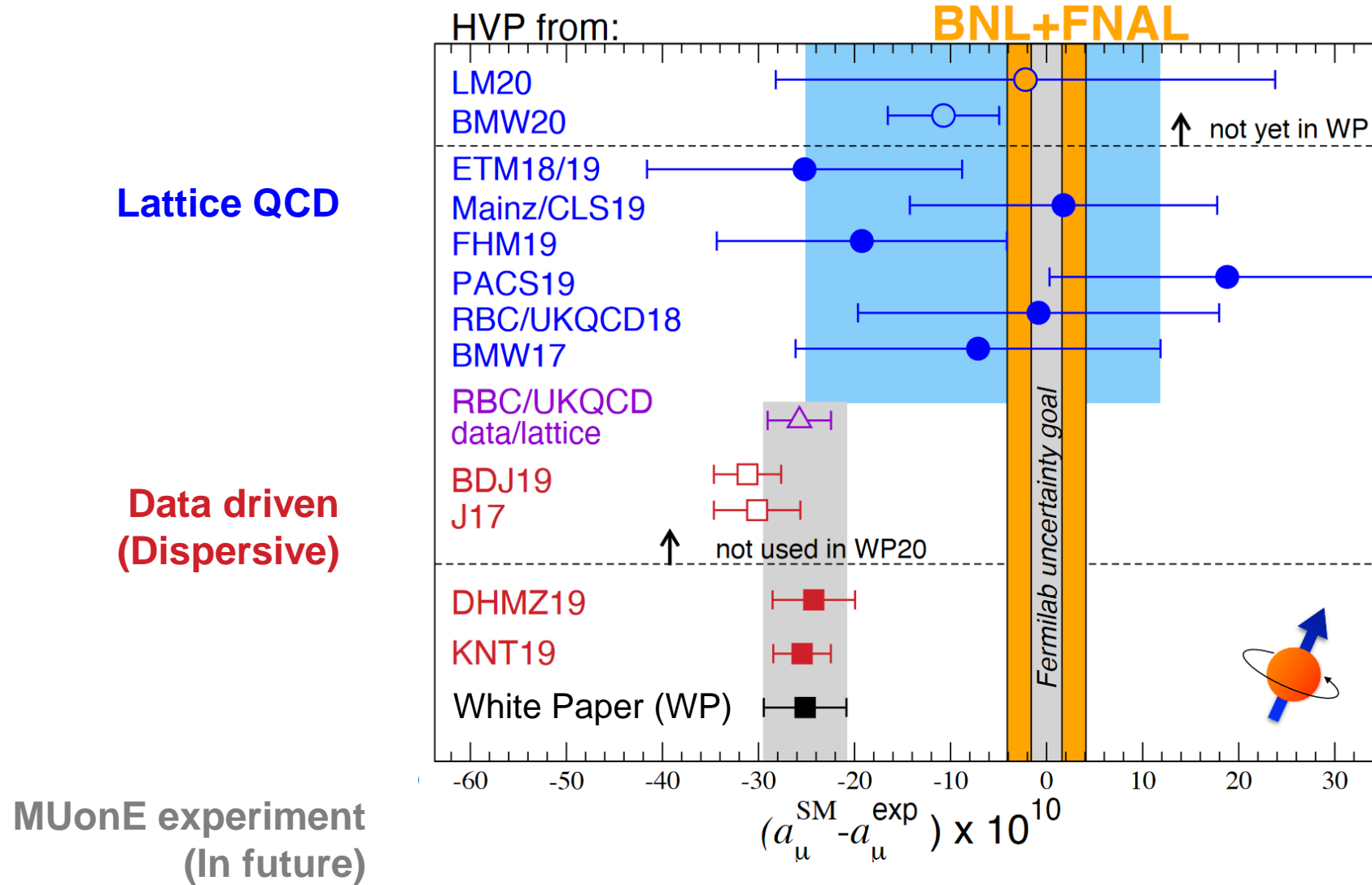
G. Colangelo *et al.*, arXiv:2203.15810 [hep-ex] (2022)



Hadron vacuum polarization (HVP) contribution

G. Colangelo *et al.*, arXiv:2203.15810 [hep-ex] (2022)

C. Alexandrou *et al.*,
arXiv:2206.15084 [hep-lat] (2022)

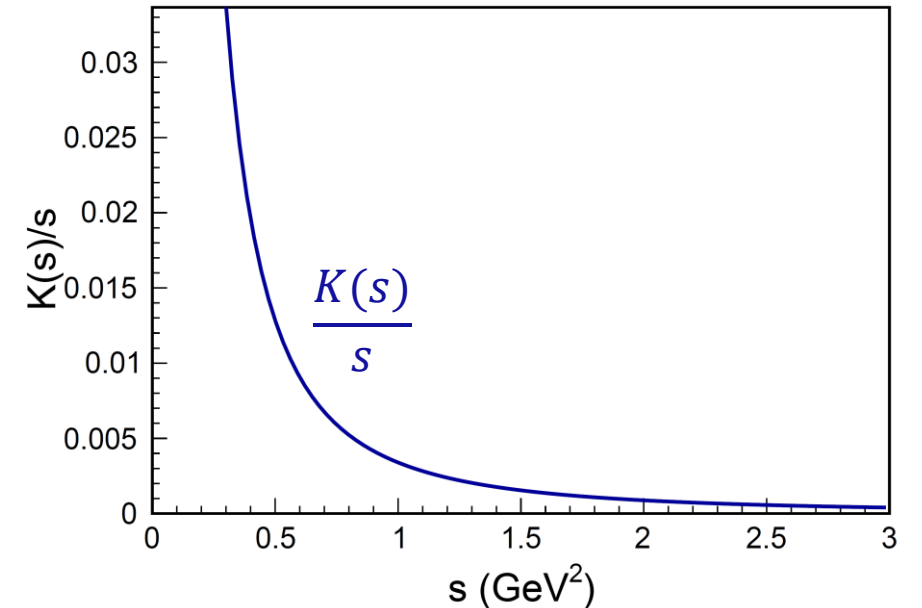


Data-driven method and R measurement

Leading order HVP contribution using dispersion relation

$$a_{\mu}^{\text{HVP,LO}} = \frac{\alpha^2}{3\pi^2} \int_{m_{\pi}^2}^{\infty} \frac{K(s)}{s} R(s) ds$$

$$R(s) = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$$



Kernel function

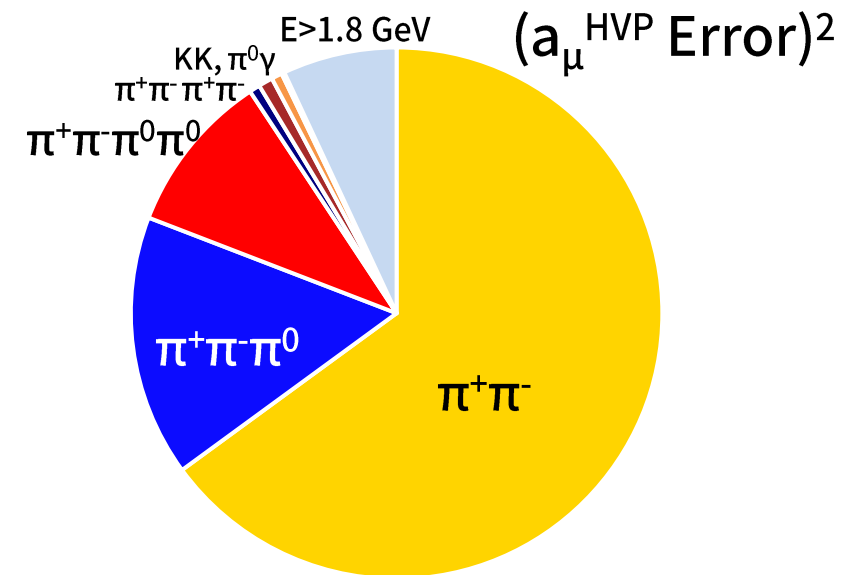
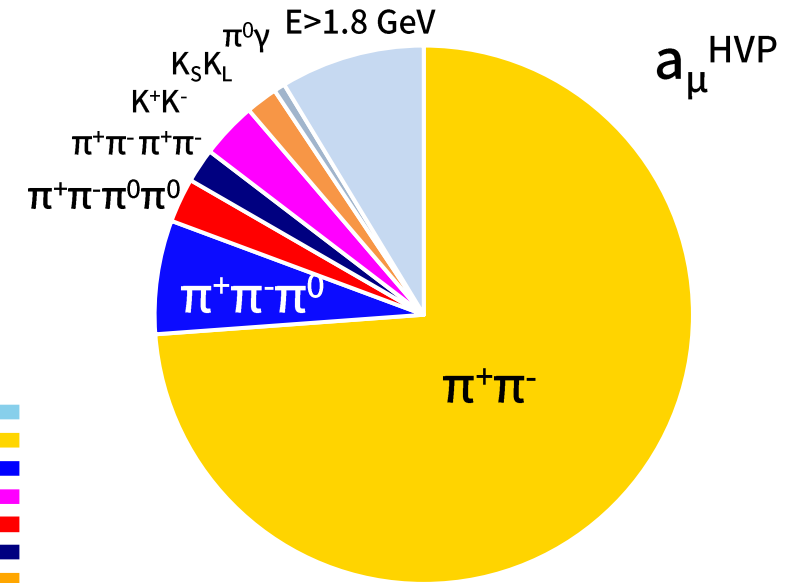
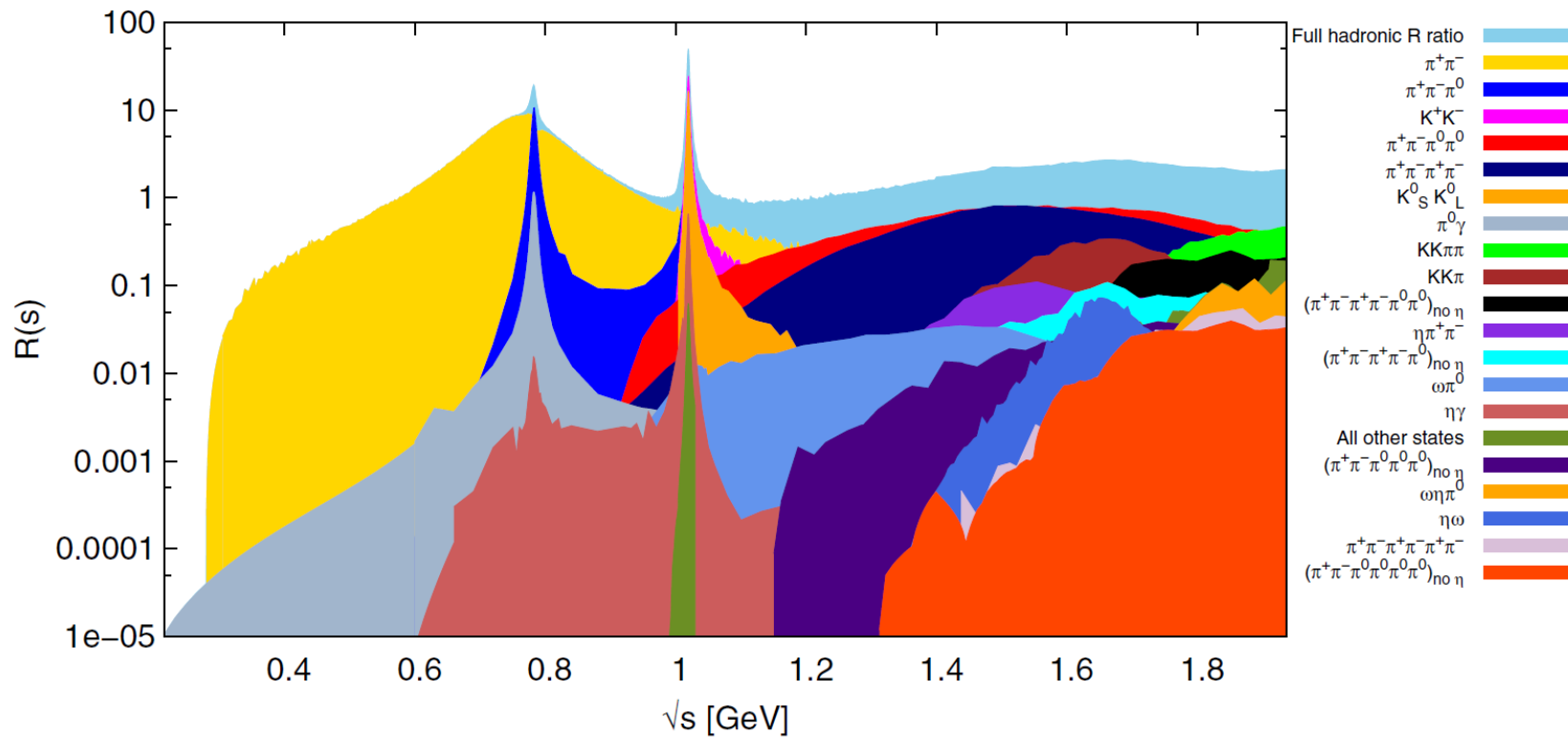
$$K(s) = \frac{x^2}{2} (2 - x) + \frac{(1 + x^2)(1 + x)^2}{x^2} \left(\ln(1 + x) - x + \frac{x^2}{2} \right) + \frac{1 + x}{1 - x} x^2 \ln(x)$$

