# Precise calculation of muon g-2 based on lattice QCD 

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Part of related calculation are done by resources from USQCD (DOE), XSEDE, ANL BG/Q Mira (DOE, ALCC), Edinburgh BG/Q, BNL BG/Q, RIKEN BG/Q and Cluster (RICC, HOKUSAI)

## Reference

- g-2 HVP Phys. Rev. Lett. 121 (2018) 022003
- g-2 Hadronic Light-by-Light (HLbL) Phys. Rev. D96 (2017) 034515 Phys. Rev. Lett. 118 (2017) 022005
- Tau input for g -2 PoS Lattice 2018 (2018) 135


## muon anomalous magnetic moment



BNL g-2 till 2004: ~ $3.7 \sigma$ larger than SM prediction

| Contribution | Value $\times 10^{10}$ | Uncertainty $\times 10^{10}$ |
| :--- | ---: | ---: |
| QED (5 loops) | 11658471.895 | 0.008 |
| EW | 15.4 | 0.1 |
| HVP LO | 692.3 | 4.2 |
| HVP NLO | -9.84 | 0.06 |
| HVP NNLO | 1.24 | 0.01 |
| Hadronic light-by-light | 10.5 | $\mathbf{2 . 6}$ |
| Total SM prediction | 11659181.5 | 4.9 |
| BNL E821 result | 11659209.1 | 6.3 |
| FNAL E989/J-PARC E34 goal |  | $\approx \mathbf{1 . 6}$ |

$$
a_{\mu}^{\mathrm{EXP}}-a_{\mu}^{\mathrm{SM}}=27.4 \underbrace{(2.7)}_{\text {HVP }} \underbrace{(2.6)}_{\text {HLbL }} \underbrace{(0.1)}_{\text {other }} \underbrace{(6.3)}_{\text {EXP }} \times 10^{-10}
$$



FNAL E989 (began 2017-)
move storage ring from BNL
x4 more precise results, 0.14 ppm

J-PARC E34
ultra-cold muon beam
0.37 ppm then 0.1 ppm , also EDM

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[Luchang Jin's analogy]

## Precession of Mercury and GR

| Amount (arc- <br> sec/century) | Cause |
| :---: | :--- |
| 5025.6 | Coordinate (due to precession of equinoxes) |
| 531.4 | Gravitational tugs of the other planets |
| $\mathbf{0 . 0 2 5 4}$ | Oblateness of the sun (quadrupole moment) |
| $\mathbf{4 2 . 9 8} \pm 0.04$ | General relativity |
| 5600.0 | Total |
| 5599.7 | Observed |

discrepancy recognized since 1859

## Known physics

1915 by-then New physics GR revolution
http://worldnpa.org/abstracts/abstracts_6066.pdf precession of perihelion


# Hadronic Vacuum Polarization (HVP) contribution to g-2 




- From experimental e+e-inclusive hadron decay cross section $\sigma_{\text {total }}(\mathrm{s})$ in time-like $\mathrm{s}=\mathrm{q}^{2}>0$, and dispersion relation, optical theorem

$$
a_{\mu}^{\mathrm{HVP}}=\frac{1}{4 \pi^{2}} \int_{\mathrm{s}_{\mathrm{th}}}^{\infty} d s K(s) \sigma_{\mathrm{total}}(s) \quad \sim_{\text {had }}
$$

## Dispersive methods 2018

[ D. Nomura's talk ]

- KNT18 (PRD97,114025, arXiv:1802.02995)
- DHMZ17 (Eur. Phys. J. C77:827)

| Channel | This work (KNT18) | DHMZ17 [78] | Difference |
| :--- | :---: | :---: | :---: |
| Data based channels $(\sqrt{s} \leq 1.8 \mathrm{GeV})$ |  |  | 0.29 |
| $\pi^{0} \gamma($ data +ChPT$)$ | $4.58 \pm 0.10$ | $4.29 \pm 0.10$ | -3.40 |
| $\pi^{+} \pi^{-}($data +ChPT$)$ | $503.74 \pm 1.96$ | $507.14 \pm 2.58$ | 1.50 |
| $\pi^{+} \pi^{-} \pi^{0}($ data +ChPT$)$ | $47.70 \pm 0.89$ | $46.20 \pm 1.45$ | 0.31 |
| $\pi^{+} \pi^{-} \pi^{+} \pi^{-}$ | $13.99 \pm 0.19$ | $13.68 \pm 0.31$ |  |


| Total | $693.3 \pm 2.5$ | $693.1 \pm 3.4$ | 0.2 |
| :--- | :---: | :---: | :---: |

