B⁺ leptonic decays: review and prospects

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Hints for New Physics in Heavy Flavor @ KMI, Nagoya

Outline

- ► Motivations and features
 - * To tag, or not to tag
- $ightharpoonup B^+ o au^+
 u$
- $ightharpoonup B^+ o \ell^+ \nu(\gamma)$
- ► Prospects (Belle II)

Features of $R^+ \rightarrow \ell^+ \nu$

SM predictions

$$\Gamma(B^+ o \ell^+
u) = rac{G_F^2 m_B m_\ell^2}{8\pi} \left(1 - rac{m_\ell^2}{m_B^2}
ight)^2 f_B^2 |V_{ub}|^2$$

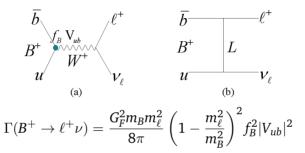
- $\triangleright B(B^+ \to \tau^+ \nu) \sim 10^{-4}$
- ▶ $\mathcal{B}(B^+ \to \mu^+ \nu) \sim \mathcal{B}(B^+ \to \tau^+ \nu)/300$
- $\triangleright \mathcal{B}(B^+ \to e^+ \nu) \sim \mathcal{B}(B^+ \to \tau^+ \nu)/10^7$

Experimental features

 $R^+ \rightarrow \tau^+ \nu$

- large BF, but multiple ν 's
- $ightharpoonup B^+ o \ell^+ \nu \; (\ell \neq \tau) \qquad E_{\ell} \sim M_B/2$, but small BF

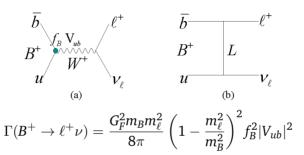
Motivations for $B^+ \rightarrow \ell^+ \nu$



B⁺ leptonic decays: review and prospects

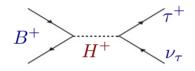
 \triangleright very clean place to measure $f_B|V_{ub}|$ and/or search for new physics (e.g. H^+ , LQ)

Motivations for $B^+ \to \ell^+ \nu$



- ▶ very clean place to measure $f_B|V_{ub}|$ and/or search for new physics (e.g. H^+ , LQ)
- ultimate test of LUV $\Gamma(B^+ \to \ell^+ \nu)/\Gamma(B^+ \to \tau^+ \nu) = f(m_\ell^2, m_\tau^2)$, and all other parameters cancel!

$B^+ \to \tau^+ \nu^-$ by new physics, e.g. H^+



▶ $B^+ \to \tau^+ \nu$ can be affected by new physics effects For instance, H^+ of 2-Higgs doublet model (type II)

$$\mathcal{B}(B^+\to\tau^+\nu)=\mathcal{B}_{\rm SM}(B^+\to\tau^+\nu)\times r_H$$
 where $r_H=\left[1-(m_B^2/m_H^2)\tan^2\beta\right]^2$ W.S. Hou, PRD 48, 2342 (1993)

B⁺ leptonic decays: review and prospects

$B^+ \to \tau^+ \nu$ for new physics

Two useful (for NP) ratios

$$R_{\rm ps} = \frac{\tau_{B^0}}{\tau_{B^+}} \frac{\mathcal{B}(B^+ \to \tau^+ \nu)}{\mathcal{B}(B^0 \to \pi^- \ell^+ \nu)}$$
$$R_{\rm pl} = \frac{\mathcal{B}(B^+ \to \tau^+ \nu)}{\mathcal{B}(B^+ \to \mu^+ \nu)}$$

$$R_{\rm ps}^{\rm NP} = (0.539 \pm 0.043) |1 + r_{\rm NP}^{\tau}|^2,$$

PTEP (2017), 1608.05207

$$R_{\rm pl}^{\rm NP} = \frac{m_{\tau}^2}{m_{\tau}^2} \frac{(1 - m_{\tau}^2/m_B^2)^2}{(1 - m_{\tau}^2/m_D^2)^2} \left| 1 + r_{\rm NP}^{\tau} \right|^2 \simeq 222.37 \left| 1 + r_{\rm NP}^{\tau} \right|^2$$

Tanaka & Watanabe.