$$x = \frac{1 - \beta_{\mu}}{1 + \beta_{\mu}}, \quad \beta_{\mu} = \sqrt{1 - \frac{4m_{\mu}^2}{s}}$$

Data-driven method and R measurement

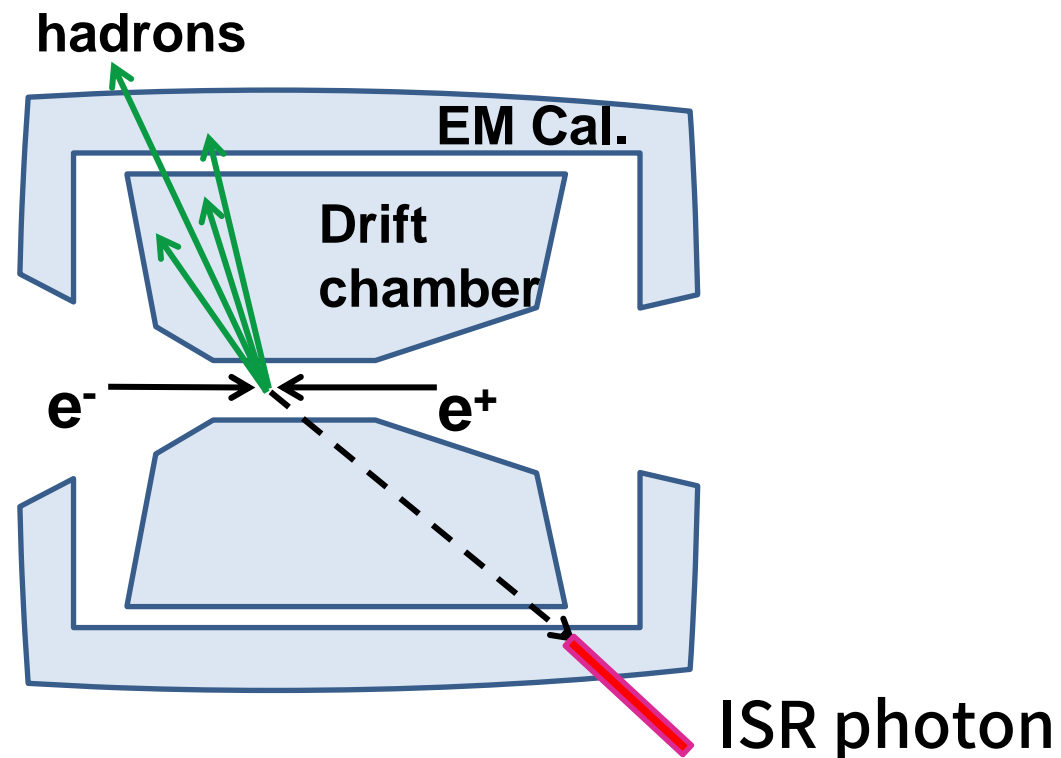
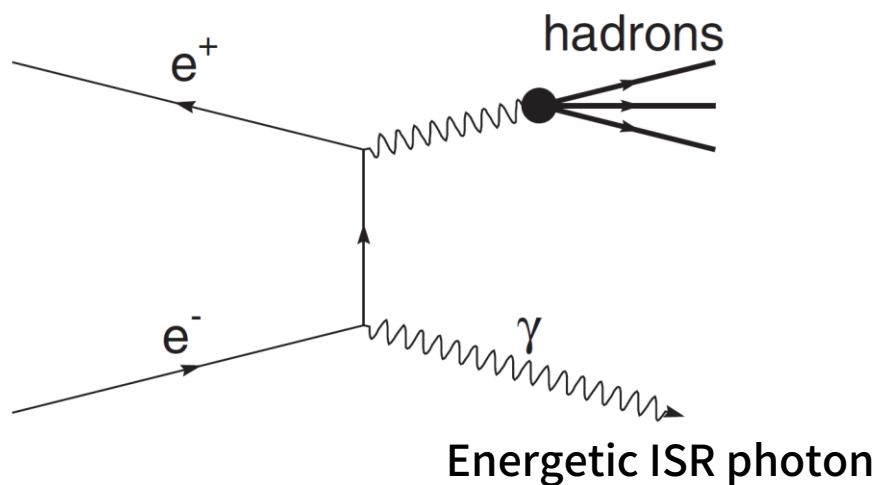
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Radiative return method

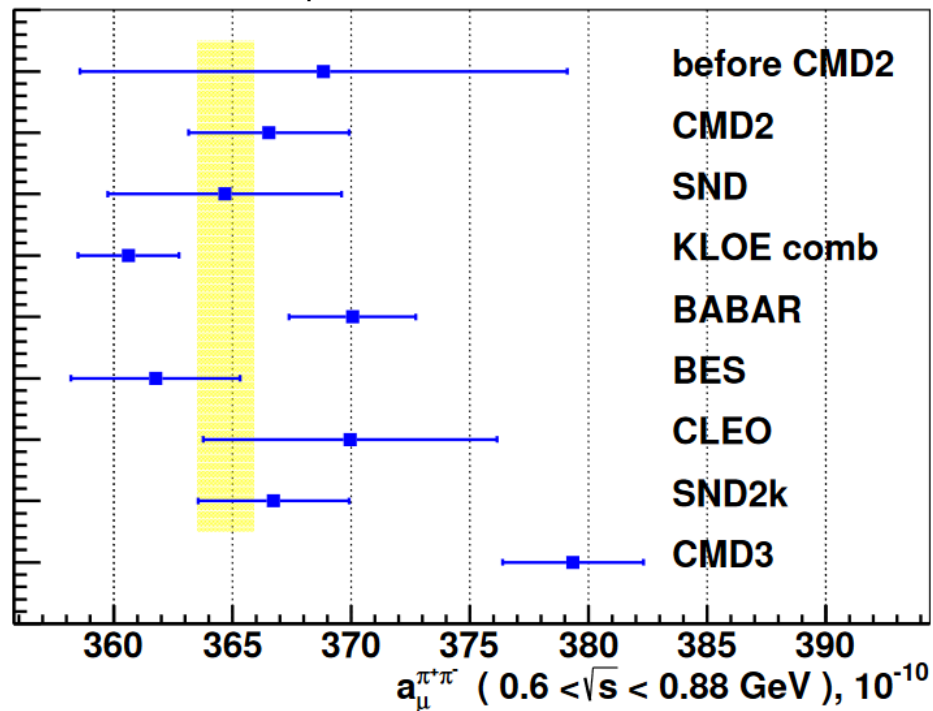
- Radiative return : BaBar, KLOE, BESIII (\leftrightarrow Direct scan : e.g. Novosibirsk)
 - Scan the energy of hadronic system at fixed energy using ISR.
 - Access to the entire hadronic mass range with single dataset
 - Boosted final hadrons



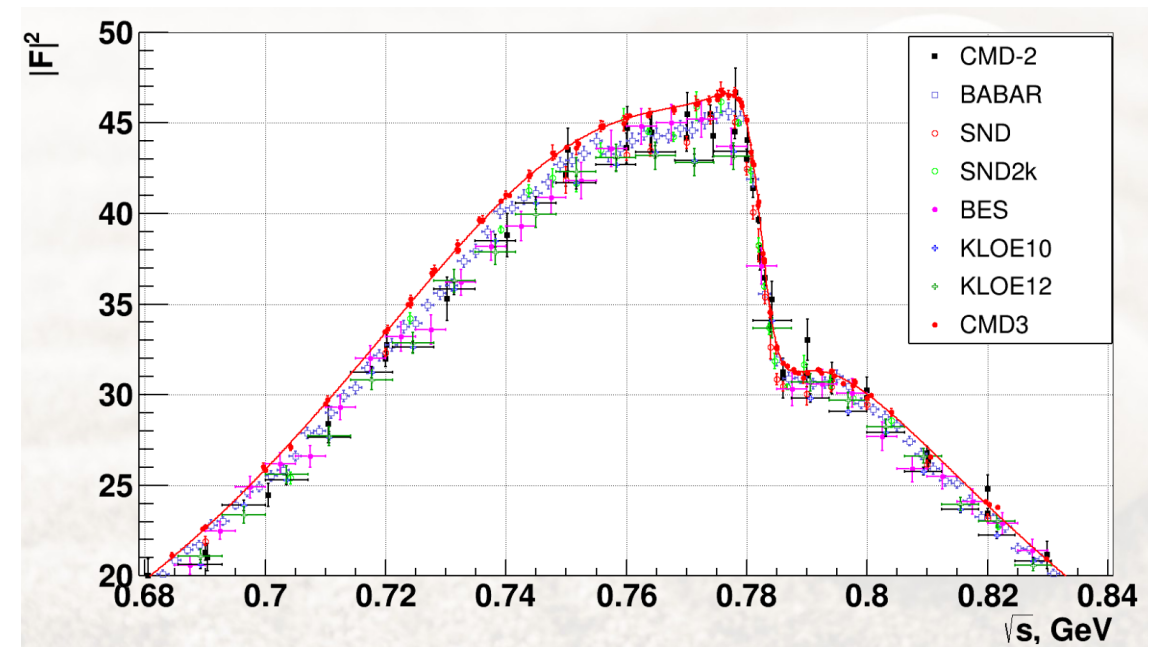
Present experimental status : $e^+e^- \rightarrow \pi^+\pi^-$