- Very small error, KNT18: $2.5 \times 10^{-10}$ [ $0.37 \%$ ] and DHMZ17 $3.4 \times 10^{-10}$ [ $0.49 \%$ ]
- Good agreement for total, individual channels have a tention.
- Difference in how to combine experiments and energy bins, correlations among them


## Dispersive method status

- BaBar and KLOE $2 \pi$ contribution differ $\sim 10(4) \times 10^{-10}$ compared with quoted uncertainties, $\{2.5$ or 3.4$\} \times 10^{-10}$

[B. Malaescu's talk @Mainz g-2 2018]




## HVP from Lattice

- Analytically continue to Euclidean/space-like momentum $\mathrm{K}^{2}=-\mathrm{q}^{2}>0$
- Vector current 2 pt function
$a_{\mu}=\frac{g-2}{2}=\left(\frac{\alpha}{\pi}\right)^{2} \int_{0}^{\infty} d K^{2} f\left(K^{2}\right) \hat{\Pi}\left(K^{2}\right) \quad \Pi^{\mu \nu}(q)=\int d^{4} x e^{i q x}\left\langle J^{\mu}(x) J^{\nu}(0)\right\rangle$
- Low Q2, or long distance, part of $\Pi$ (Q2) is relevant for g-2





## Euclidean Time Momentum Representation

[Bernecker Meyer 2011 , Feng et al. 2013]
In Euclidean space-time, project verctor 2 pt to zero spacial momentum, $\vec{p}=0$ :

$$
C(t)=\frac{1}{3} \sum_{x, i}\left\langle j_{i}(x) j_{i}(0)\right\rangle
$$

g-2 HVP contribution is

$$
\begin{gathered}
a_{\mu}^{H V P}=\sum_{t} w(t) C(t) \\
w(t)=2 \int_{0}^{\infty} \frac{d \omega}{\omega} f_{\mathrm{QED}}\left(\omega^{2}\right)\left[\frac{\cos \omega t-1}{\omega^{2}}+\frac{t^{2}}{2}\right] \\
\mathrm{w}(\mathrm{t}) \sim \mathrm{t}^{4}
\end{gathered}
$$



- Subtraction $\Pi(0)$ is performed. Noise/Signal $\sim e^{\left(E_{\pi \pi}-m_{\pi}\right) t}$, is improved [Lehner et al. 2015].
- Corresponding $\hat{\Pi}\left(Q^{2}\right)$ has exponentially small volume error [Portelli et al. 2016] . $w(t)$ includes the continuum QED part of the diagram


## DWF light HVP [ 2016 Christoph Lehner ]



120 conf ( $a=0.11 \mathrm{fm}$ ), 80 conf ( $a=0.086 \mathrm{fm}$ ) physical point $\mathrm{Nf}=2+1$ Mobius DWF $4 D$ full volume LMA with 2,000 eigen vector (of e/o preconditioned zMobius $D^{+} D$ ) EV compression ( $1 / 10$ memory) using local coherence [ C. Lehner Lat2017 Poster ] In addition, 50 sloppy / conf via multi-level AMA more than $\times 1,000$ speed up compared to simple CG

## disconnected quark loop contribution

- [ C. Lehner et al. (RBC/UKQCD 2015, arXiv:1512.09054, PRL) 1
- Very challenging calculation due to statistical noise
- Small contribution, vanishes in SU(3) limit, Qu+Qd+Qs = 0
- Use low mode of quark propagator, treat it exactly ( all-to-all propagator with sparse random source )
- First non-zero signal


