

HVP contribution to $(g - 2)_\mu$: status of the Mainz calculation

Antoine Gérardin

In collaboration with M. Cè, G. von Hippel, B. Hoerz, T. San José, S. Kuberski,
H.B. Meyer, K. Miura, D. Mohler, K. Ottnad, J. Wilhelm, H. Wittig

30 June 2021

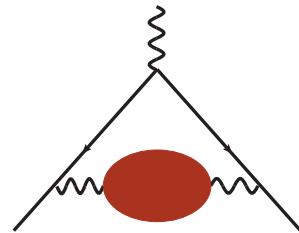


Specificities of our setup

- Time-momentum representation [Bernecker, Meyer '11 '13]

$$a_\mu^{\text{hvp}} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^\infty dt \ \tilde{K}(t) \ G(t)$$

$$G(t) = -\frac{1}{3} \sum_{k=1}^3 \sum_{\vec{x}} \langle V_k(x) V_k(0) \rangle$$



Electromagnetic current : $V_\mu(x) = \frac{2}{3}\bar{u}(x)\gamma_\mu u(x) - \frac{1}{3}\bar{d}(x)\gamma_\mu d(x) - \frac{1}{3}\bar{s}(x)\gamma_\mu s(x) + \dots$

- $N_f = 2 + 1$ with Wilson fermions : $\mathcal{O}(a)$ -improved (action + currents)
- Two different discretizations of the vector current

$$J_\mu^{(l),a}(x) = \bar{\psi}(x)\gamma_\mu \frac{\lambda^a}{2} \psi_j(x),$$

$$J_\mu^{(c),a}(x) = \frac{1}{2} \left(\bar{\psi}(x + a\hat{\mu})(1 + \gamma_\mu) U_\mu^\dagger(x) \frac{\lambda^a}{2} \psi(x) - \bar{\psi}(x)(1 - \gamma_\mu) U_\mu(x) \frac{\lambda^a}{2} \psi(x + a\hat{\mu}) \right)$$

→ combined continuum extrapolation

Window quantities

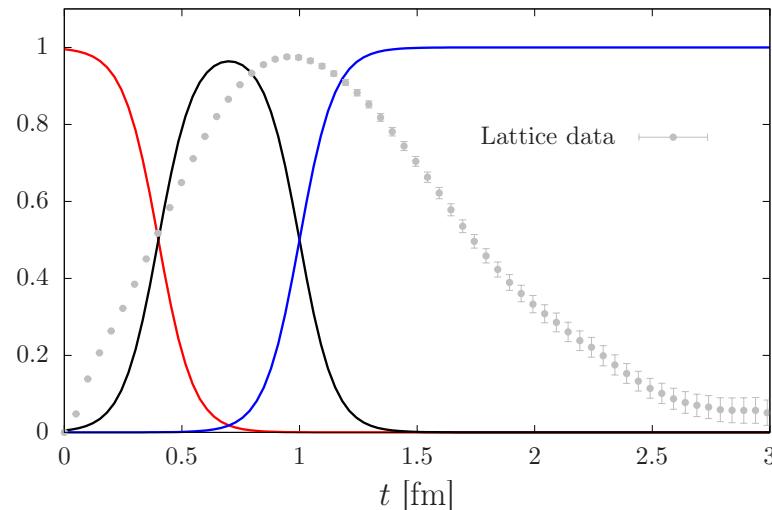
► Time-momentum representation

[Bernecker, Meyer '11 '13]

$$a_\mu^{\text{win}} = \left(\frac{\alpha}{\pi}\right)^2 \sum_t G(t) \tilde{K}(t) W(t; t_0, t_1),$$

- Short distances : $W^{\text{SD}}(t; t_0) = [1 - \Theta(t, t_0, \Delta)]$
- Intermediate distances : $W^{\text{ID}}(t; t_0, t_1) = [\Theta(t, t_0, \Delta) - \Theta(t, t_1, \Delta)]$
- Long distances : $W^{\text{LD}}(t; , t_1) = \Theta(t, t_1, \Delta)$