- $e^+e^- \rightarrow \pi^+\pi^-$ channel is the largest contribution and uncertainty.
 - Already measured by several experiments with $\approx 1\%$ precision.
 - Small discrepancy among measurements.
- Follow-up tests in Belle II would be important.

a_μ at ρ resonance

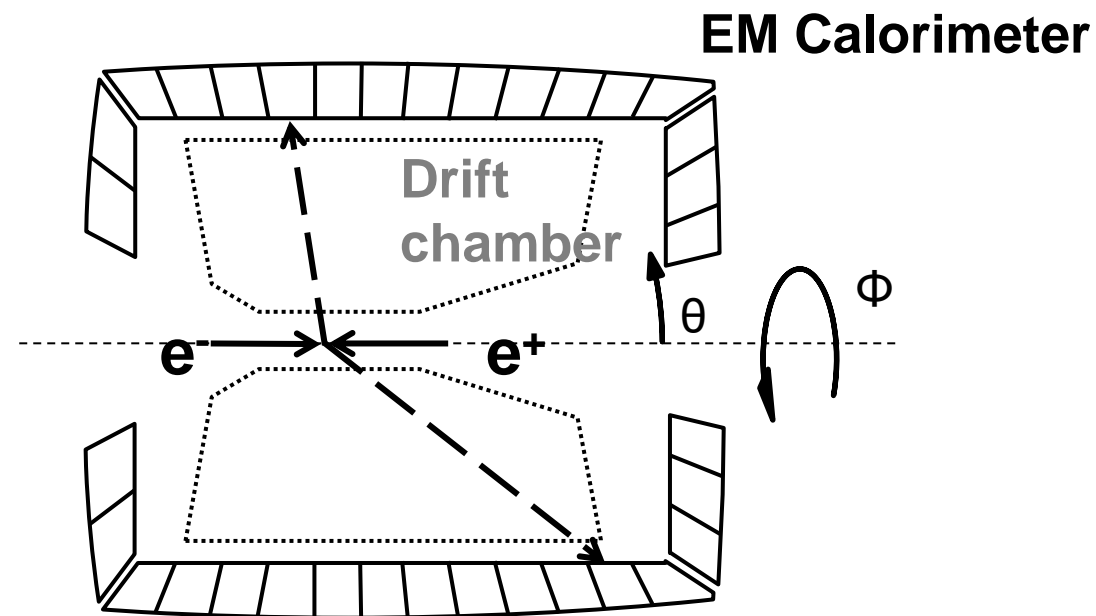


Form factor at ρ resonance



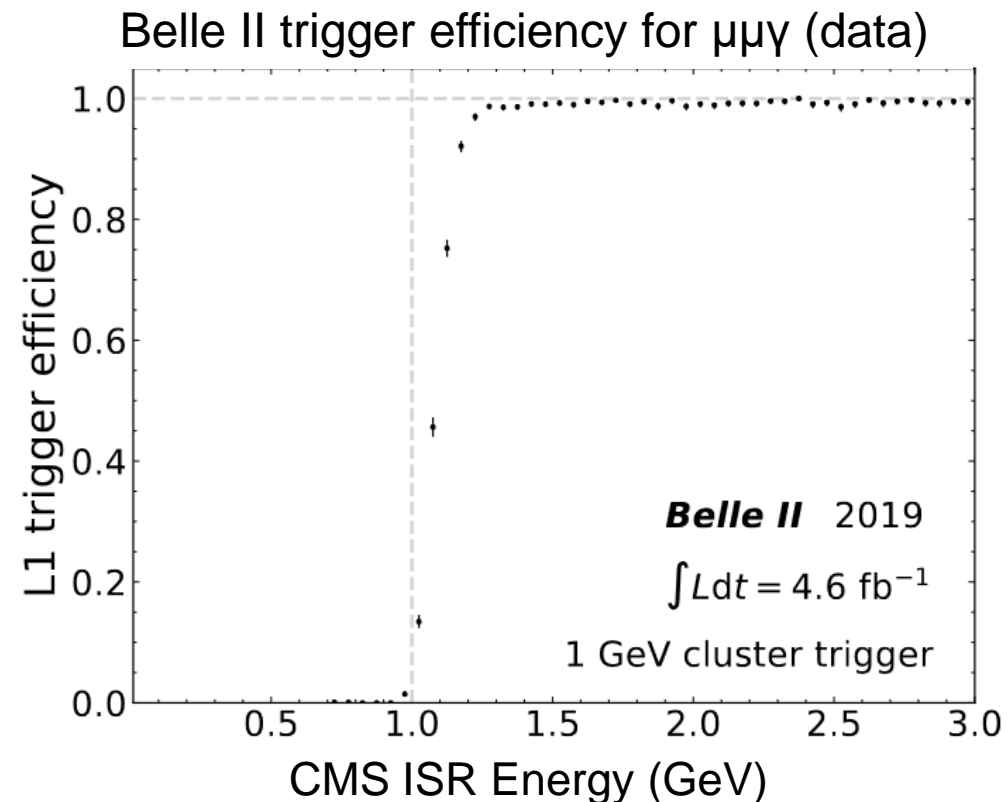
Trigger challenge at Belle II

- Light hadron cross section measurement at BELLE was suffered from the trigger efficiency.
 - The measurement for $\sigma(e^+e^- \rightarrow \pi^+\pi^-\pi^0)$ was attempted, but could not be published.
[J. Crnkovic, PhD thesis, Illinois U. (2013)]
- Bhabha veto has been upgraded to avoid the inefficiency and uncertainty.
 - BELLE bhabha veto was based on only θ angle.
 - Belle II 3D bhabha veto uses θ and Φ angle.
- The trigger efficiency of EM Calorimeter triggers for energetic ISR can be measured by making the orthogonal tracking trigger a reference.



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 - Belle II 3D bhabha veto uses θ and Φ angle.
- The trigger efficiency of EM Calorimeter triggers for energetic ISR can be measured by making the orthogonal tracking trigger a reference.
 - Efficiency for energetic ISR > 99%
 - Event loss due to 3D bhabha veto is suppressed in $\mu\mu\gamma$.
- The high trigger efficiency for energetic ISR is beneficial for most light hadron cross section measurements in the radiative return method.



$e^+e^- \rightarrow \pi^+\pi^-$: Status at Belle II

- Target precision : 0.5% of $a_\mu(2\pi)$
- Trying to follow BaBar methods as a base line.
- Systematics uncertainty dominant analysis
 - BaBar : 232 /fb [Phys. Rev. D 86 (2012), 032013]
 - We can use large statistics to control systematic uncertainties.
- Implementation of kinematic fitting tools
 - Useful for reducing background and correction for tracking efficiency.
 - Implementation of basic fitter has been completed.

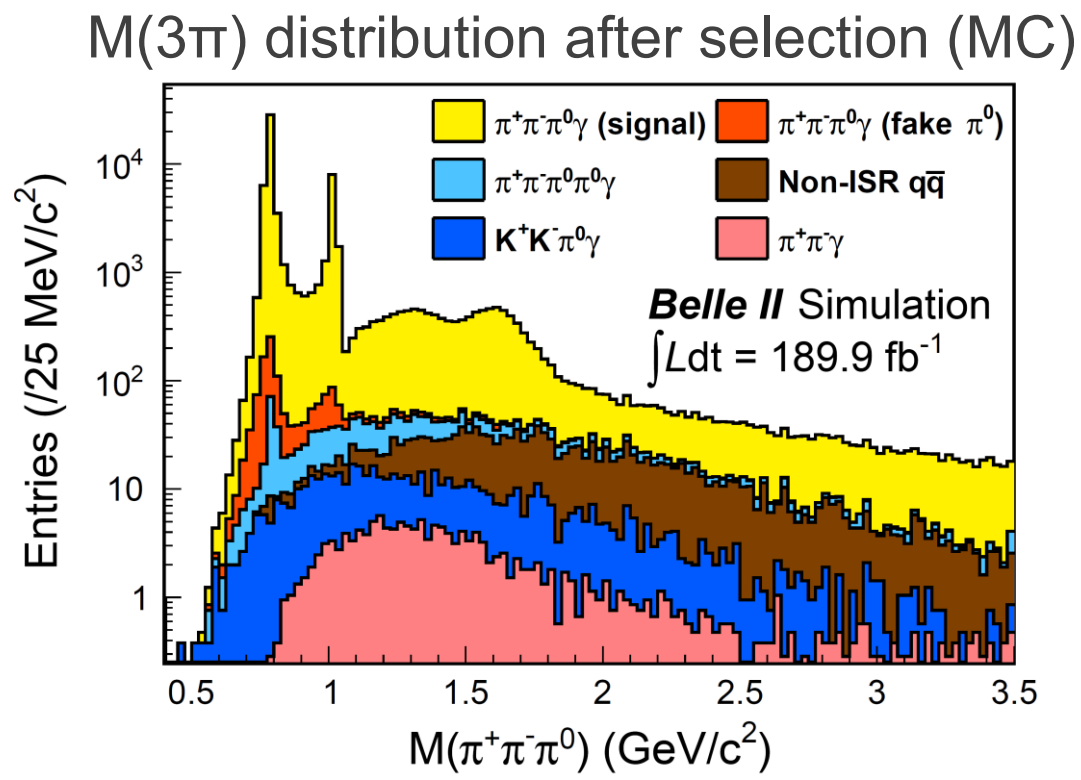
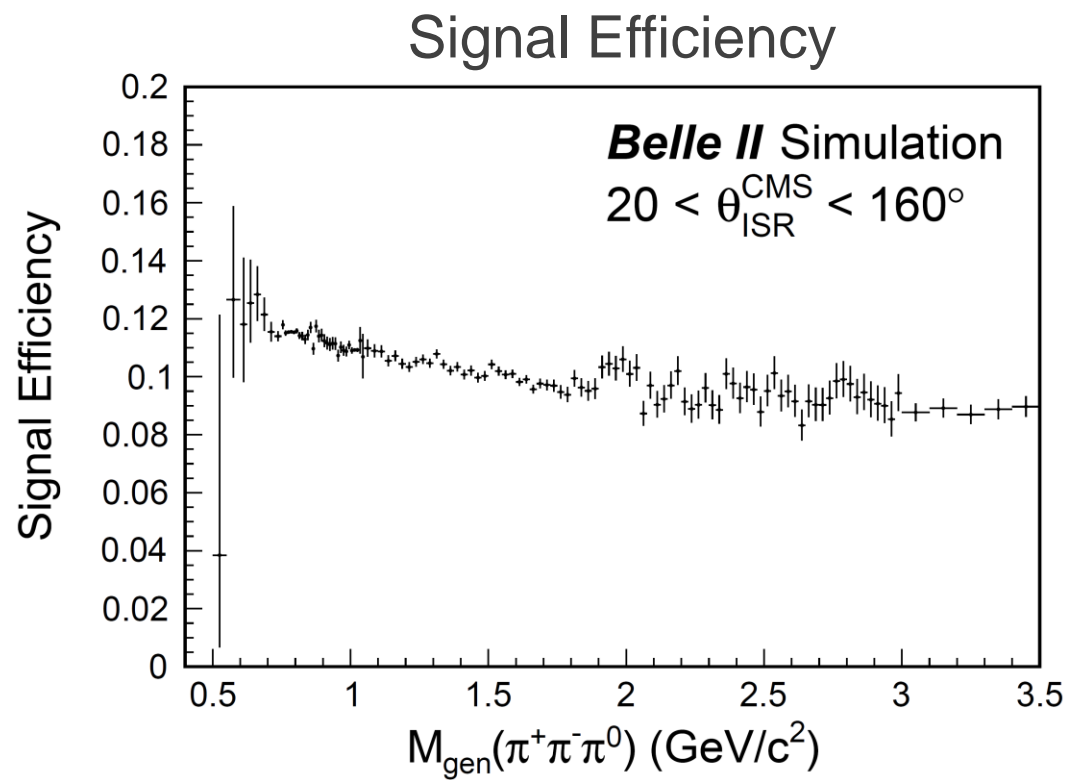
- Sanity check on signal generator and background MC using $< 2 \text{ fb}^{-1}$ data .
- Design of data-driven efficiency corrections for tracking, trigger and $\pi/\mu/K$ ID is ongoing.