To tag, or not to tag

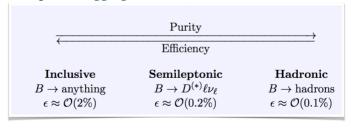
- ▶ Why bother?
 - * $B^+ \to \tau^+ \nu$ has multiple ν 's in the final state
 - * need extra kinematic constraints to improve sensitivity
 - * exploit $\Upsilon(4S)$ producing $B\bar{B}$ and nothing else

$$e^+e^- o \Upsilon(4S) o B_{\mathrm{sig}} \overline{B}_{\mathrm{tag}}$$

► How to tag?

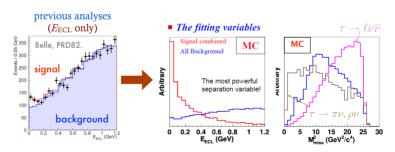
Y. Kwon (Yonsei Univ.)

- * "hadronic tagging" full reconstruction of the decay chain of B_{tag}
- * "semileptonic tagging" use $B^+ \to \overline{D}^{(*)} \ell^+ \nu$



$B^+ \to \tau^+ \nu$ (Belle, had) – signal extraction

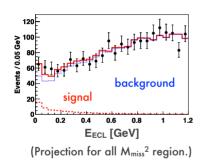
- ► Signal τ modes: $\tau^+ \to e^+ \nu_e \overline{\nu}_\tau$, $\mu^+ \nu_\mu \overline{\nu}_\tau$, $\pi^+ \overline{\nu}_\tau$, $\rho^+ \overline{\nu}_\tau$
- \blacktriangleright π^0 , K_L^0 veto demand no trace of π^0 , K_L^0 after reconstructing B_{tag} and B_{sig}
 - K_L^0 gives $\sim 5\%$ improvement in the expected sensitivity
- ▶ 2D fitting to $E_{\text{ECL}} \& M_{\text{miss}}^2$
 - improve sensitivity by $\sim 20\%$; more robust against peaking backgs. in E_{ECL}

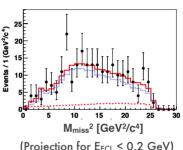


 $E_{\rm ECL}=$ residual energy in the EM calorimeter (ECL) that has not been attributed to either $B_{
m sig}$ or $B_{
m tag}$

$B^+ \to \tau^+ \nu$ (Belle, had) – Result

Simultaneous fit to different τ decay modes Figures below shown for the sum of different τ decay modes





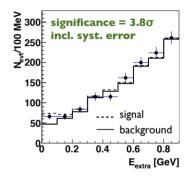
(Projection for E_{ECL} < 0.2 GeV)

- ► Signal yield: $62^{+23}_{-22} \pm 6$ significance = 3.0σ incl. systematic error Major sources of systematic error are: background PDF (8.8%), K_L^0 efficiency (7.3%), and B_{tag} efficiency (7.1%).
- \triangleright $\mathcal{B}(B^+ \to \tau^+ \nu) = (0.72^{+0.27}_{-0.25} \pm 0.11) \times 10^{-4}$

PRL 110, 131801 (2013)

$B^+ \to \tau^+ \nu$ (BABAR, had) – Result

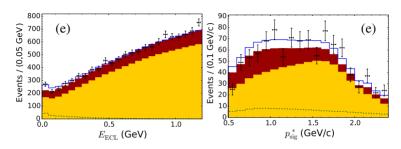
- ► Hadronic *B*-tagging analysis with $N_{B\bar{B}} = 468 \times 10^6$
- Signal τ modes: $\tau^+ \to e^+ \nu_e \overline{\nu}_\tau$, $\mu^+ \nu_\mu \overline{\nu}_\tau$, $\pi^+ \overline{\nu}_\tau$, $\rho^+ \overline{\nu}_\tau$
- Signal extraction via E_{extra} (= E_{ECL}) $N_{\text{sig}} = 62.1 \pm 17.3$ from simultaneous fit to the four τ modes
- $\triangleright \mathcal{B}(B^+ \to \tau^+ \nu) = (1.83^{+0.53}_{-0.49} \pm 0.24) \times 10^{-4}$
- ► Major systematic uncertainties are from background PDF's (10%), *B*-tag efficiency (5%), etc.