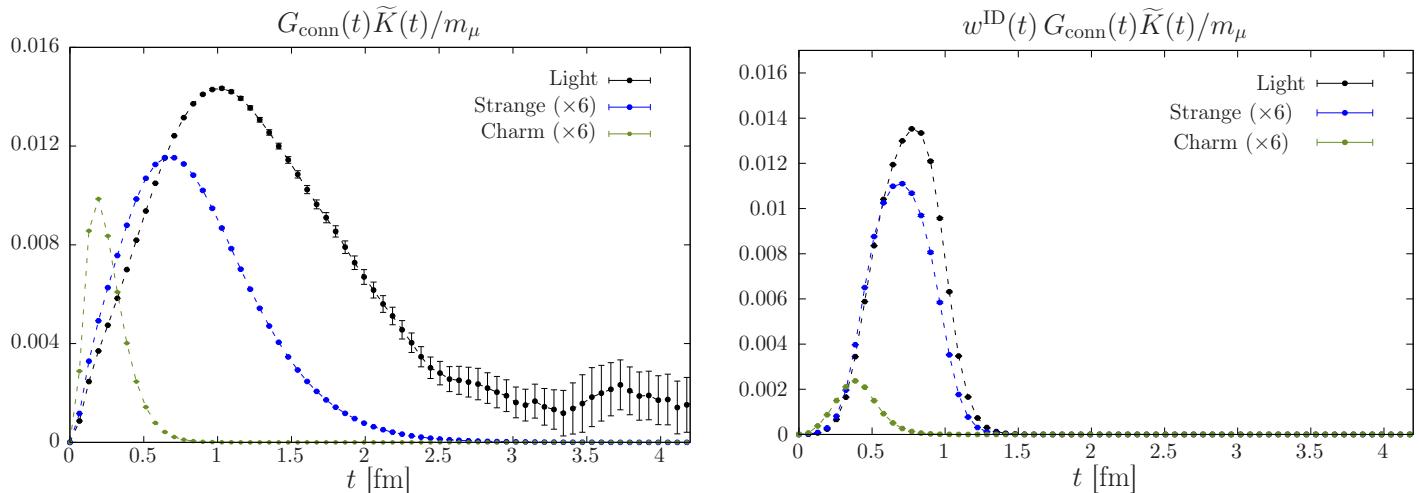
Smooth step function : $\Theta(t, t', \Delta) = [1 + \tanh[(t - t')/\Delta]] / 2$



► This talk focuses on the intermediate window

- no signal-to-noise ratio issue

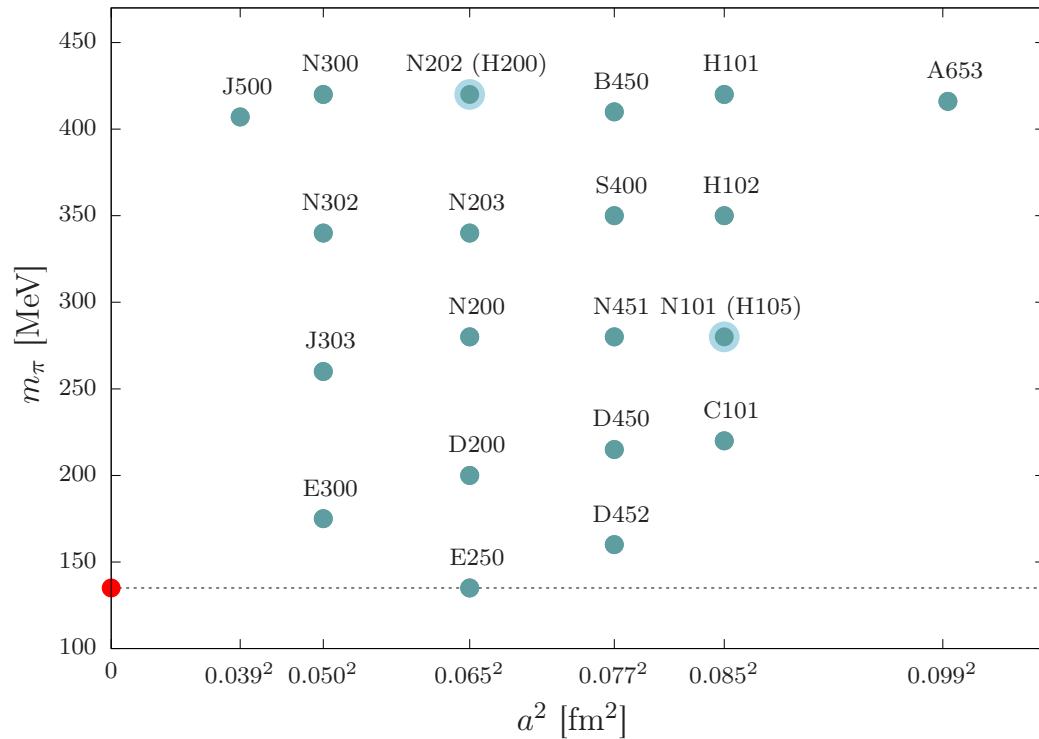
→ 1-2 permille statistical precision can be reached on the integrand
 → the treatment of the noisy long-distance part of the correlator is irrelevant here
 → it means that the statistical error is not the main source of uncertainty



