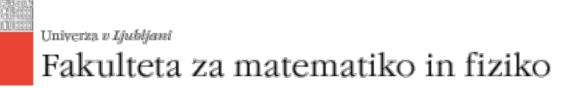


Interplay Between LHC and Flavor Physics

Jernej F. Kamenik





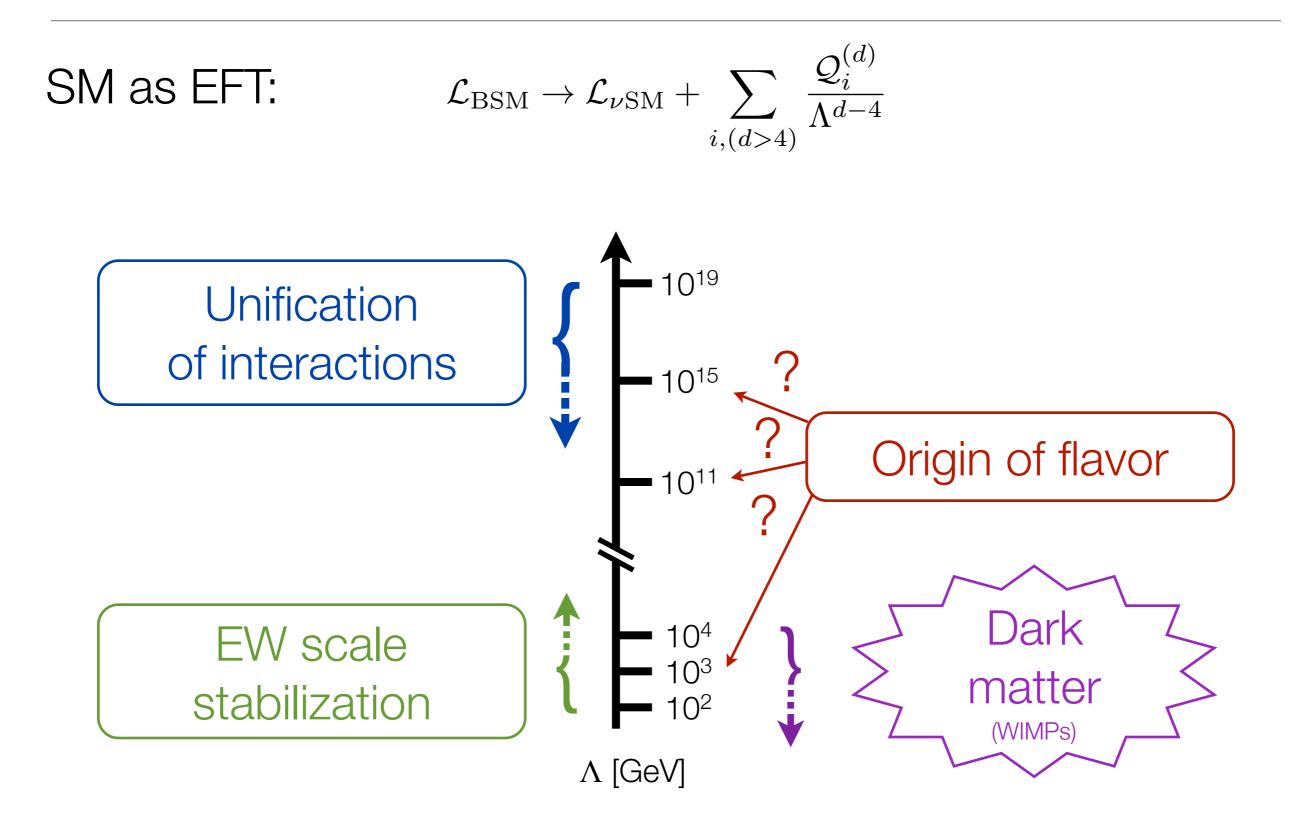
Nagoya 7/1/2017

Outline

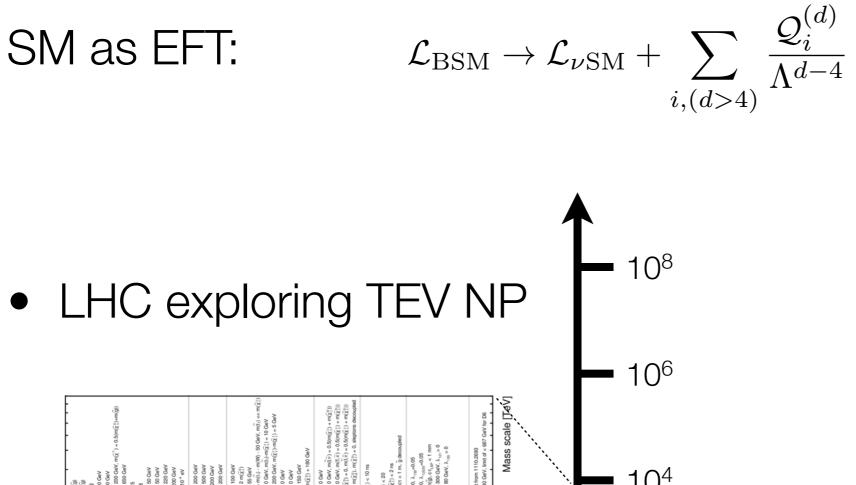
Flavor & high p_T physics interplay in several ways