PRD 88, 031102(R) (2013)

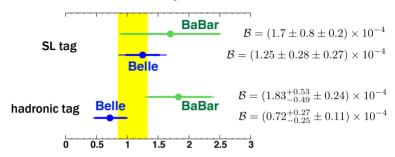
$B^+ \to \tau^+ \nu$ (Belle, SL-tag)





- ▶ tagged by $B^- \to D^{(*)0} \ell^- \overline{\nu}$
- Signal extraction by 2D-fitting (E_{ECL}, p_{sig}^*) $N_{sig} = 222 \pm 50$ events
- $\mathcal{B}(B^+ \to \tau^+ \nu) = (1.25 \pm 0.28 \pm 0.27) \times 10^{-4}$ 4.6 σ significance by combining had-tag and SL-tag analyses of Belle

$B^+ \to \tau^+ \nu$ Summary



Belle combined:
$$\mathcal{B}(B^+ \to \tau^+ \nu) = (0.91 \pm 0.22) \times 10^{-4}$$

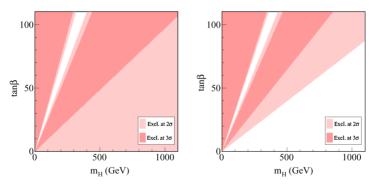
BABAR combined: $\mathcal{B}(B^+ \to \tau^+ \nu) = (1.79 \pm 0.48) \times 10^{-4}$
World avg: $\mathcal{B}(B^+ \to \tau^+ \nu) = (1.06 \pm 0.19) \times 10^{-4}$ HFLAV (2017)

- ▶ Belle vs. BABAR consistent within $\sim 1.7\sigma$
- ► The average is consistent with SM

$B^+ \to \tau^+ \nu$ constraints on charged Higgs

▶ With 2-Higgs doublet model (type II),

$$\mathcal{B}(B^+ \to \tau^+ \nu) = \mathcal{B}_{\rm SM}(B^+ \to \tau^+ \nu) \times \left[1 - (m_B^2/m_H^2) \tan^2 \beta \right]^2$$



Plots are from PRD 88, 031102(R) (2013), by BABAR, based on BABAR's combined $\mathcal{B}(B^+ \to \tau^+ \nu)$.

Search for $B^+ \to \ell^+ \nu$

- ► (experimental) very clean
 - * just a mono-energetic charged lepton and nothing else
- (theoretical) very small branching fraction compared to $B^+ \to \tau^+ \nu$
 - * helicity suppression: $\Gamma \propto m_\ell^2$
- ► Tagged vs. Untagged for $B^+ \to \ell^+ \nu$,
 - * tagging is not really necessary $\cdot \cdot \cdot$ mono-energetic ℓ^+ in the final state
 - * Nonetheless, analyses with tagging have also been tried

$\Gamma(B^+ \to e^+ \nu_e) / \Gamma_{\text{total}}$

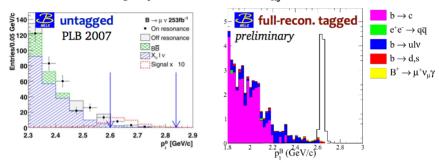
$VALUE (10^{-6})$	CL%		DOCUME	NT ID	TECN	COMMENT	_
< 0.98	90	1	SATOYAN	MA 2007	BELL	$e^+~e^-\to \Upsilon(4S)$	untagged
*** We do not u	se the	foll	owing data	for average	es, fits, li	mits, etc ***	
<3.5	90	2	YOOK	2015	BELL	$e^+~e^-\to \Upsilon(4S)$	had tag
<8	90	1	AUBERT	2010E	BABR	$e^+~e^-\to \Upsilon(4S)$	SL tag
<1.9	90	1	AUBERT	2009V	BABR	$e^+~e^-\to \Upsilon(4S)$	untagged
<5.2	90	1	AUBERT	2008AD	BABR	$e^+~e^-\to \Upsilon(4S)$	had tag