- need to have a good control over other systematics :
 continuum extrapolation, finite-size effects

CLS ensembles used in this work

- > 35 000 gauge field configurations for 22 ensembles
- include the physical pion mass
- 6 lattice spacings



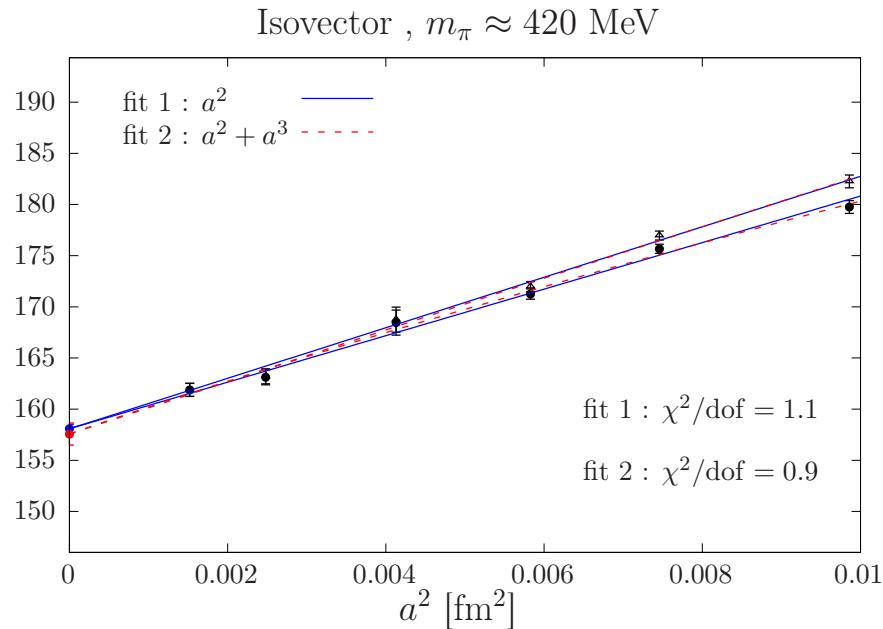
Continuum extrapolation

- The continuum extrapolation is one of the main sources of error

→ Dedicated study performed at $m_\pi = m_K \approx 420 MeV}$

→ 6 lattice spacings < 0.1 fm

→ 2 different discretizations of the vector current to check $O(a)$ -improvement



- We observe a very good scaling within a large range of lattice spacings

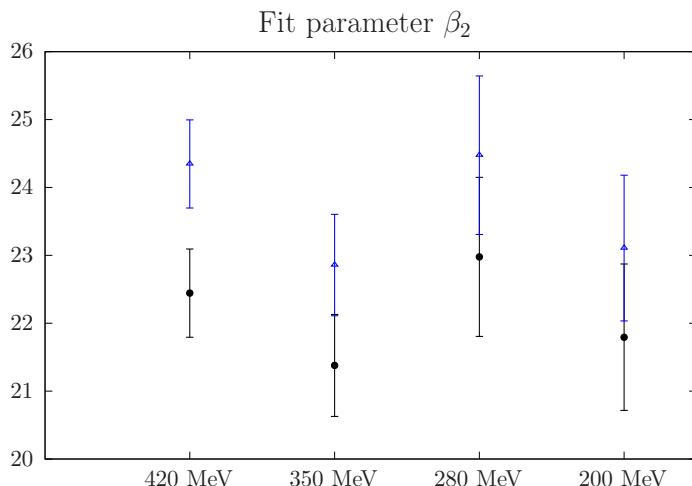
Continuum extrapolation

- What about lighter pion masses ?