$ee \rightarrow \pi\pi$ uncertainty (10^{-3}) at BaBar

Sources	0.3–0.4	0.4–0.5	0.5–0.6	0.6–0.9	0.9–1.2
Trigger/filter	5.3	2.7	1.9	1.0	0.7
Tracking	3.8	2.1	2.1	1.1	1.7
π -ID	10.1	2.5	6.2	2.4	4.2
Background	3.5	4.3	5.2	1.0	3.0
Acceptance	1.6	1.6	1.0	1.0	1.6
Kinematic fit (χ^2)	0.9	0.9	0.3	0.3	0.9
Correl. $\mu\mu$ ID loss	3.0	2.0	3.0	1.3	2.0
$\pi\pi/\mu\mu$ non-cancel.	2.7	1.4	1.6	1.1	1.3
Unfolding	1.0	2.7	2.7	1.0	1.3
ISR luminosity	3.4	3.4	3.4	3.4	3.4
Sum (cross section)	13.8	8.1	10.2	5.0	6.5

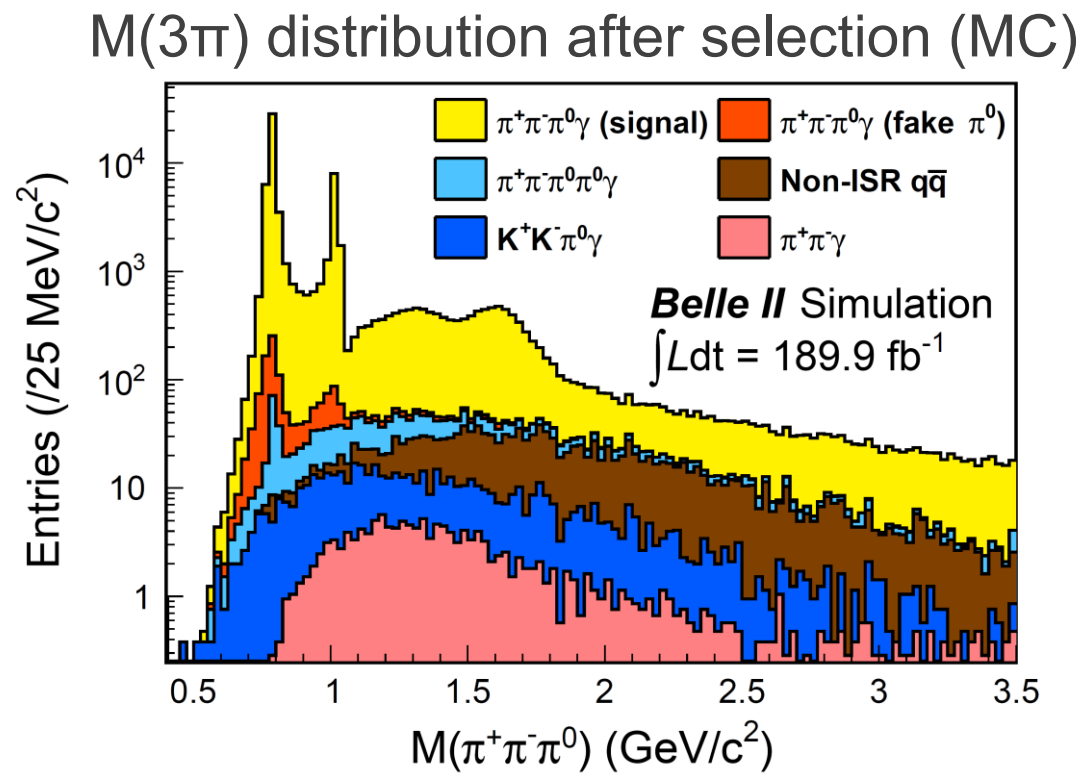
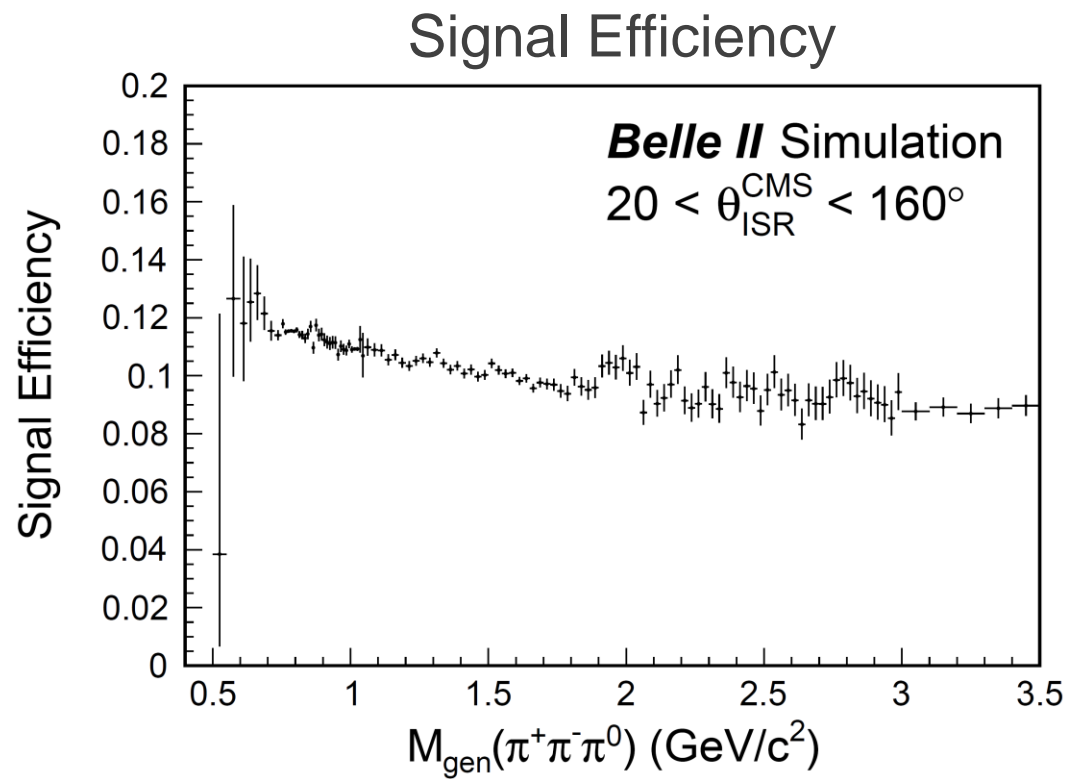
$e^+e^- \rightarrow \pi^+\pi^-\pi^0$ at Belle II

- $e^+e^- \rightarrow \pi^+\pi^-\pi^0$ channel is the 2nd largest contribution to HVP term.
- Aim $\sim 2\%$ precision measurement using 190 fb^{-1} data
- Most analysis methods are fixed and are in final confirmation with 10% data