## Sensitive to $\mathrm{m}_{\pi}$

 crucial to compute at physical mass



## HVP QED+ strong IB corrections

- HVP is computed so far at Iso-symmetric quark mass, needs to compute isospin breaking corrections: Qu, Qd, mu-md $\neq 0$
- u,d,s quark mass and lattice spacing are re-tuned using \{charge, neutral\} $\times\{$ pion,kaon\} and ( Omega baryon masses )
- For now, V, S, F, M are computed : assumes EM and IB of sea quark and also shift to lattice spacing is small (correction to disconnected diagram)
- Point-source method : stochastically sample pair of 2 EM vertices a la important sampling with exact photon

(a) V

(c) T
(b) S

(d) D1

(e) D2

(f) F

(g) D3



## Tau input for HVP IB+QED corrections

- Could also compute the difference IB correction of
$\Delta \mathrm{a}_{\mu}=\mathrm{a}_{\mu}(\mathrm{e}+\mathrm{e}-)-\mathrm{a}_{\mu}(\tau)$

$\pi^{+} \pi^{-}, \cdots[I=1]$
isospin rotation

- I=0 to I=1 contribution from Strong IB+EM effect (left), I=1 contribution EM effects (right)




## Interplay between Lattice and Experiment

- Check consistency between Lattice and R-ratio
- Short distance from Lattice, Long distance from R-ratio :
error already 0.5-1.2\% around $\mathrm{t}_{\text {lat } / \text { exp }}=2 \mathrm{fm}$

$$
a_{\mu}^{\mathrm{HVP}}=\left[\sum_{t=0}^{t_{\text {taterexp }}} w(t) C(t)\right]^{\mathrm{LAT}}+\left[\int_{t_{\text {tatetexp }}}^{\infty} d t w(t) C(t)\right]^{\operatorname{EXP}}
$$




## Euclidean time correlation from $e^{+} e^{-} R(s)$ data

From $e^{+} e^{-} R(s)$ ratio, using disparsive relation, zero-spacial momentum projected Euclidean correlation function $C(t)$ is obtained

$$
\begin{aligned}
\hat{\Pi}\left(Q^{2}\right) & =Q^{2} \int_{0}^{\infty} d s \frac{R(s)}{s\left(s+Q^{2}\right)} \quad \begin{array}{l}
\text { Lattice can compute Integral of } \\
\text { Inclusive cross sections accurately }
\end{array} \\
C^{\mathrm{R} \text {-ratio }(t)} & =\frac{1}{12 \pi^{2}} \int_{0}^{\infty} \frac{d \omega}{2 \pi} \hat{\Pi}\left(\omega^{2}\right) e^{i \omega t}=\frac{1}{12 \pi^{2}} \int_{0}^{\infty} d s \sqrt{s} R(s) e^{-\sqrt{s} t}
\end{aligned}
$$

- $C(t)$ or $w(t) C(t)$ are directly comparable to Lattice results with the proper limits ( $m_{q} \rightarrow m_{q}^{\text {phys }}, a \rightarrow 0, V \rightarrow \infty$, QED ...)
- Lattice: long distance has large statistical noise, (short distance: discretization error, removed by $a \rightarrow 0$ and/or pQCD )
- R-ratio : short distance has larger error

$\hat{\Pi}\left(Q^{2}\right)=Q^{2} \int_{0}^{\infty} d s \frac{R(s)}{s\left(s+Q^{2}\right)}$
( $1 / a=1.78 \mathrm{GeV}, \quad$ Relative statistical error)



## Comparison of R-ratio and Lattice [ F. Jegerlehner alphaQED 2016 ]

- Covariance matrix among energy bin in R-ratio is not available, assumes $100 \%$ correlated