→ same fit ansatz :

$$a_\mu^{\text{hvp,I1}}(a, d) = a_\mu^{\text{hvp,I1}}(0) + \beta_2^{(d)} \tilde{a}^2 + \beta_3^{(d)} \tilde{a}^3$$

→ dimensionless parameter $\tilde{a} = a/a_{\beta=3.34}$ and $d = (ll), (lc)$



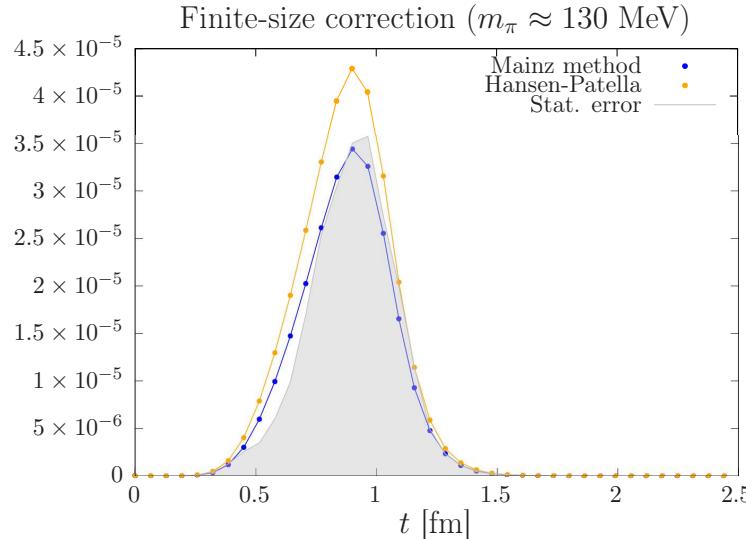
Global fit (22 ensembles)

$$\begin{aligned}\beta_2^{(ll)} &= 27(5), & \beta_3^{(ll)} &= -5(5) \\ \beta_2^{(lc)} &= 25(5), & \beta_3^{(lc)} &= -0(5)\end{aligned}$$

→ compatible values
→ β_3 compatible with zero

- No significant pion mass dependence of slope parameter β_2

- We are currently accumulating statistics on a second ensemble at our finest lattice spacing with smaller pion mass



- Mainz method [D.Bernecker, H.B.Meyer '11]
 - used in [Phys.Rev.D 100 (2019)]
 - NLO ChPT for $t < 1.5$ fm
- Hansen-Patella method [M.T Hansen, A. Patella '19 '20]

► Finite-volume correction for full a_μ^{hyp} :

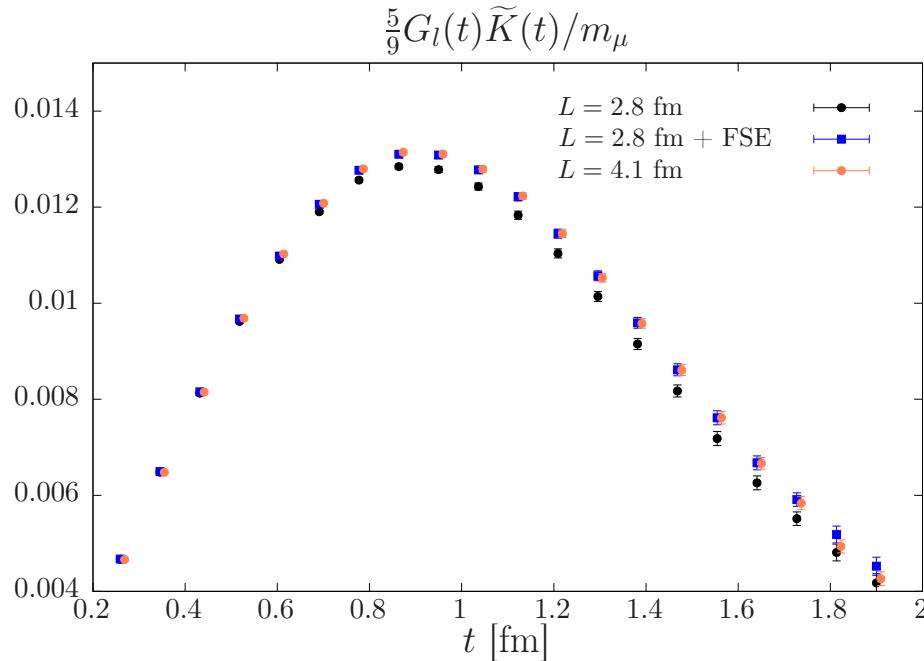
$$m_\pi L = 4 \quad (L = 6.2 \text{ fm}) \quad \Rightarrow \quad \Delta a_\mu = 22.6 \times 10^{-10} \quad (3\%)$$

► Intermediate window : $\Delta a_\mu = 0.6 \times 10^{-10}$ (0.25%)

→ FSE are much smaller for the intermediate window

→ but not negligible compared to the statistical precision ($\sim 0.2\%$)

► Wilson quarks : this is the **only** correction applied to the raw lattice data !