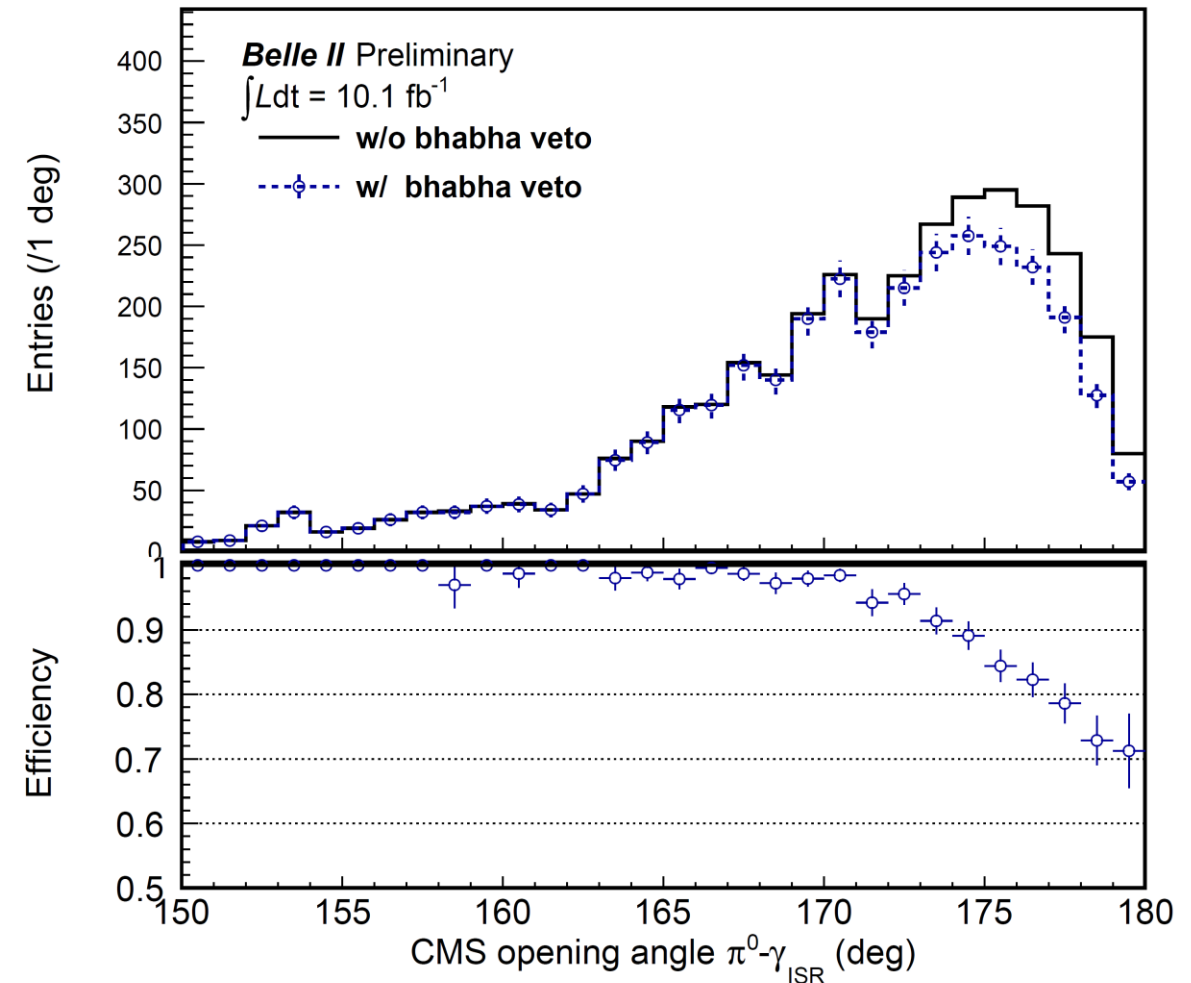
$e^+e^- \rightarrow \pi^+\pi^-\pi^0$ at Belle II

- Selection : two charged tracks + three photon
 - Use kinematic fit χ^2 probability to select events consistent with signal topology
 - Prioritising the reduction of systematic errors.
- Signal efficiency of $\sim 10\%$ is expected.
- Main remaining background : π^0 combinatorial, $\pi^+\pi^-\pi^0\pi^0\gamma$, non-ISR $q\bar{q}$
- MC study using 10 times more than data.



$e^+e^- \rightarrow \pi^+\pi^-\pi^0$: Efficiency correction

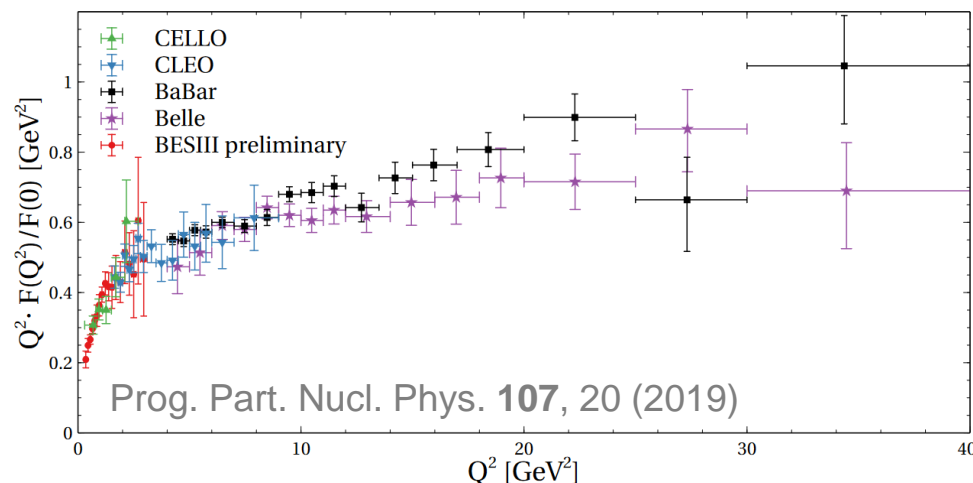
- Photon, tracking and trigger efficiency are confirmed.
- The additional data-driven corrections being evaluated: π^0 efficiency, correlated tracking inefficiency, and background rejection criteria.
- Study for trigger bhabha veto :
 - For $ee \rightarrow \pi^+\pi^-\pi^0$ high energy $\pi^0 \rightarrow \gamma\gamma$ emitted back-to-back to ISR induces 10-15% loss.
 - Almost half of the data affected
 - 100 /fb : without Bhabha veto
 - 90 /fb : with Bhabha veto



Other channels

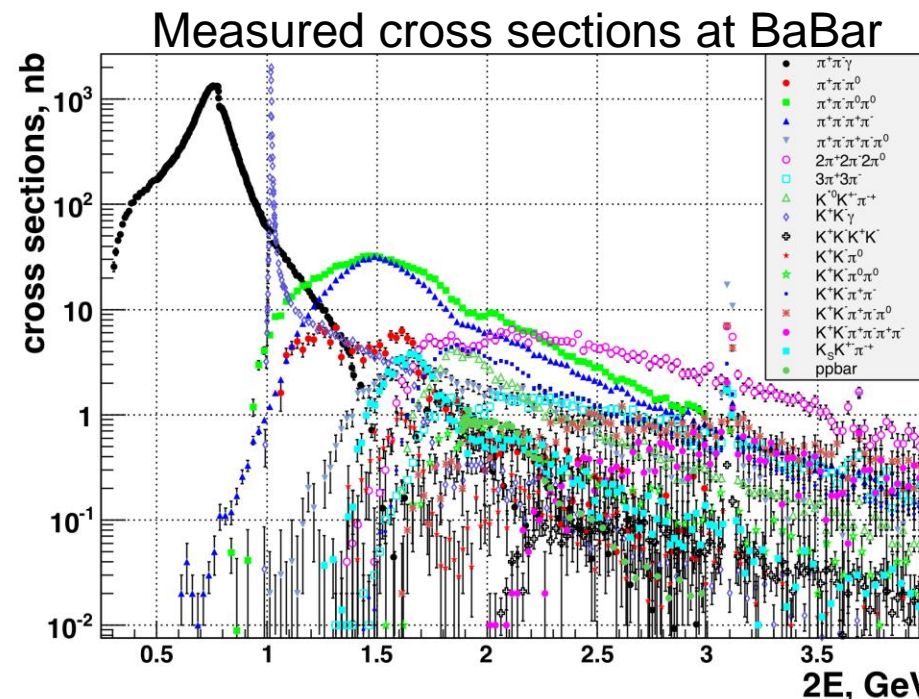
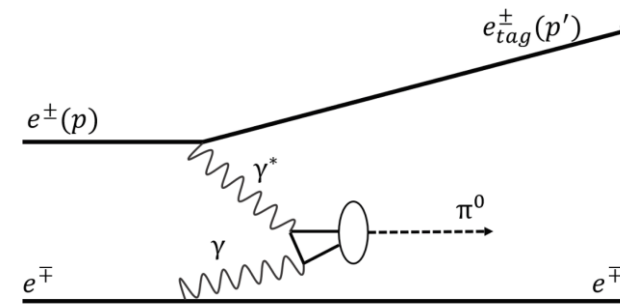
Ongoing channels :

- $\gamma\gamma^* \rightarrow \pi^0$ (Not HVP but Hadron Light-by-Light contribution)
 - Preliminary check using 12 /fb data is done.
 - Further analysis is underway for results using larger dataset.



Further final states can be explored.

- >20 exclusive channels were studied in the BABAR.
 - $\pi^+\pi^-$, $\pi^+\pi^-\pi^0$, $\pi^+\pi^-\pi^0\pi^0$, $\pi^+\pi^-\pi^+\pi^-$, $2\pi^+2\pi^-2\pi^0$, $3\pi^+3\pi^-$, $K^+K^-\pi^0$, $K^0K^\pm\pi^\mp$, K^+K^- , $K^+K^-K^+K^-$, $K^+K^-\pi^0\pi^0$, pp...
- Trigger upgrade allows us to measure other final states.