## Combine R-ratio and Lattice [ Christoph Lehner et al PRL18]

- Use short and long distance from R-ratio using smearing function, and mid-distance from lattice

$$
\begin{aligned}
& \Theta(t, \mu, \sigma) \equiv[1+\tanh [(t-\mu) / \sigma]] / 2 \\
& a_{\mu}=\sum_{t} w_{t} C(t) \equiv a_{\mu}^{\mathrm{SD}}+a_{\mu}^{\mathrm{W}}+a_{\mu}^{\mathrm{LD}} \\
& {a_{\mu}}_{\mathrm{SD}}=\sum_{t} C(t) w_{t}\left[1-\Theta\left(t, t_{0}, \Delta\right)\right], \\
& a_{\mu}^{\mathrm{W}}=\sum_{t} C(t) w_{t}\left[\Theta\left(t, t_{0}, \Delta\right)-\Theta\left(t, t_{1}, \Delta\right)\right], \\
& \partial_{\mu}^{\mathrm{LD}}=\sum_{t} C(t) w_{t} \Theta\left(t, t_{1}, \Delta\right)
\end{aligned}
$$

How does this translate to the time-like region?


Most of $\pi \pi$ peak is captured by window from $t_{0}=0.4 \mathrm{fm}$ to $t_{1}=1.5 \mathrm{fm}$, so replacing this region with lattice data reduces the dependence on BaBar versus KLOE data sets.


## Continuum limit of $\mathbf{a}^{\mathbf{w}}$

Continuum limit of $a_{\mu}^{\mathrm{W}}$ from our lattice data; below $t_{0}=0.4 \mathrm{fm}$ and $\Delta=0.15 \mathrm{fm}$


## RBC/UKQCD [C. Lehner Lat17]

Continuum extrapolation is mild
c.f BMWc [K. Miura Lat17]


## R-ratio + Lattice


t1 dependence is flat => a consistency between R-ratio and Lattice $\mathrm{t} 1=1.2 \mathrm{fm}$, R-ratio : Lattice $=50: 50$
$\mathrm{t} 1=1.2 \mathrm{fm}$ current error (note $100 \%$ correlation in R -ratio) is minimum

## HVP results



- Significant improvements is in progress for statistical error using $2 \pi$ and $4 \pi$ (!) states in addition to EM current (GEVP, GS-parametrization)
- Checking finite volume and discretization error as well as Isospin V effects