- Our procedure works remarkably well at $m_\pi = 280$ MeV [PRD 100 (2019) 014510]

$$m_\pi L = 4 \quad (L = 6.2 \text{ fm}) \quad \Rightarrow \quad \Delta a_\mu = 22.6 \times 10^{-10} \quad (3\%)$$

- In agreement with other direct lattice calculations [E. Shintani, Y. Kuramashi '19], [BMW '20]

Isospin decomposition : error budget

► Scale setting

→ one of the main challenges for a_μ^{hyp} : $\frac{\delta a_\mu^{\text{hyp}}}{a_\mu^{\text{hyp}}} = \left| \frac{m_\mu}{a_\mu^{\text{hyp}}} \frac{da_\mu^{\text{hyp}}}{dm_\mu} \right| \frac{\delta a}{a} \approx 1.8 \frac{\delta a}{a}$

→ window quantities : the absolute scale also enters the definition of the window

$$\frac{\delta a_\mu^{\text{hyp,win}}}{a_\mu^{\text{hyp,win}}} \approx \frac{m_\mu}{a_\mu^{\text{hyp,win}}} \frac{da_\mu^{\text{hyp,win}}}{dm_\mu} \frac{\delta a}{a} + \left(\frac{\alpha}{\pi} \right)^2 \int_0^\infty dt \frac{\delta w(t)}{\delta a} G(t) \tilde{K}(t) \times \frac{\delta a}{a_\mu^{\text{hyp,win}}}$$

- significantly reduces the sensitivity to the scale for the isovector contribution
- but makes it worse for the (small !) charm quark contribution

► Preliminary error budget

Table – Isovector

Type	%
Renormalization	< 0.01
Correlator	0.2
Finite-volume	≤ 0.2
Extrap.	0.61
Scale setting	0.38
Total	0.77

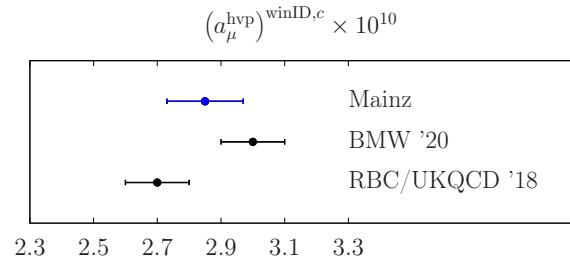
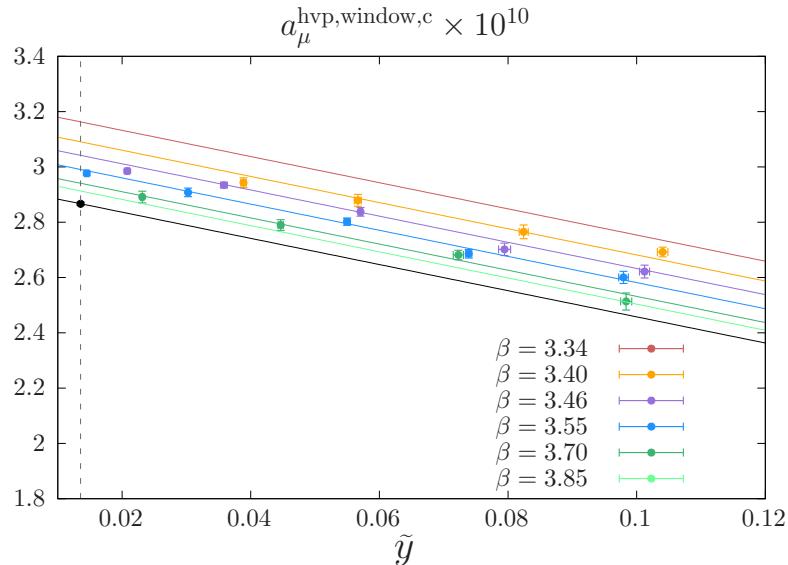
Table – Isoscalar