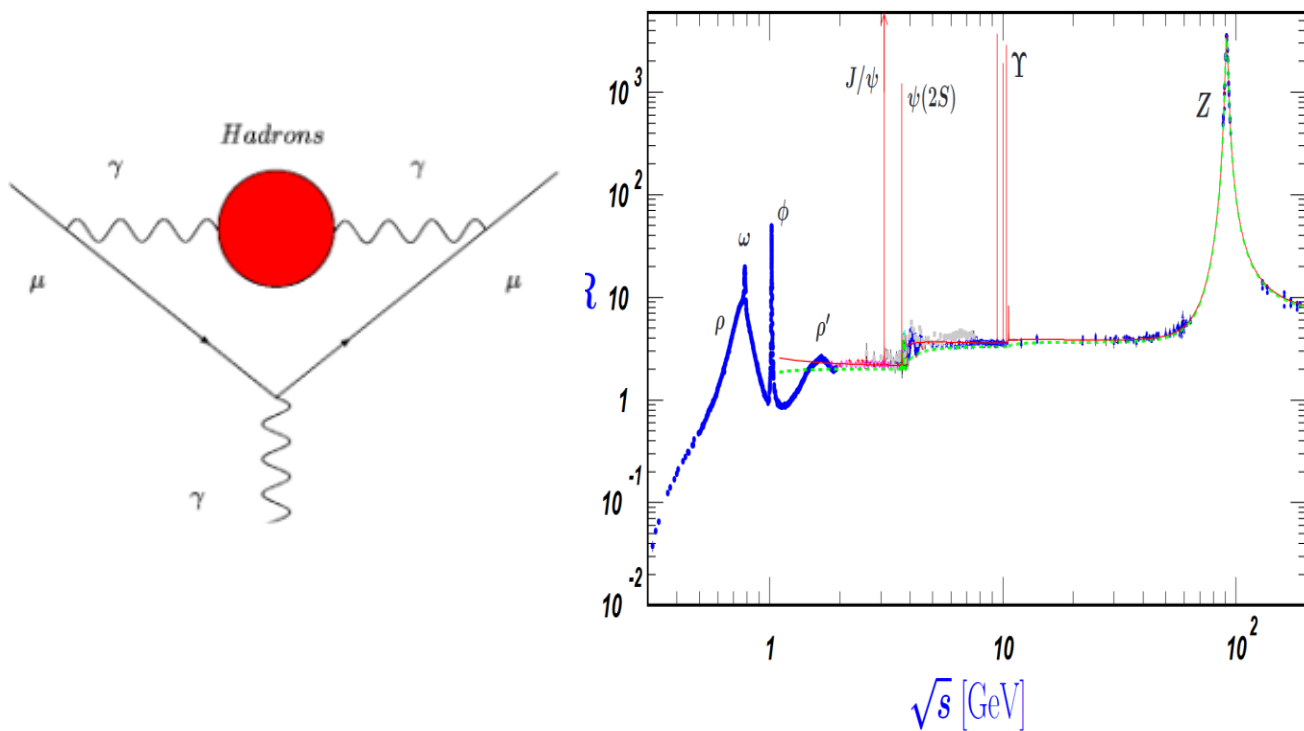
Conclusion

- The light hadron cross section is important in the data-driven method for calculating the HVP contribution of muon $g-2$.
- The trigger upgrade provides us very good efficiency for the cross section measurement.
- Analysis relating to muon $g-2$ are active and in progress.
 - $\pi^+\pi^-$
 - Aim high precision measurement of 0.5%.
 - A methodology based on the BABAR is being established.
 - Focusing on data/MC sanity checks using tiny data of less than 2/fb.
 - $\pi^+\pi^-\pi^0$
 - Aim to release result with $\sim 2\%$ precision using 190/fb data.
 - Most analysis methods are fixed and are in final confirmation with 10% of the data.
 - $\gamma\gamma^*\rightarrow\pi^0$
 - Preliminary check using 12 /fb data is done.
 - Further analysis is underway for results using larger dataset.
- Further channel analysis can be expected in the future.

Backup

Time-like formula

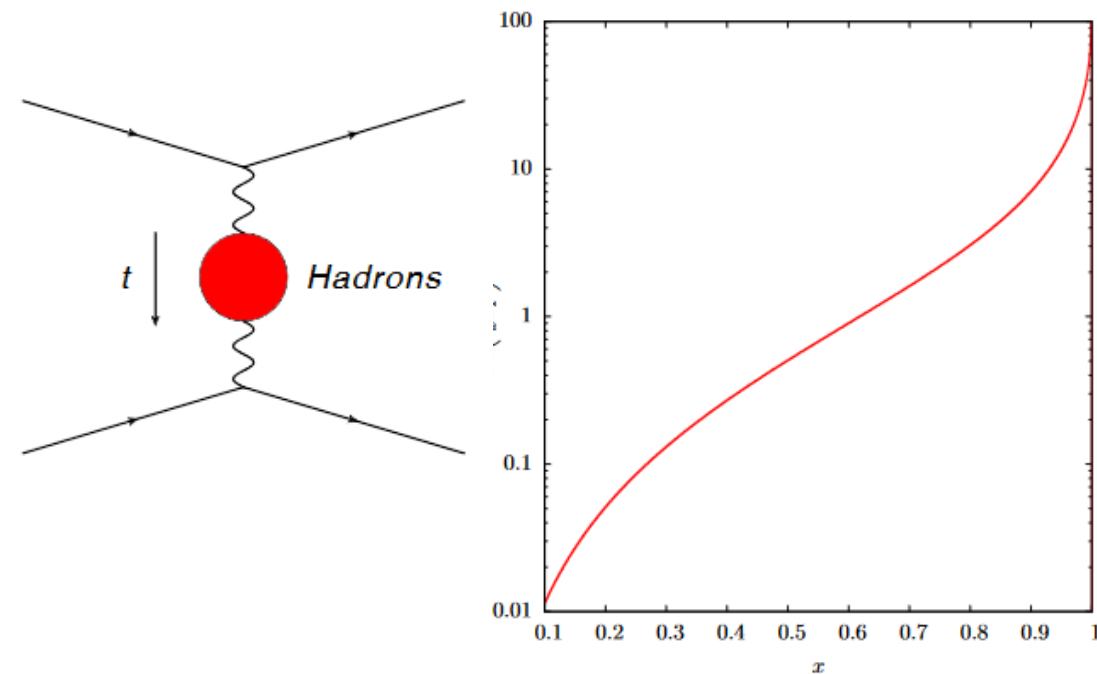
$$a_{\mu}^{\text{HVP,LO}} = \frac{\alpha^2}{3\pi^2} \int_{m_{\pi}^2}^{\infty} \frac{K(s)}{s} R(s) ds$$



Space-like formula

$$a_{\mu}^{\text{HVP,LO}} = \frac{\alpha}{\pi} \int_0^1 dx (1-x) \Delta\alpha_{\text{had}} [t(x)]$$

$$t(x) = \frac{x^2 m_{\mu}^2}{x-1}$$



HVP lattice QCD : Window method

- The window method divides this Euclidean time integral interval by a window function.
- Conversely, it can also be calculated from the R-ratio.

$$a_\mu^{\text{HVP}} = \sum_{t=0}^{\infty} w_t C(t)$$

Lattice QCD+QED correlator $C(t) = \frac{1}{3} \sum_{i, \vec{x}} \langle V_i(\vec{x}, t) V_i(\vec{0}, 0) \rangle$

R-ratio correlator $C(t) = \frac{1}{12\pi^2} \int_0^\infty d(\sqrt{s}) R(s) s e^{-\sqrt{s}t}$

$$a_\mu = a_\mu^{\text{SD}} + a_\mu^{\text{W}} + a_\mu^{\text{LD}}$$

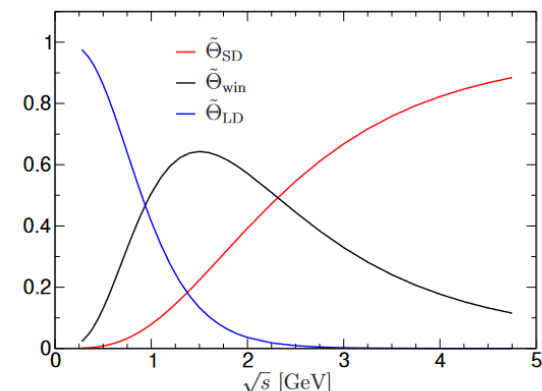
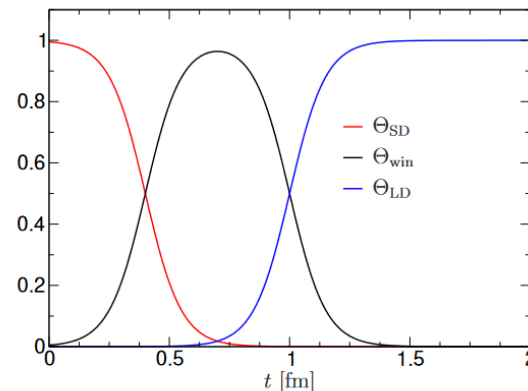
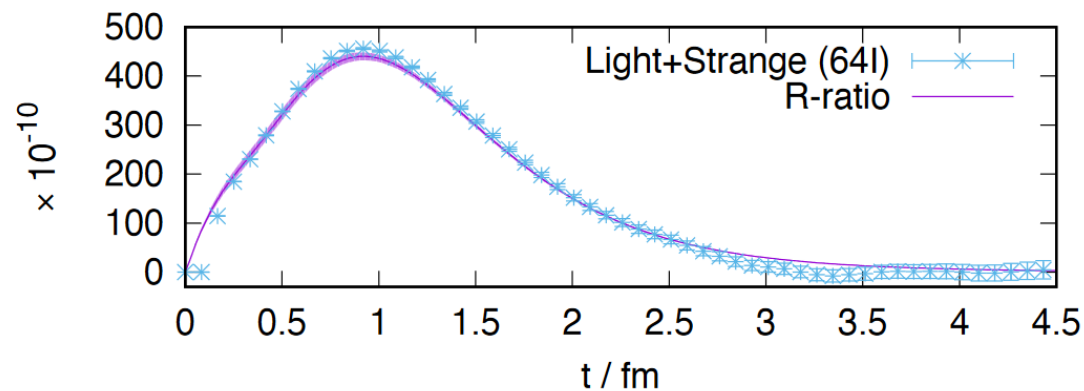
$$a_\mu^{\text{SD}} = \sum_t C(t) w_t [1 - \Theta(t, t_0, \Delta)],$$

$$a_\mu^{\text{W}} = \sum_t C(t) w_t [\Theta(t, t_0, \Delta) - \Theta(t, t_1, \Delta)],$$

$$a_\mu^{\text{LD}} = \sum_t C(t) w_t \Theta(t, t_1, \Delta),$$

$$\Theta(t, t', \Delta) = [1 + \tanh[(t - t')/\Delta]] / 2.$$

$w_t C(t)$



HVP lattice QCD : Window method

- The major systematic error sources differ for each Euclidean time domain.
- The main error sources, short distance (SD) and long distance (LD), can be eliminated.
- Detailed comparisons can be made by extracting only the accurate part of the values: a_μ^{win} .