## Example error budget from RBC/UKQCD 2018 (Fred’s alphaQED17 results used for window result)

| Window $t=[0.4,1 \mathrm{fm}]$ |  | Pure Lattice |
| :---: | :---: | :---: |
| $\overline{a_{\mu}}$ ud, conn, isospin | $202.9(1.4)_{\mathrm{S}}(0.2)_{\mathrm{C}}(0.1)_{\mathrm{V}}(0.2)_{\mathrm{A}}(0.2)_{\mathrm{Z}}$ | $649.7(14.2)_{\mathrm{S}}(2.8)_{\mathrm{C}}(3.7)_{\mathrm{V}}(1.5)_{\mathrm{A}}(0.4)_{\mathrm{Z}}(0.1)_{\mathrm{E} 48}(0.1)_{\mathrm{E} 64}$ |
| $a_{\mu}{ }^{\text {s, conn, isospin }}$ | $27.0(0.2)_{\mathrm{S}}(0.0)_{\mathrm{C}}(0.1)_{\mathrm{A}}(0.0)_{\mathrm{Z}}$ | $53.2(0.4)_{\mathrm{S}}(0.0)_{\mathrm{C}}(0.3)_{\mathrm{A}}(0.0)_{\mathrm{Z}}$ |
| $a_{\mu}{ }^{\mathrm{c}}$, conn, isospin | $3.0(0.0)_{\mathrm{S}}(0.1)_{\mathrm{C}}(0.0)_{\mathrm{Z}}(0.0)_{\mathrm{M}}$ | $14.3(0.0)_{\mathrm{S}}(0.7)_{\mathrm{C}}(0.1)_{\mathrm{Z}}(0.0)_{\mathrm{M}}$ |
| $a_{\mu}^{\text {uds, disc, isospin }}$ | $-1.0(0.1)_{\mathrm{S}}(0.0)_{\mathrm{C}}(0.0)_{\mathrm{V}}(0.0)_{\mathrm{A}}(0.0)_{\mathrm{Z}}$ | $-11.2(3.3)_{\mathrm{S}}(0.4)_{\mathrm{V}}(2.3)_{\mathrm{L}}$ |
| $a_{\mu}^{\text {QED, conn }}$ | $0.2(0.2)_{\mathrm{S}}(0.0)_{\mathrm{C}}(0.0)_{\mathrm{V}}(0.0)_{\mathrm{A}}(0.0)_{\mathrm{Z}}(0.0)_{\mathrm{E}}$ | $5.9(5.7)_{\mathrm{S}}(0.3)_{\mathrm{C}}(1.2)_{\mathrm{V}}(0.0)_{\mathrm{A}}(0.0)_{\mathrm{Z}}(1.1)_{\mathrm{E}}$ |
| $a_{\mu}^{\text {SIB }}$, disc | $-0.2(0.1)_{\mathrm{S}}(0.0)_{\mathrm{C}}(0.0)_{\mathrm{V}}(0.0)_{\mathrm{A}}(0.0)_{\mathrm{Z}}(0.0)_{\mathrm{E}}$ | $-6.9(2.1)_{\mathrm{S}}(0.4)_{\mathrm{C}}(1.4)_{\mathrm{V}}(0.0)_{\mathrm{A}}(0.0)_{\mathrm{Z}}(1.3)_{\mathrm{E}}$ |
| $a_{\mu}{ }^{\text {SIB }}$ | $0.1(0.2)_{\mathrm{S}}(0.0)_{\mathrm{C}}(0.2)_{\mathrm{V}}(0.0)_{\mathrm{A}}(0.0)_{\mathrm{Z}}(0.0)_{\mathrm{E} 48}$ | $10.6(4.3)_{\mathrm{S}}(0.6)_{\mathrm{C}}(6.6)_{\mathrm{V}}(0.1)_{\mathrm{A}}(0.0)_{\mathrm{Z}}(1.3)_{\mathrm{E} 48}$ |
| $a_{\mu}^{\text {udsc, isospin }}$ | $231.9(1.4)_{\mathrm{S}}(0.2)_{\mathrm{C}}(0.1)_{\mathrm{V}}(0.3)_{\mathrm{A}}(0.2)_{\mathrm{Z}}(0.0)_{\mathrm{M}}$ | $\begin{aligned} & 705.9(14.6)_{\mathrm{S}}(2.9)_{\mathrm{C}}(3.7)_{\mathrm{V}}(1.8)_{\mathrm{A}}(0.4)_{\mathrm{Z}}(2.3)_{\mathrm{L}}(0.1)_{\mathrm{E} 48} \\ & \quad(0.1)_{\mathrm{E} 64}(0.0)_{\mathrm{M}} \end{aligned}$ |
| $\begin{aligned} & a_{\mu}^{\mathrm{QED}, \mathrm{SIB}} \\ & a_{\mu}^{\mathrm{R}-\text { ratio }} \end{aligned}$ | $\frac{0.1(0.3)_{\mathrm{S}}(0.0)_{\mathrm{C}}(0.2)_{\mathrm{V}}(0.0)_{\mathrm{A}}(0.0)_{\mathrm{Z}}(0.0)_{\mathrm{E}}(0.0)_{\mathrm{E} 48}}{460.4(0.7)_{\mathrm{RST}}(2.1)_{\mathrm{RSY}}}$ | $9.5(7.4)_{\mathrm{S}}(0.7)_{\mathrm{C}}(6.9)_{\mathrm{V}}(0.1)_{\mathrm{A}}(0.0)_{\mathrm{Z}}(1.7)_{\mathrm{E}}(1.3)_{\mathrm{E} 48}$ |
| $a_{\mu}$ | $\begin{aligned} & 692.5(1.4)_{\mathrm{S}}(0.2)_{\mathrm{C}}(0.2)_{\mathrm{V}}(0.3)_{\mathrm{A}}(0.2)_{\mathrm{Z}}(0.0)_{\mathrm{E}}(0.0)_{\mathrm{E} 48} \\ & (0.0)_{\mathrm{b}}(0.1)_{\mathrm{c}}(0.0)_{\overline{\mathrm{S}}}(0.0)_{\overline{\mathrm{Q}}}(0.0)_{\mathrm{M}}(0.7)_{\mathrm{RST}}(2.1)_{\mathrm{RSY}} \end{aligned}$ | $\begin{gathered} 715.4(16.3)_{\mathrm{S}}(3.0)_{\mathrm{C}}(7.8)_{\mathrm{V}}(1.9)_{\mathrm{A}}(0.4)_{\mathrm{Z}}(1.7)_{\mathrm{E}}(2.3)_{\mathrm{L}} \\ (1.5)_{\mathrm{E} 48}(0.1)_{\mathrm{E} 64}(0.3)_{\mathrm{b}}(0.2)_{\mathrm{c}}(1.1)_{\mathrm{S}}(0.3)_{\overline{\mathrm{Q}}}(0.0)_{\mathrm{M}} \end{gathered}$ |