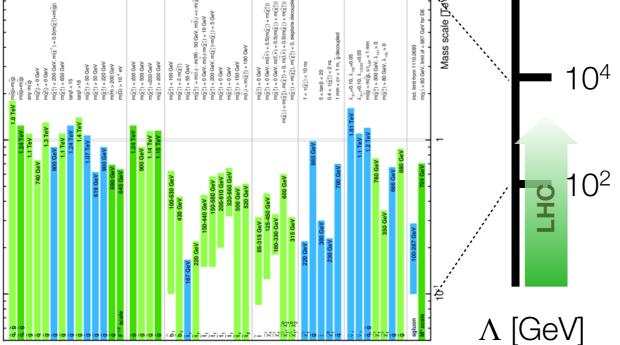
- Complementary constraints on NP models from low energy precision observables vs. high p_T searches
- Nontrivial flavor structure modifies signatures at LHC
- Anomalies in *B/D/K* physics motivate NP searches at high *p_T*
 - goes also the other way (e.g., $h \rightarrow \tau \mu$)

Flavor bounds on NP vs. LHC reach



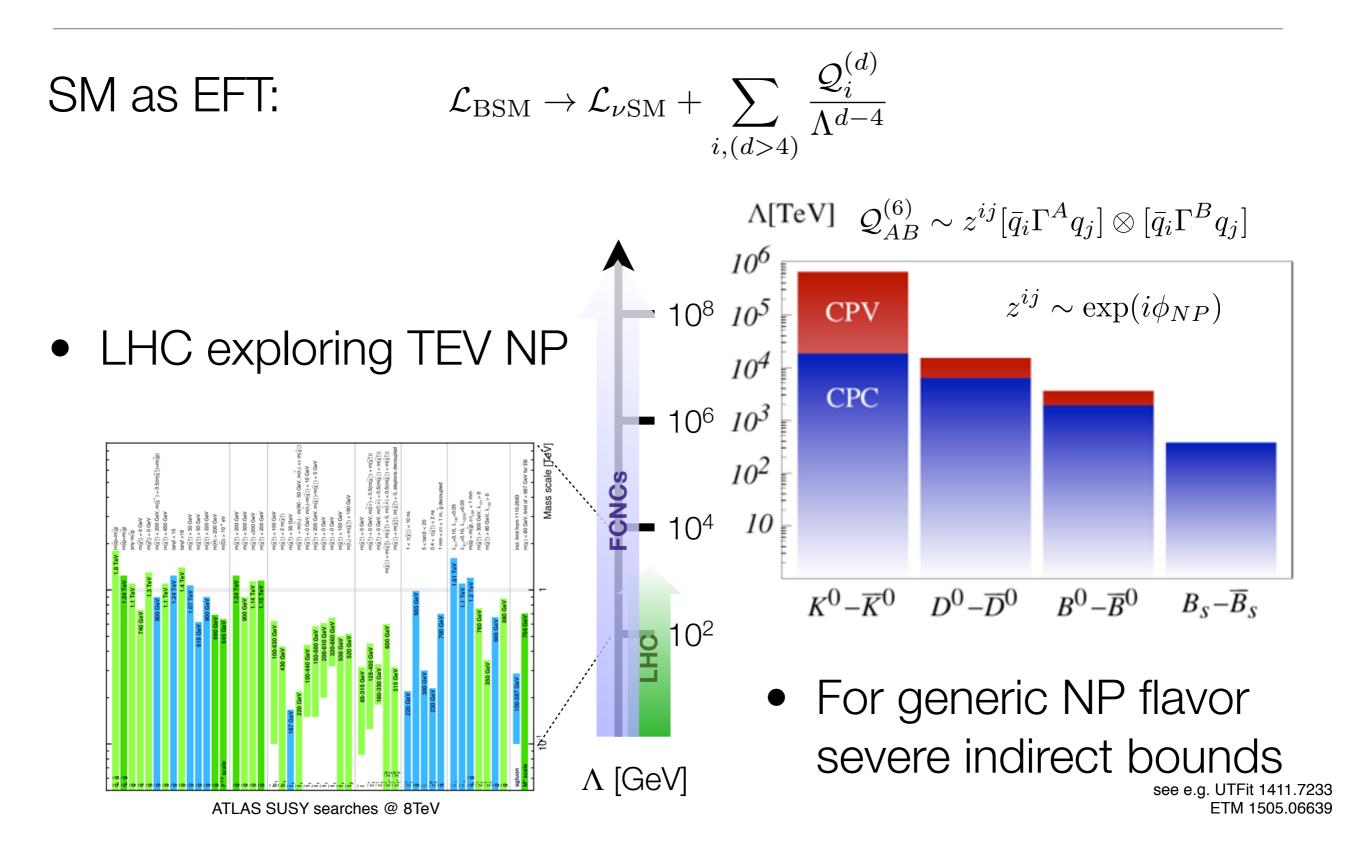
Flavor bounds on NP vs. LHC reach



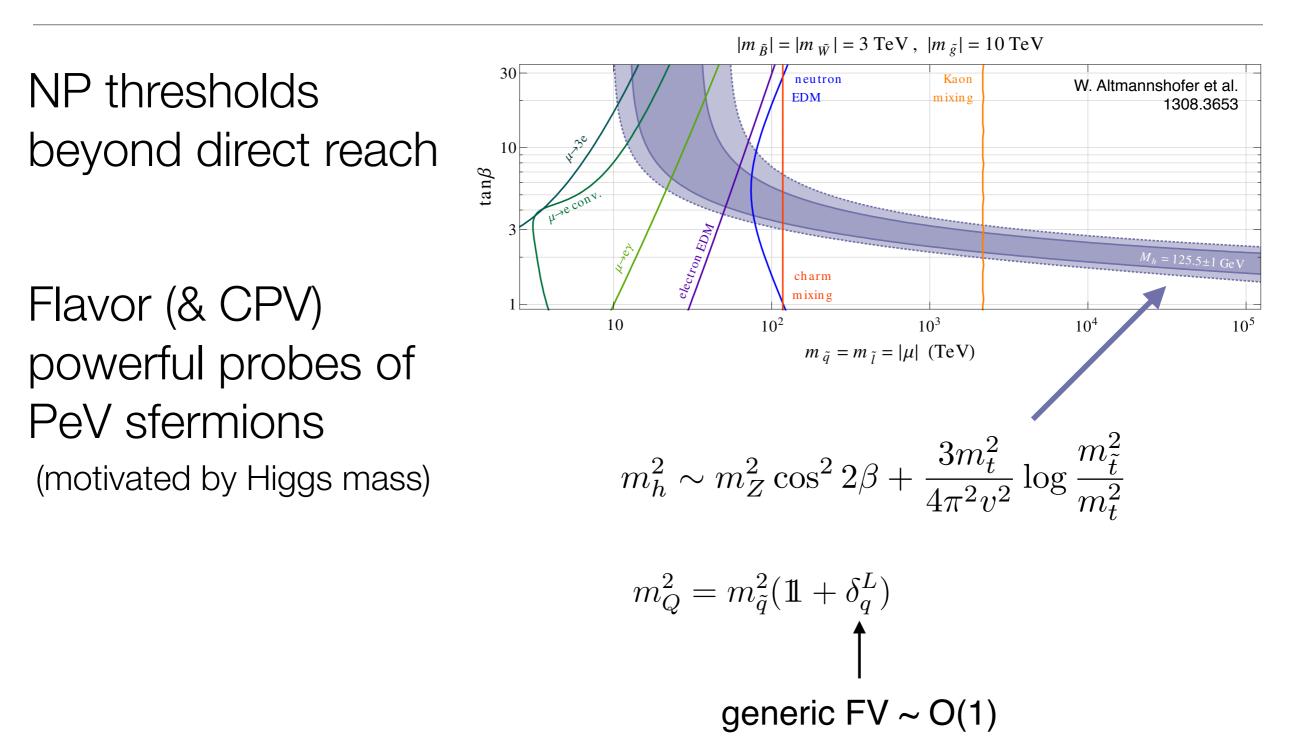


ATLAS SUSY searches @ 8TeV

Flavor bounds on NP vs. LHC reach



LHC bad dream scenario: (mini)split SUSY



LHC bad dream scenario: (mini)split SUSY

30 NP thresholds neutron Kaon W. Altmannshofer et al. EDM m ixin g 1308.3653 beyond direct reach 10 Now $\tan\beta$ EDM electron] charm Flavor (& CPV) m ixin g 30 neutron electron EDM EDM powerful probes of 10 PeV sfermions Kaon ~2025 $an\beta$ mixing charm mixing Significant 10^{2} 10^{5} 10 10^{3} 