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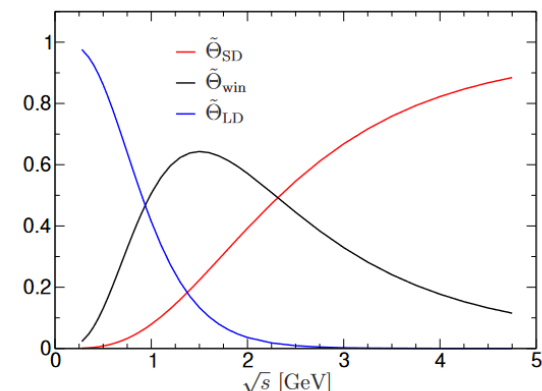
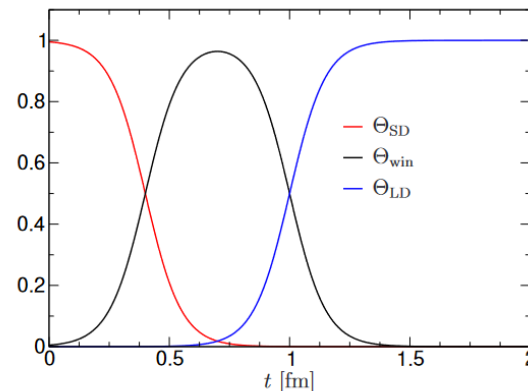
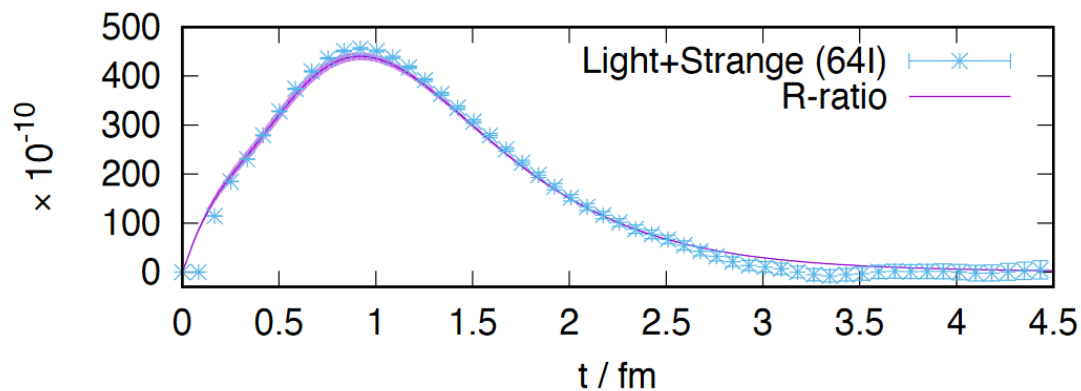
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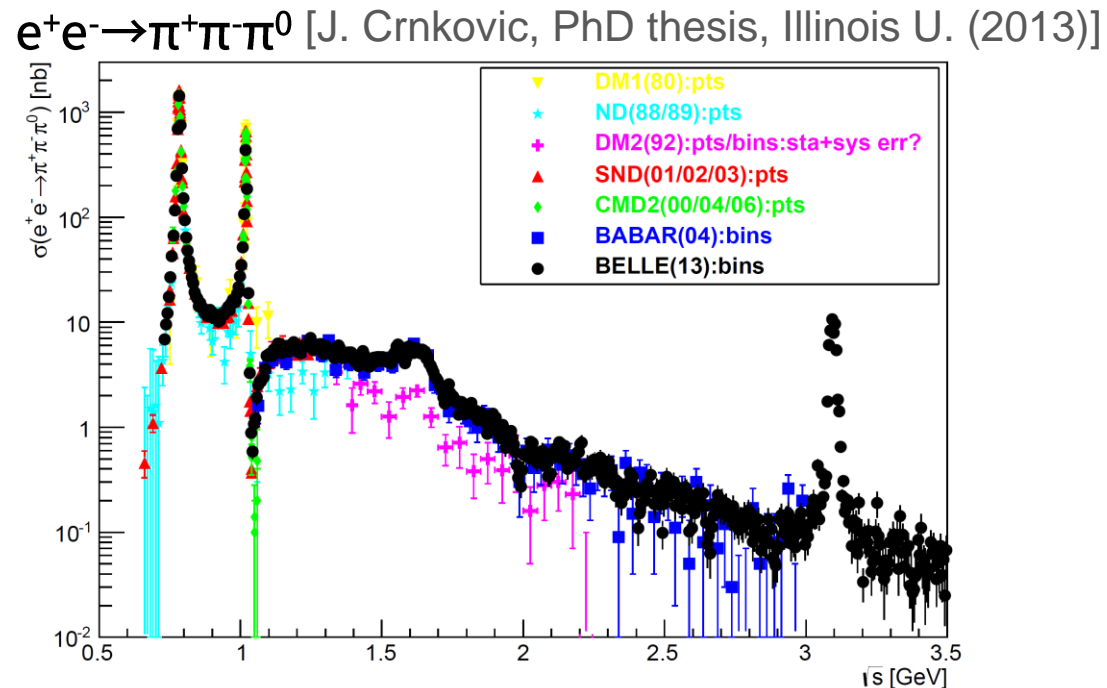
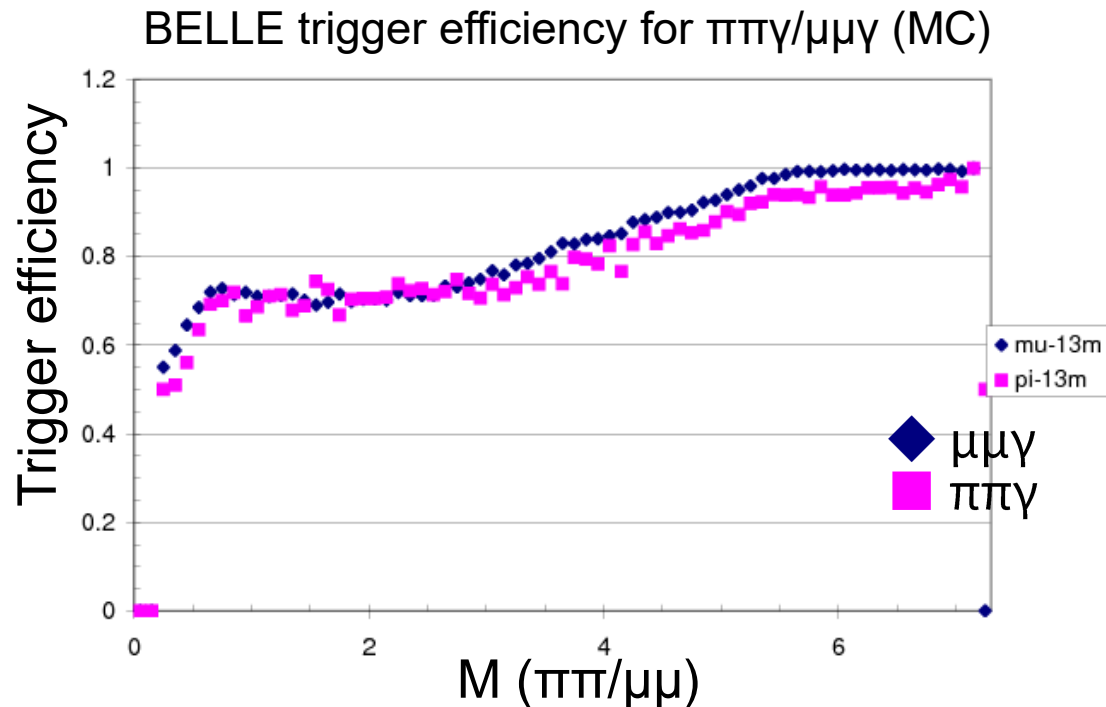
$$\Theta(t, t', \Delta) = [1 + \tanh[(t - t')/\Delta]] / 2.$$

$w_t C(t)$



Light hadron cross section measurement at BELLE

- The measurement for $\sigma(e^+e^- \rightarrow \pi^+\pi^-\pi^0)$ was attempted, but could not be published.
 - Systematic uncertainty goal was 5%.
 - Large uncertainty of level-1 trigger efficiency prevents publication, and the preliminary result is recorded in a PhD thesis.
- A study for $ee \rightarrow \pi\pi$ was also conducted but the triggering efficiency was critical for systematic.
- Other $ee \rightarrow$ hadron cross section has also not been measured at BELLE.



Belle II Detector

Particle Identification

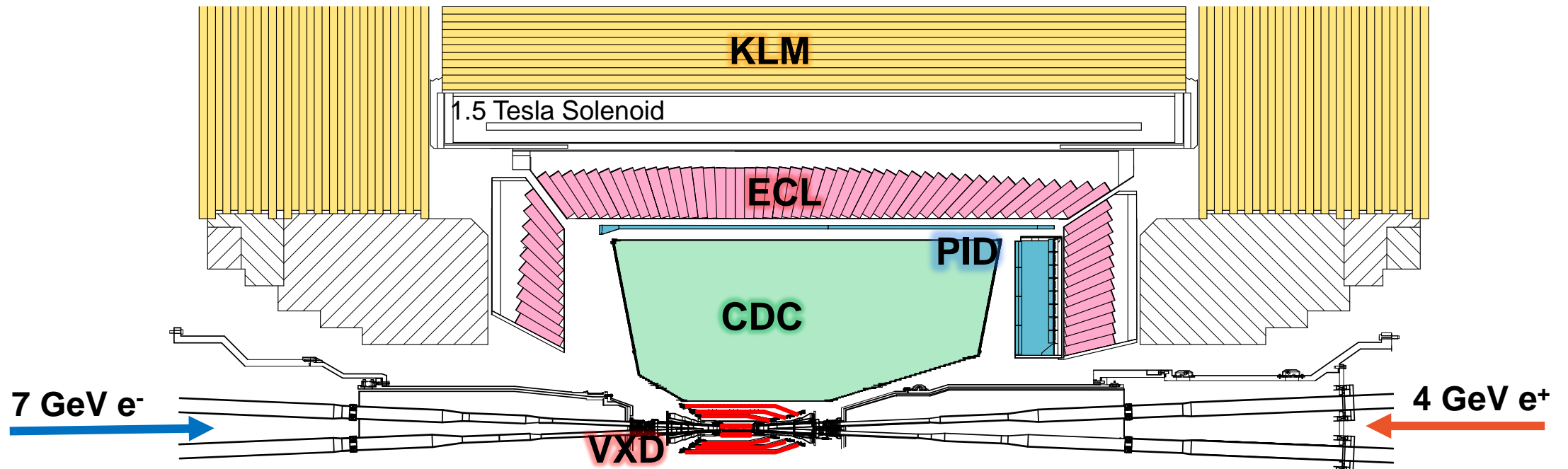
- Aerogel RICH in the forward endcap
- Time-of-Propagation counter in the barrel
- **K/ π ID : K efficiency 90% at 1.8% π fake**

Electromagnetic Calorimeter (ECL)

- CsI(Tl) crystals + Waveform sampling
- **Electron ID eff. 90% at <0.1% fake**
- **Energy resolution 1.6-4%**
- **94% of solid angle coverage**

K-long and Muon Detector (KLM)

- Alternating iron and detector plates
- Scintillator / Resistive Plate Chamber
- **Muon ID efficiency 90% at 2% fake**



Vertex Detector (VXD)

- Inner 2 layer : Pixel
- Outer 4 layer : Double side strip
- vertex resolution 20-30 μm

Central Drift Chamber (CDC)