TABLE I. Individual and summed contributions to $a_{\mu}$ multiplied by $10^{10}$. The left column lists results for the window method with $t_{0}=0.4 \mathrm{fm}$ and $t_{1}=1 \mathrm{fm}$. The right column shows results for the pure first-principles lattice calculation. The respective uncertainties are defined in the main text.

For the pure lattice number the dominant errors are (S) statistics, (V) finite-volume errors, and (C) the continuum limit extrapolation uncertainty.
For the window method there are additional R -ratio systematic (RSY) and R-ratio statistical (RST) errors.

## Hadronic Light-by-Light (HLbL) contributions



## HLbL from Models

- Model estimate with non-perturbative constraints at the chiral / low energy limits using anomaly : (9-12) x $10^{-10}$ with $25-40 \%$ uncertainty

$$
a_{\mu}^{\exp }-a_{\mu}^{\mathrm{SM}}=28.8(6.3)_{\exp }(4.9)_{\mathrm{SM}} \times 10^{-10} \quad[3.6 \sigma]
$$

F. Jegerlehner, x $10^{11}$


| Contribution | BPP | HKS | KN | MV | PdRV | N/JN |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\pi^{0}, \eta, \eta^{\prime}$ | $85 \pm 13$ | $82.7 \pm 6.4$ | $83 \pm 12$ | $114 \pm 10$ | $114 \pm 13$ | $99 \pm 16$ |
| $\pi, K$ loops | $-19 \pm 13$ | $-4.5 \pm 8.1$ | - | $0 \pm 10$ | $-19 \pm 19$ | $-19 \pm 13$ |
| axial vectors | $2.5 \pm 1.0$ | $1.7 \pm 1.7$ | - | $22 \pm 5$ | $15 \pm 10$ | $22 \pm 5$ |
| scalars | $-6.8 \pm 2.0$ | - | - | - | $-7 \pm 7$ | $-7 \pm 2$ |
| quark loops | $21 \pm 3$ | $9.7 \pm 11.1$ | - | - | 2.3 | $21 \pm 3$ |
| total | $83 \pm 32$ | $89.6 \pm 15.4$ | $80 \pm 40$ | $136 \pm 25$ | $105 \pm 26$ | $116 \pm 39$ |

## Coordinate space Point photon method

[ Luchang Jin et all. , PRD93, 014503 (2016) ]

- Treat all 3 photon propagators exactly (3 analytical photons), which makes the quark loop and the lepton line connected :
disconnected problem in Lattice QED+QCD -> connected problem with analytic photon
- QED 2-loop in coordinate space. Stochastically sample, two of quark-photon vertex location $x, y, z$ and $x_{o p}$ is summed over space-time exactly


- Short separations, $\operatorname{Min}[|x-z|,|y-z|,|x-y|]<R \sim O(0.5) f m$, which has a large contribution due to confinement, are summed for all pairs
- longer separations, Min $[|x-z|,|y-z|,|x-y|]>=R$, are done stochastically with a probability shown above ( Adaptive Monte Carlo sampling )