10^{4} $m_{\tilde{q}} = m_{\tilde{i}} = |\mu|$ (TeV) improvements (+generic FV) **MEG upgrade** Mu3e Mu2e expected in next LHCb 25fb⁻¹ decade

 $|m_{\tilde{B}}| = |m_{\tilde{W}}| = 3 \text{ TeV}, \ |m_{\tilde{e}}| = 10 \text{ TeV}$

Flavor & high-p₇ as complementary NP probes

Flavor safe NP? Flavor already broken in SM (Higgs).

$$\mathcal{H}_{mat} = \underbrace{\frac{c_{RL}^{IJ}}{\Lambda^n} H^{\dagger} \bar{D}^I Q^J \times X}_{An} + \frac{c_{LR}^{IJ}}{\Lambda^n} H \bar{Q}^I D^J \times X}_{An} + \underbrace{\frac{c_{LL}^{IJ}}{\Lambda^n} \bar{Q}^I Q^J \times X}_{An} + \underbrace{\frac{c_{RR}^{IJ}}{\Lambda^n} \bar{D}^I D^J \times X}_{An} +$$

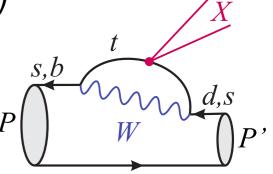
Any (additional) scalar coupling to SM fermions introduces additional breaking (can be aligned with Higgs)

New (massive) vectors coupling to SM fermionic currents can preserve flavor

 \Rightarrow FCNCs loop & GIM suppressed (as in SM)