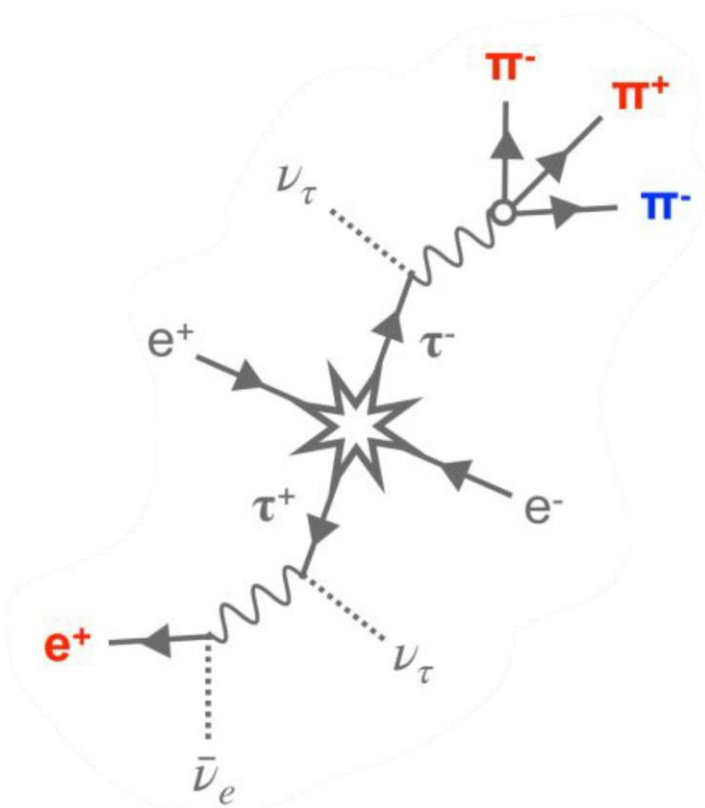
- **91% of solid angle coverage**
- **p_T resolution $\sim 0.4\%/p_T$**
- **dE/dx resolution 5% (low-p PID)**

Trigger and DAQ

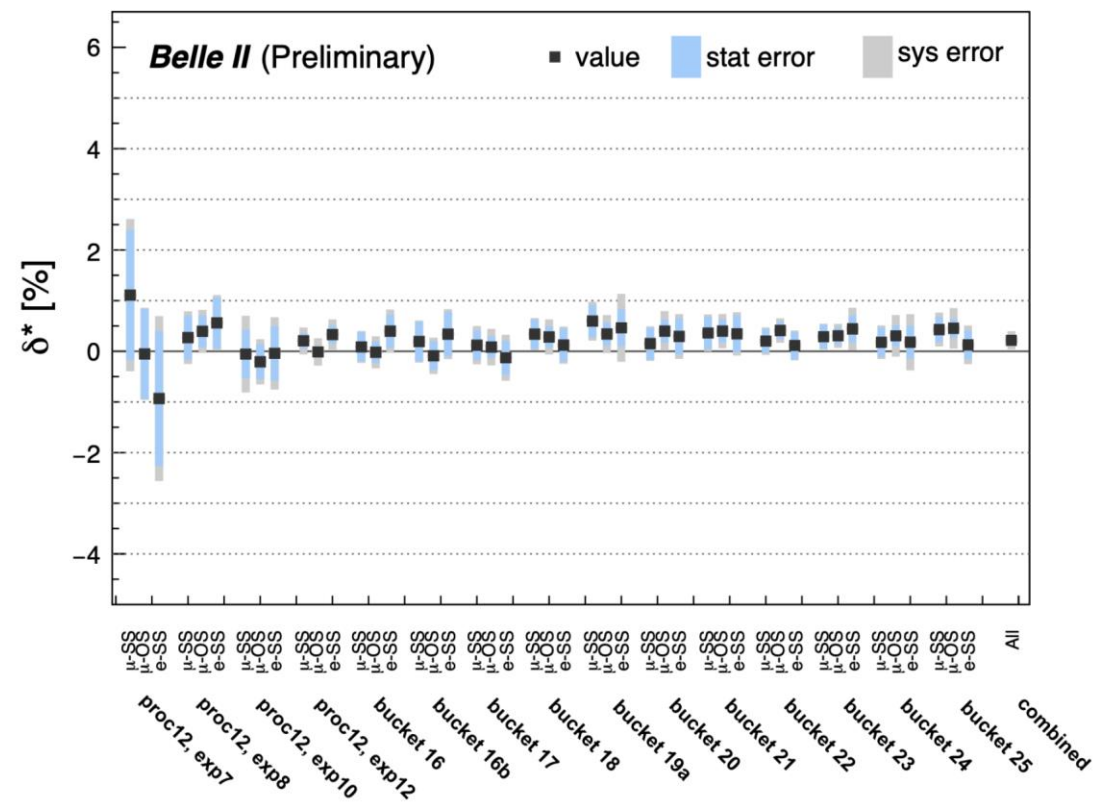
- L1 Trigger rate 30 kHz (design)
- **New trigger line for low-multiplicity events**
- Constant improvements of trigger algorithm

Performance : Tracking Efficiency

- Tracking efficiency is measured by tag-and-probe method on $ee \rightarrow \tau\tau \rightarrow 1 \times 3$ prong.
 - 3 good quality tracks for **tag**
 - Look for 4th track for **probe**
- Uncertainty for tracking efficiency is 0.30% per track.

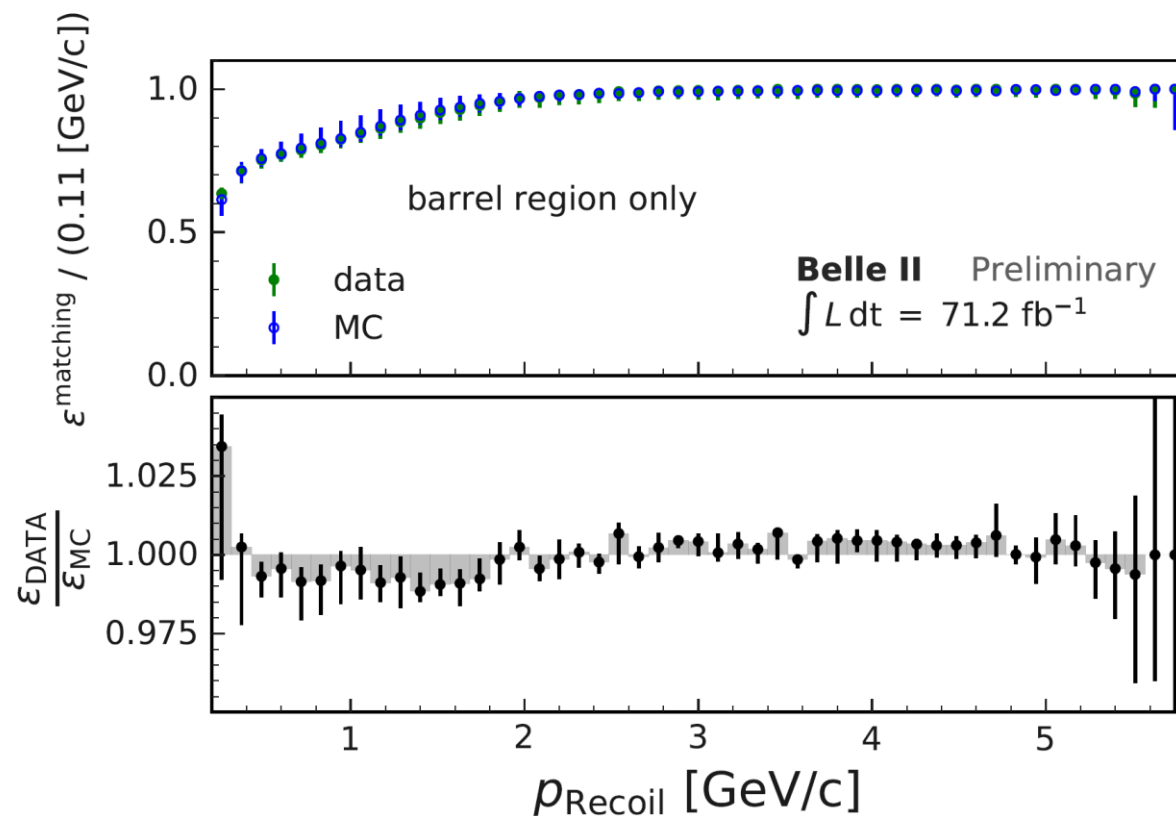
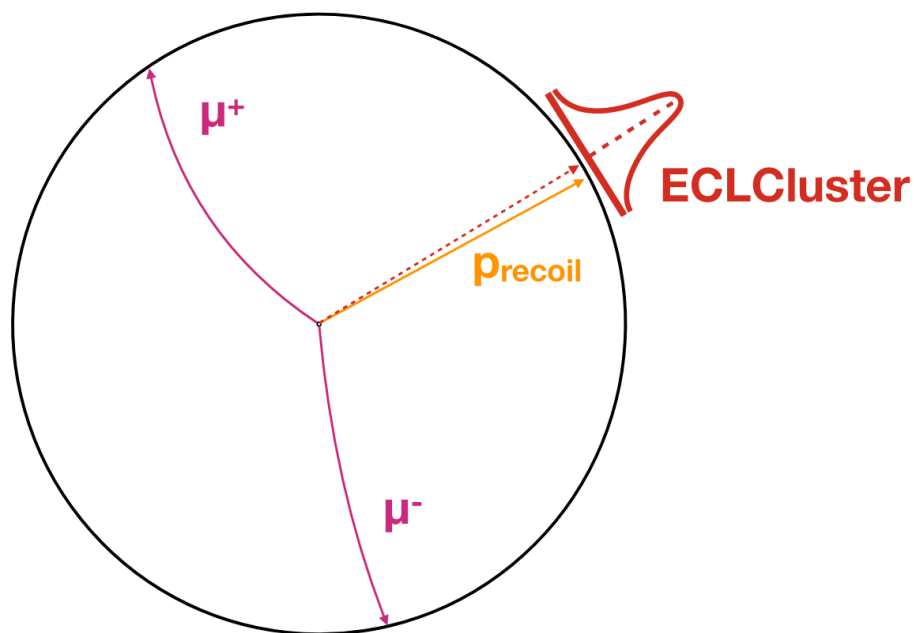


Data/MC discrepancy of tracking efficiency



Performance : Photon Detection Efficiency

- Photon detection efficiency is measured using $ee \rightarrow \mu\mu\gamma$ events.
 - Detection efficiency is estimated by taking match between a ECL cluster and the missing momentum of dimuon system.
- Data/MC agreement is good. Uncertainty for photon detection efficiency is 0.30%.



$e^+e^- \rightarrow \pi^+\pi^-$: Analysis strategy

- Loose selection :two tracks + one hard photon
- Kinematic fitting for Signal/Background separation
 - One fit assuming FSR and one assuming ISR
- PID selaration $\pi\pi/\mu\mu/KK$
- Efficiency corrections

$$\varepsilon^{\text{data}} = \varepsilon^{\text{MC}} \left(\frac{\varepsilon^{\text{data}}}{\varepsilon^{\text{MC}}} \right)_{\text{trigger}} \left(\frac{\varepsilon^{\text{data}}}{\varepsilon^{\text{MC}}} \right)_{\text{PID}} \left(\frac{\varepsilon^{\text{data}}}{\varepsilon^{\text{MC}}} \right)_{\text{tracking}} \left(\frac{\varepsilon^{\text{data}}}{\varepsilon^{\text{MC}}} \right)_{\text{KFit } \chi^2}$$