## Dramatic Improvement! Luchang Jin

$a=0.11 \mathrm{fm}, 24^{3} \times 64(2.7 \mathrm{fm})^{3}$,
$\mathrm{m}_{\pi}=329 \mathrm{MeV}, \quad \mathrm{m}_{\mu}=\sim 190 \mathrm{MeV}, \mathrm{e}=1$

$$
\begin{array}{r}
q=2 \pi / L N_{\text {prop }}=81000 \longmapsto \vdash \\
q=0 N_{\text {prop }}=26568 \longmapsto \bigcirc
\end{array}
$$



## SU(3) hierarchies for d-HLbL

- At $\mathrm{m}_{\mathrm{s}}=\mathrm{m}_{\mathrm{ud}}$ limit, following type of disconnected HLbL diagrams survive $Q_{u}+Q_{d}+Q_{s}=0$
- Physical point run using similar techniques to c-HLbL.
- other diagrams suppressed by $O\left(m_{s}-m_{u d}\right) / 3$ and $O\left(\left(m_{s}-m_{u d}\right)^{2}\right)$



## 140 MeV Pion, connected and disconnected LbL results

[ Luchang Jin et al. , Phys.Rev.Lett. 118 (2017) 022005]

- left: connected, right : leading disconnected

- Using AMA with 2,000 zMobius low modes, AMA
( statistical error only )

$$
r=|x-y|
$$

$$
\begin{array}{ll}
\left.\frac{g_{\mu}-2}{2}\right|_{\mathrm{cHLbL}} & =(0.0926 \pm 0.0077) \times\left(\frac{\alpha}{\pi}\right)^{3}=(11.60 \pm 0.96) \times 10^{-10} \\
\left.\frac{g_{\mu}-2}{2}\right|_{\mathrm{dHLbL}} & =(-0.0498 \pm 0.0064) \times\left(\frac{\alpha}{\pi}\right)^{3}=(-6.25 \pm 0.80) \times 10^{-10} \\
\left.\frac{g_{\mu}-2}{2}\right|_{\mathrm{HLbL}} & =(0.0427 \pm 0.0108) \times\left(\frac{\alpha}{\pi}\right)^{3}=(5.35 \pm 1.35) \times 10^{-10}
\end{array}
$$

## Lattice 2017 Updates from PRL (2017)

- Discretization error
$\rightarrow$ a scaling study for $1 / \mathrm{a}=2.7,1.4,1.0 \mathrm{GeV}$ at physical quark mass for both connected and disconnected is being finalized
- Finite volume

QED_L (photon/lepton in a box) [ 08 Hayakawa Uno ]
Infinite Volume and continuum lepton + photon diagrams



?

## Summary

- Lattice calculation for g-2 calculation is improved very rapidly
- HVP
- New methods using low mode for connected at physical quark mass,
- disconnected quark loop at physical quark mass, QED and IB studies are included
- Combining with R-ratio experiment data for cross-check and improvement => $0.4 \%$ error
- Eventually the window will be enlarged for a pure LQCD prediction (currently 2.6 \% error)
- Significant improvements is in progress for statistical error using $2 \pi$ and $4 \pi$ (!) states in addition to EM current (GEVP, GS-parametrization)
- Checking finite volume and discretization error as well as Isospin V effects
- We could compute Inclusive hadron cross sections at Euclidean $q^{\wedge} 2$ from the first principle Lattice QCD with Isospin breaking effects ! e+e- -> hadron tau -> nu + hadrons tau inclusive decay and |Vus| arXiv:1803.07228 (to appear in PRL)
- HLbL
- computing connected and leading disconnected diagrams :
-> $8 \%$ stat error in connected, 13 \% stat error in leading disconnected
- coordinate-space integral using analytic photon propagator with adaptive probability (point photon method), config-by-config conserved external current
- Improving statistics right now.
- Various size of Lattice ensemble / method for systematic error as well as higher disconnected diagram Comparing with Mainz group's results (for connected at heavy pion mass)
- Goal : HVP sub 1\% (then 0.25\%), HLbL 10\% error

Can we see the next physics Revolution (c.f GW ) ?