see e.g. J.F.K. & C. Smith 1111.6402

$$c^{IJ} \to (g/(4\pi))^2 V_{tI}^* V_{tJ} \times c^{33}$$



D'Ambrosio et al. hep-ph/0207036

Flavor & high-p_T as complementary NP probes

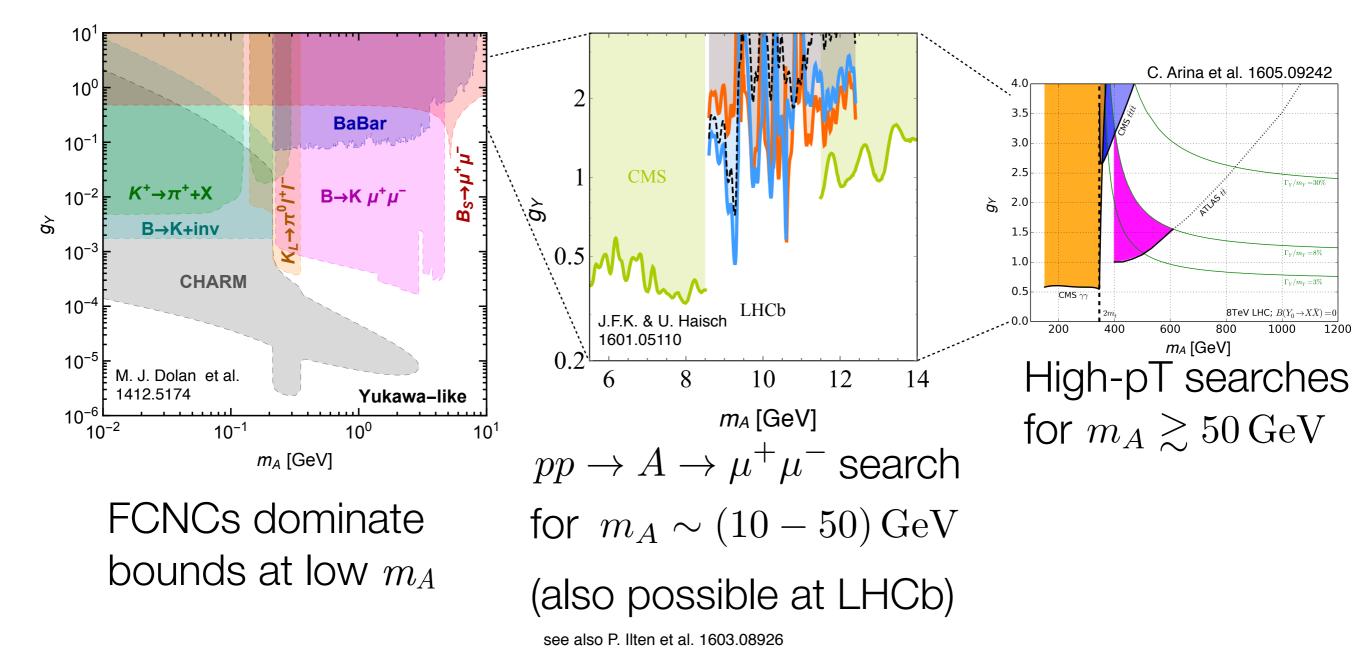
Example: simplified DM models with (pseudo)scalar mediators

$$\mathcal{L}_{\rm DM} = i \, g_{\chi} \, A \, \bar{\chi} \gamma^5 \chi \, + \sum_{f=q,\ell,\nu} i \, g_f \, A \, \bar{f} \gamma^5 f$$

- Direct DM detection suppressed
- SM Yukawa-like couplings: $g_f = \sqrt{2} g_Y m_f / v$
- No missing E_T signals for $m_{\chi} > m_A/2$