Type	%
Renormalization	< 0.01
Correlator	0.3
Extrap.	1.3
Scale setting	0.86
Total	1.59

► Final stage of the analysis

Charm quark contribution

- Extrapolation to the physical point $\tilde{y} \sim m_\pi^2$

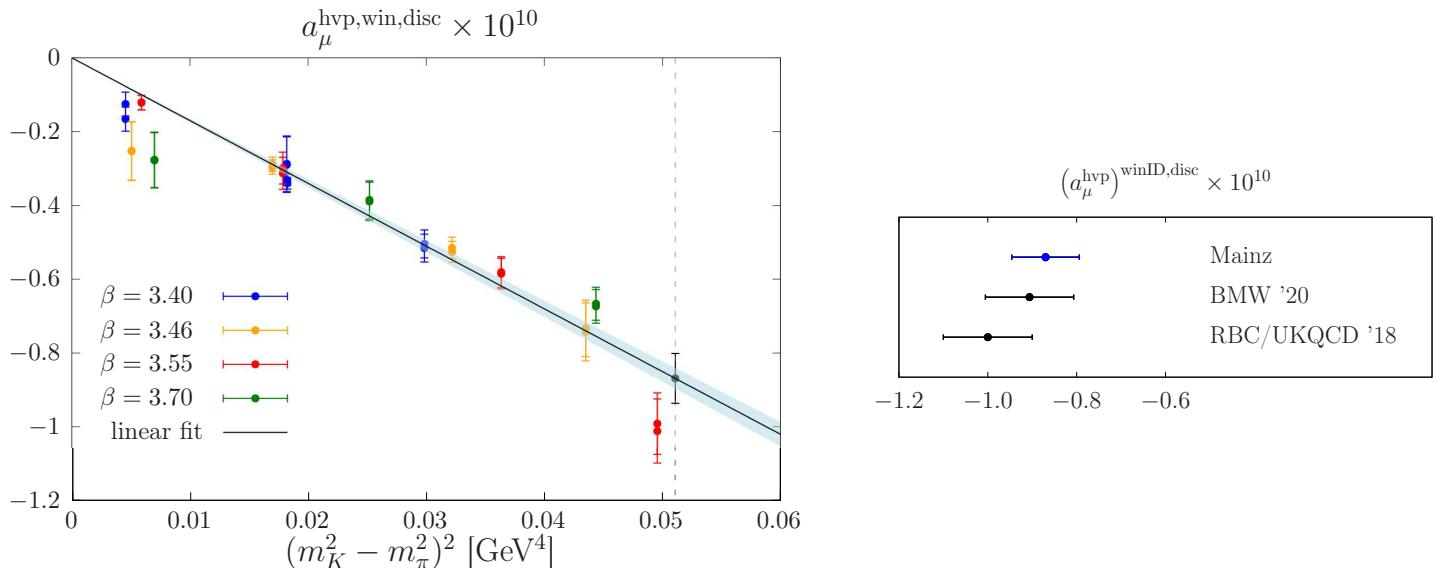


- Error budget for the charm quark contribution

Type	%
Renormalization	< 0.05
Tuning of κ_c	< 0.2
Correlator	< 0.05
Extrap.	1.7
Scale setting	3.9
Total	4.3

→ much more sensitive to the scale than $a_\mu^{\text{hvp,c}}$

Quark disconnected contribution



- as for the light quark contribution : there is no signal/noise problem here
- small contribution (about 7% of the total disconnected contribution to full a_μ^{hvp})
- small discretization effects

- ▶ **Large statistics** : 22 ensembles including the physical pion mass
 - statistical error is sub-dominant for the intermediate window
 - the treatment of the tail of the correlator is also irrelevant here
- ▶ **Continuum limit**
 - 6 lattice spacings. **All of them $< 0.1 \text{ fm}$.**
 - our 2 finest lattice spacings are smaller than any other lattice collaboration (for $g - 2$)
 - fully non-perturbative $\mathcal{O}(a)$ -improvement
 - **two discretizations of the vector current** to test the consistency of the continuum limit and $\mathcal{O}(a)$ -improvement
- ▶ **Ensembles with different volumes to check FSE**
 - this is the **only correction** applied to the raw lattice data
- ▶ **Ongoing works**
 - estimate of the quenching of the charm
 - isospin-breaking corrections
 - running of alpha / weak mixing angle

Thank you.