## Simulation details [RBC/UKQCD 2015]

two gauge field ensembles generated by RBC/UKQCD collaborations

Domain wall fermions: chiral symmetry at finite $a$

Iwasaki Gauge action (gluons)

- pion mass $m_{\pi}=139.2(2)$ and 139.3(3) $\mathrm{MeV}\left(m_{\pi} L \lesssim 4\right)$
- lattice spacings $a=0.114$ and 0.086 fm
- lattice scale $a^{-1}=1.730$ and 2.359 GeV
- lattice size $L / a=48$ and 64
- lattice volume $(5.476)^{3}$ and $(5.354)^{3} \mathrm{fm}^{3}$

Use all-mode-average (AMA) [Blum et al 2012] and low-mode- averaging (LMA) [Giusti et al, 2004, Degrand et al 2005, Lehner 2016 for HVP] techniques for improved statistics by more than three orders of magnitudes compared to basic CG, and $\times 10$ smaller memory via multigrid-Lanczos [Lehner 2017] .

## Conserved current \& moment method

- [conserved current method at finite q2] To tame UV divergence, one of quark-photon vertex (external current) is set to be conserved current (other three are local currents). All possible insertion are made to realize conservation of external currents config-by-config.

- [moment method, $\mathrm{q} 2 \rightarrow 0$ ] By exploiting the translational covariance for fixed external momentum of lepton and external EM field, $q->0$ limit value is directly computed via the first moment of the relative coordinate, $\mathrm{xop}-(\mathrm{x}+\mathrm{y}) / 2$, one could show

$$
\left.\frac{\partial}{\partial q_{i}} \mathcal{M}_{\nu}(\vec{q})\right|_{\vec{q}=0}=i \sum_{x, y, z, x_{\mathrm{op}}}\left(x_{\mathrm{op}}-(x+y) / 2\right)_{i} \times
$$

to directly get $F_{2}(0)$ without extrapolation.


$$
\text { Form factor: } \Gamma_{\mu}(q)=\gamma_{\mu} F_{1}\left(q^{2}\right)+\frac{i \sigma^{\mu \nu} q_{\nu}}{2 m_{l}} F_{2}\left(q^{2}\right)
$$

## Current conservation \& subtractions

- conservation => transverse tensor

$$
\Pi^{\mu \nu}(q)=\left(\hat{q}^{2} \delta^{\mu \nu}-\hat{q}^{\mu} \hat{q}^{\nu}\right) \Pi\left(\hat{q}^{2}\right)
$$

- In infinite volume, $\mathrm{q}=0, \Pi_{\mu \mathrm{v}}(\mathrm{q})=0$
- For finite volume, $\Pi_{\mu \mathrm{v}}(0)$ is exponentially small (L.Jin, use also in HLbL)

$$
\begin{aligned}
& \int_{V} d x^{4}\left\langle V_{\mu}(x) \mathcal{O}(0)\right\rangle=\int_{V} d x^{4} \partial_{x}\left(x\left\langle V_{\mu}(x) \mathcal{O}(0)\right\rangle\right) \\
= & \int_{\partial V} d x^{3} x\left\langle V_{\mu}(x) \mathcal{O}(0)\right\rangle \propto L^{4} \exp (-M L / 2) \rightarrow 0
\end{aligned}
$$

- e.g. DWF L=2, 3, $5 \mathrm{fm} \quad \Pi_{\mu v}(0)=8(3) \mathrm{e}-4,2(13) \mathrm{e}-5,-1(5) \mathrm{e}-8$
- Subtract $\Pi_{\mu \mathrm{v}}(0)$ alternates FVE, and reduce stat error "-1" subtraction trick :
$\Pi^{\mu \nu}(q)-\Pi^{\mu \nu}(0)=\int d^{4} x\left(e^{i q x}-1\right)\left\langle J^{\mu}(x) J^{\nu}(0)\right\rangle$