Flavor & high- p_T as complementary NP probes

Example: simplified DM models with (pseudo)scalar mediators

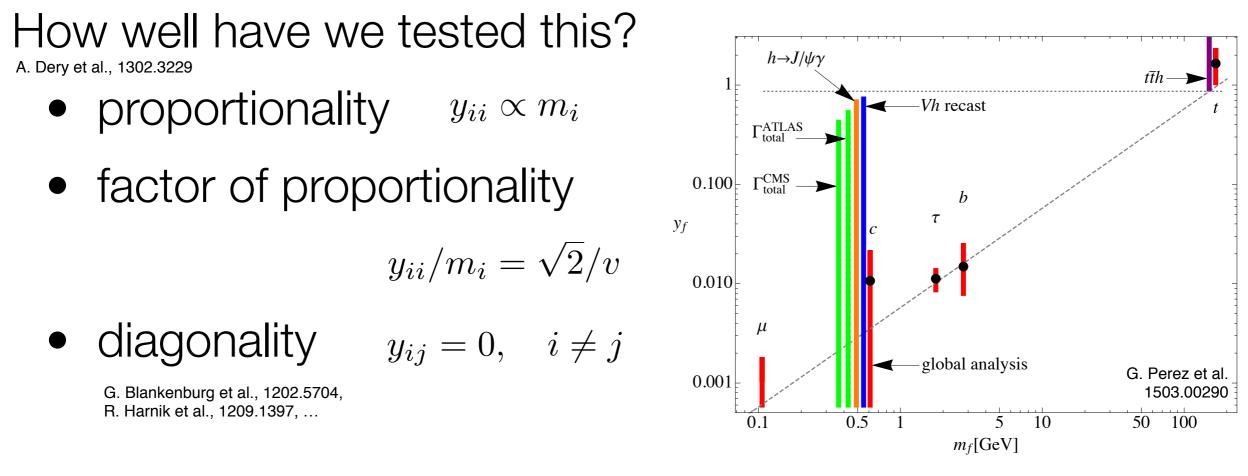


Flavor probes of the Higgs sector

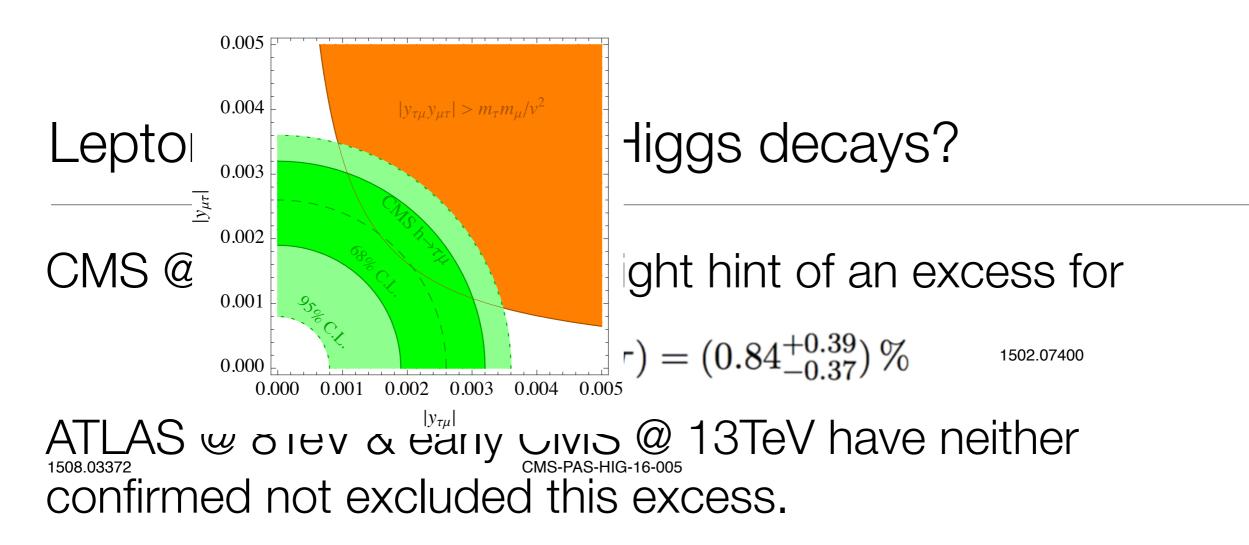
generation of masses in SM through Higgs mechanism

 \Rightarrow Higgs has hierarchical couplings to fermions

$$y_f^{\rm SM} = \sqrt{2}m_f/v$$



Many recent proposals...



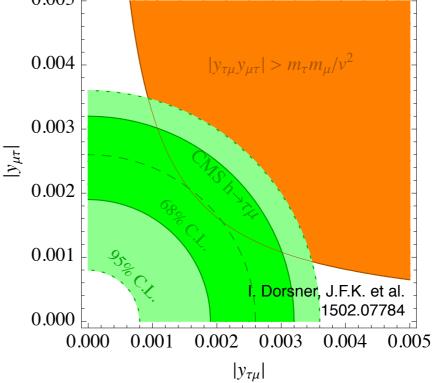
Can one have lepton flavor violating Higgs decays at ~1% level in reasonable NP models?

$$\mathcal{L}_{Y_{\ell}}^{\text{eff.}} = -m_i \delta_{ij} \bar{\ell}_L^i \ell_R^j - y_{ij} \left(\bar{\ell}_L^i \ell_R^j \right) h + \ldots + \text{h.c.},$$

radiative stability of Yukawas

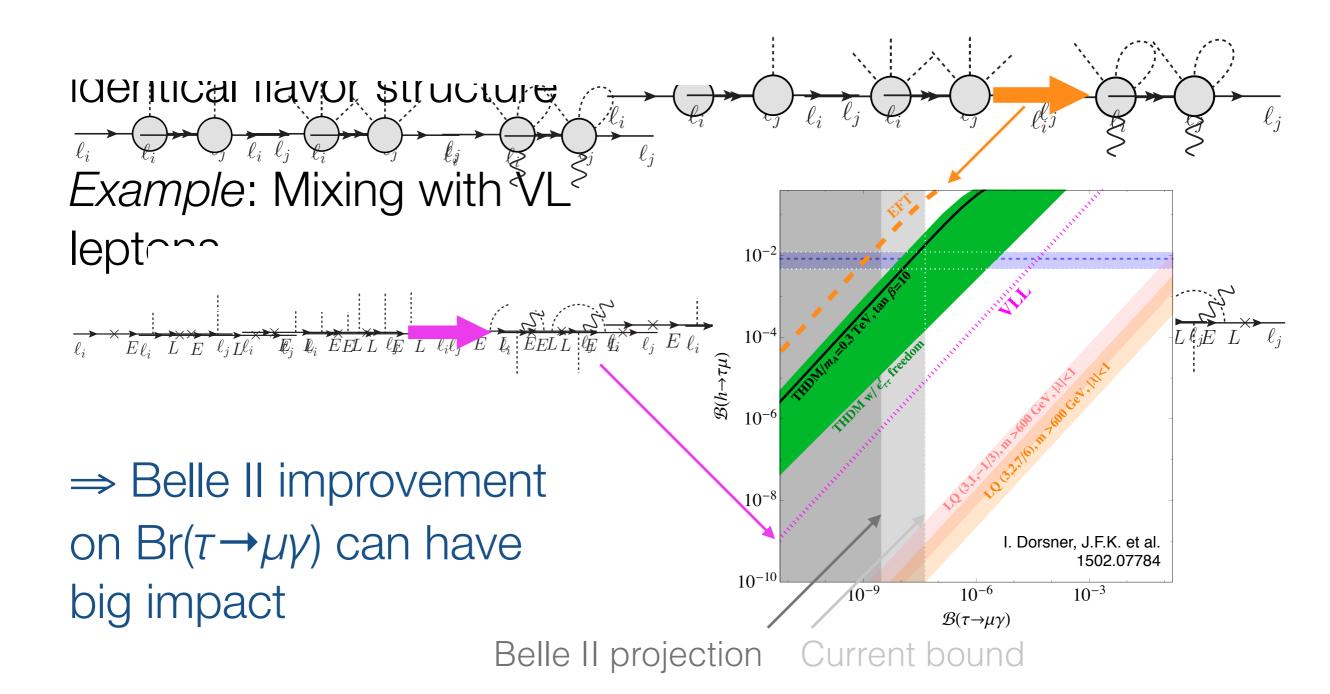
Ceng & Sher, PhysRevD.35.3484, Branco et al., 1106.0034

$$\sqrt{|y_{\tau\mu}y_{\mu\tau}|} \lesssim \frac{\sqrt{m_{\mu}m_{\tau}}}{v} = 0.0018$$



Lepton Flavor violating Higgs decays?

Generic obstacle to large $h \rightarrow \underline{-}$ is the bound on $\underline{-}$ in h and $\underline{-}$ is the bound on $\underline{-}$ in h and $\underline{-}$ is the bound on $\underline{-}$ is the bound on \underline



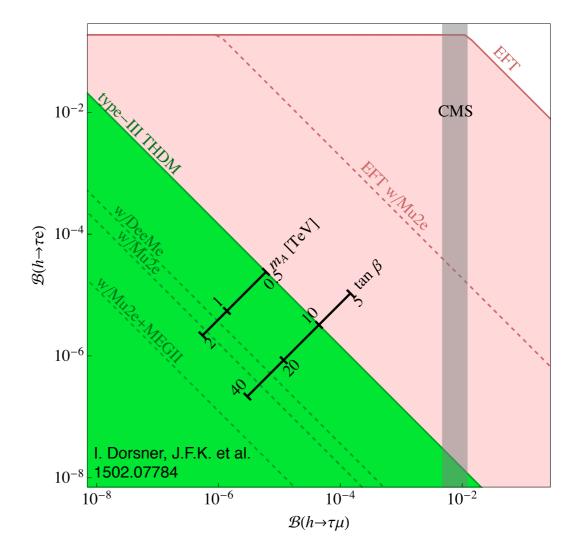
Lepton Flavor violating Higgs decays?

Solution: decouple fermion mass generation from $\tau \rightarrow \mu \gamma$

 possible if new source of EWSB (strong dynamics, multi-Higgs doublet models) ^{W. Altmannshofer et al. 1507.07927} see also A. Crivellin et al. 1611.02703

Generic implications:

- LFV tau decays close to present bounds
- New d.o.f.s within LHC energy reach
- $h \rightarrow \tau e$ constrained by $\mu \rightarrow e$ conversions



Flavor structure affecting high-p_T searches

Nontrivial flavor structure can have important implications also for on-shell searches

Example: THDM accommodating $h \rightarrow \tau \mu$

 $\Delta y_t = [0.1, 1]$ \Rightarrow Both neutral and charged 0.50 heavy scalars can dominantly $Br(H^+ \rightarrow XY)$ decay to leptons 0.10 0.05 lv tb \Rightarrow By SU(2) invariance: A. Efrati, J.F.K. & Y. Nir hW 1606.07082 0.01 0.010 0.001 0.100 $R_{hW^+Z}^{H^+/A} \simeq R_{t\bar{b}\bar{t}}^{H^+/A} = R_{\tau^+\nu\mu^-}^{H^+/A} + R_{\mu^+\nu\tau^-}^{H^+/A} = 1$ $sin \alpha_{vh}$ $R_{XYZ}^{H^+/A} \equiv \frac{\Gamma(H^+ \to XY)}{\Gamma(A \to XZ)} \quad \text{\& assuming } m_{H^+} \sim m_A \qquad v-h \text{ misalignement angle}$

Flavor anomalies motivate high- p_T searches

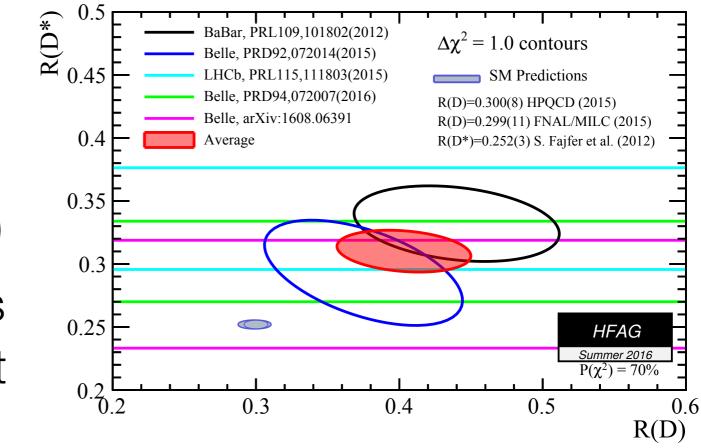
Charged-current decays $B \rightarrow D^{(*)}/v$:

- with I = e, μ used to measure CKM element V_{cb}
- $B \rightarrow D^{(*)}\tau v$ precisely predicted in SM if normalized to $I = e, \mu$ modes $R(D^{(*)})$.

Intriguing exp. situation:

 \Rightarrow 3.9 σ combined tension with SM (HFAG)

 \Rightarrow SM (FF) uncertainties insignificant at this point



Flavor anomalies motivate high- p_T searches

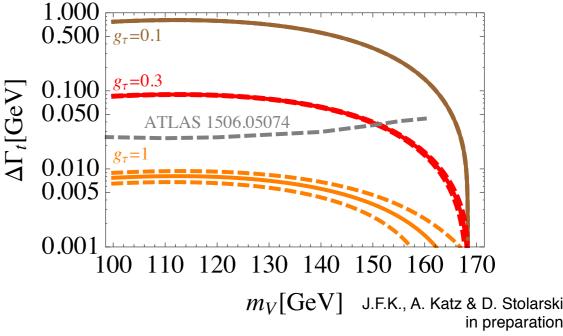
Size of effect calls for tree-level NP, needs to be EM charged.

 \Rightarrow Mass bounds from LEP M > 100 GeV

FCNC & LFU constraints require flavor alignement with the 3rd generation

$$\Rightarrow$$
 effects in top quark decays

$$\mathcal{L}^{(a)} = \mathcal{L}_{\rm SM} + \frac{1}{4} R^{+}_{\mu\nu} R^{-\mu\nu} - m^{2}_{\rho} \rho^{+}_{\mu} \rho^{-\mu} + \left[g_{b} \sum_{q} V_{qb} \bar{q} \not{\!\!\!/}^{+} P_{L} b + g_{\tau} \bar{\tau} \not{\!\!\!/}^{-} P_{L} \nu_{\tau} + \text{h.c.} \right],$$



LHC measurements starting to constrain $m_V < m_t$ region.

Flavor anomalies motivate high- p_T searches

In addition SU(2) invariance predict significant tau production at LHC

Strong constraints from existing $pp \rightarrow \tau^+\tau^-$ searches at CMS $|g_b g_{\tau}| \times v^2 / M_{Z'}^2$ $|g_b g_{\tau}| \times v^2 / M_{Z'}^2$ D. A. Faroughy, A. Greljo & J.F.K. 1609.07138 50 50 ATLAS 13 TeV, 13.2 fb⁻ ATLAS 13 TeV, 3.2 fb⁻ 40 40 0.09 0.4 $Z'/M_{Z'}$ [%] $Z'/M_{Z'}$ [%] 0.07 0.08 30 30 0.07 0.19 0.06 20 20 0.05 0.05 0.3 10 10 0.04 0.03 0.03 $\Gamma_{Z'} < \Gamma_{Z' \to bb} + \Gamma_{Z' \to \tau\tau}$ 0 0 0.3 0.4 0.5 0.6 0.7 0.6 0.8 1.0 1.2 1.4 1.6 1.8 2.0 0.2 M_Z (TeV) M_Z (TeV)

W'/Z' explanation only allowed if light (M < 500 GeV) or broad ($\Gamma/M > 30\%$)

Flavor is powerful guide to high- p_T searches at LHC:

- to ensure no stone is left unturned (and that the most interesting stones are turned first)
- in case of significant signals of NP in flavor observables can identify prospective LHC experimental targets

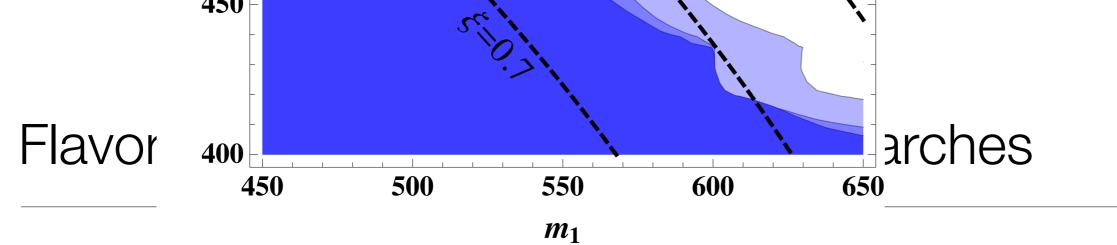
In case new phenomena are discovered at LHC, flavor physics will allow to disentangle different possible interpretations and discriminate between different proposals and scenarios

Examples: 125GeV Higgs & now defunct 750GeV diphoton resonance

In case no new d.o.f.s are seen at LHC, precision tests of flavor, CP, B & L possibly best probes forward

 ⇒ their sensitivity in many cases already (far) exceeds energies/scales attainable in present and planned collider & cosmic ray experiments.

Additional material



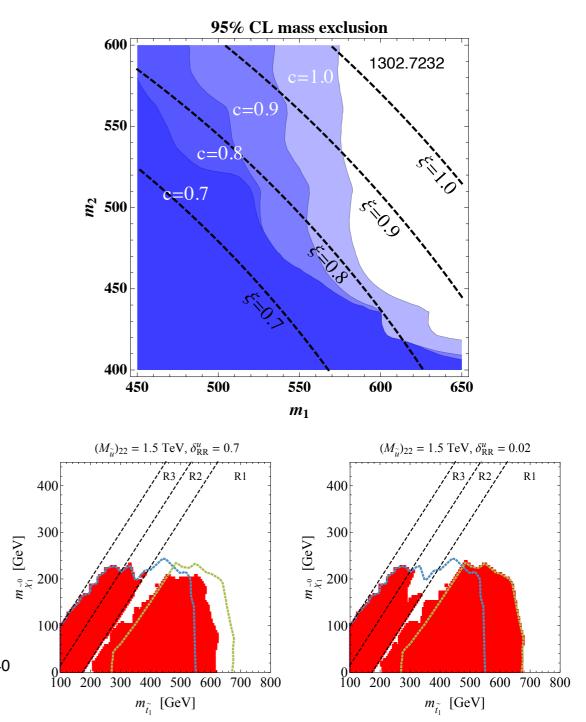
Top partners - direct test of EW naturalness

- stop in SUSY, custodians in composite Higgs
- large flavor breaking can modify exp. searches
- \Rightarrow reduction of fine-tuning
- example: large \tilde{t}_R - \tilde{c}_R mixing in MSSM $\frac{c^2 m_1^2 + s^2 m_2^2}{m_0^2}$

$$c = \cos \theta_R^{ct} \qquad \xi =$$

 \Rightarrow new signature t+c-jet+MET

A. Crivellin et al. 1604.00440



Flavor & high-p_T as complementary NP probes

Further examples: modified tbW, ttZ, tty, ttg, ttH couplings

 $Q_{\phi q,33}^{(3)} \equiv (\phi^{\dagger} i \overset{\leftrightarrow}{D}_{\mu}^{a} \phi) (\bar{Q}_{L,3} \gamma^{\mu} \sigma^{a} Q_{L,3}),$ $Q_{\phi q,33}^{(1)} \equiv (\phi^{\dagger} i \overset{\leftrightarrow}{D}_{\mu} \phi) (\bar{Q}_{L,3} \gamma^{\mu} Q_{L,3}),$ $Q_{\phi u,33} \equiv (\phi^{\dagger} i \overset{\leftrightarrow}{D}_{\mu} \phi) (\bar{t}_{R} \gamma^{\mu} t_{R}),$ $Q_{\phi,33} \equiv (\tilde{\phi}^{\dagger} i \overset{\leftrightarrow}{D}_{\mu} \phi) (\bar{t}_{R} \gamma^{\mu} b_{R}),$

$$Q_{Wd,33} \equiv \bar{Q}_{L,3}(\sigma \cdot W)b_R \phi ,$$

$$Q_{Wu,33} \equiv \bar{Q}_{L,3}(\sigma \cdot W)t_R \tilde{\phi} ,$$

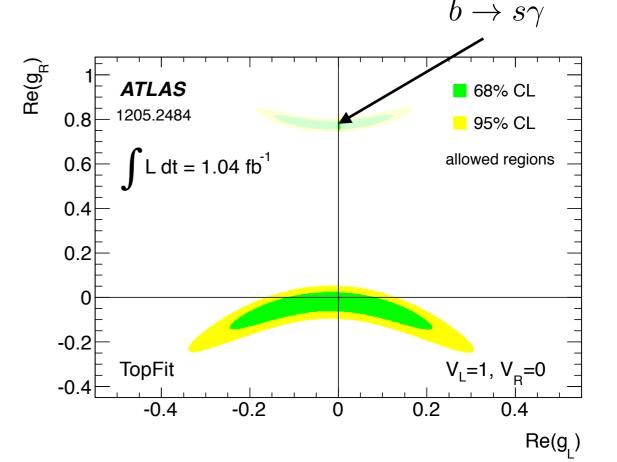
$$Q_{Bu,33} \equiv \bar{Q}_{L,3}(\sigma \cdot B)t_R \tilde{\phi} ,$$

$$Q_{Hu,33} \equiv \bar{Q}_{L,3} t_R \,\tilde{\phi} |\phi|^2 \,,$$

 $Q_{Gu,33} \equiv Q_{L,3}(\sigma \cdot G)t_R \phi$, J. Drobnak, S. Fajfer & J.F.K. 1109.2357, 1102.4347 B. Grzadkowski & M. Misiak 0802.1413

constrained by

- Probed directly through (single, pair, associate) top production and decays at LHC
- Important complementarity with low energy probes



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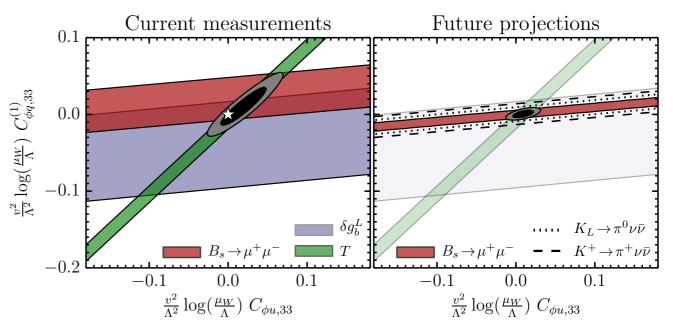
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J. Brod et al. 1408.0792

Flavor & high-p₇ as complementary NP probes

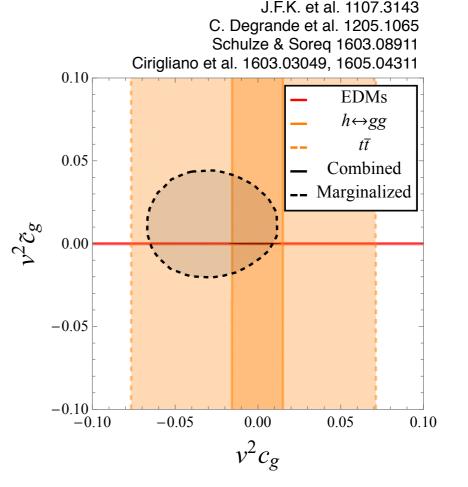
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Flavor & high-p₇ as complementary NP probes

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J. Brod et al. 1310.1385

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