$$\Gamma(B^+ \to \mu^+ \nu_\mu)/\Gamma_{\text{total}}$$

	$VALUE (10^{-6})$	CL%		DOCUMENT ID	TECN COMMENT
untagged	< 1.0	90	1	AUBERT 2009V	BABR $e^+ e^- \rightarrow \Upsilon(4S)$
	*** We do not u	se the	foll	owing data for average	es, fits, limits, etc ***
had tag	<2.7	90	2	YOOK 2015	BELL $e^+ e^- \rightarrow \Upsilon(4S)$
SL tag	<11	90	1	AUBERT 2010E	BABR $e^+ e^- \rightarrow \Upsilon(4S)$
had tag	<5.6	90	1	AUBERT 2008AD	BABR $e^+ e^- \rightarrow \Upsilon(4S)$
untagged	<1.7	90	1	SATOYAMA 2007	BELL $e^+ e^- \rightarrow \Upsilon(4S)$

Why then bother with 'tagged' for $B^+ \to \ell^+ \nu$?

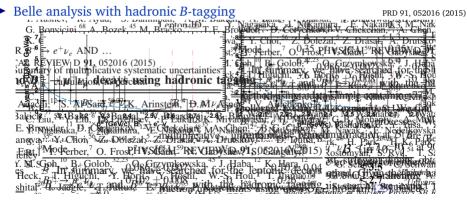
- The signal lepton candidate's momentum in B_{sig} rest frame. -



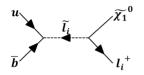
- \blacktriangleright much better resolution of p_{ℓ}^{B} with the full-recon. tagging
- ▶ But, does it make a case for 'full-recon-tagged' analysis of $B^+ \to \ell^+ \nu$?

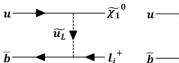
Why then bother with 'tagged' for $B^+ \to \ell^+ \nu$?

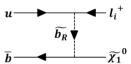
- Note: $\mathcal{B}_{SM}(B^+ \to e^+ \nu) \sim 10^{-11}$ and $\mathcal{B}_{SM}(B^+ \to \mu^+ \nu) \sim 3 \times 10^{-7}$ \Rightarrow Any signal for $B^+ \to e^+ \nu$ at the Belle sensitivity is way beyond the SM
- ► In that case, are we *sure* what we see is *really* $B^+ \to e^+ \nu$? What about $B^0 \to e^+ \tau^-$? How about $B^+ \to e^+ X^0$ where X^0 is any unknown particle from NP?
- ▶ With full-recon., we can use p_{ℓ}^{B} to discern many such cases

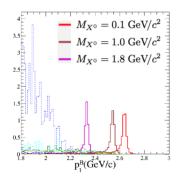


$B^+ \to \ell^+ X^0$ (Belle)



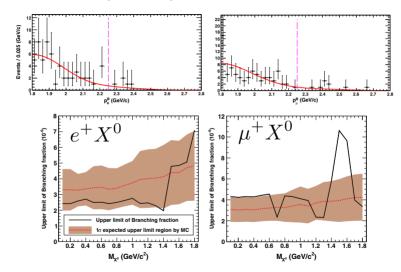






- Search for massive neutral invisible fermion " X^0 "
- a heavy neutrino, or an LSP in RPV models, or whatever
- Very similar experimental signature to $B^+ \rightarrow \ell^+ \nu$
- ▶ But, p_{ℓ}^{B} gives a handle on M_{X}

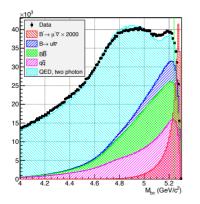
$B^+ o \ell^+ X^0$ (Belle)

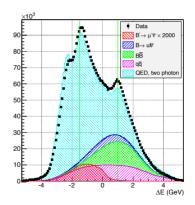


PRD 94, 012003 (2016)

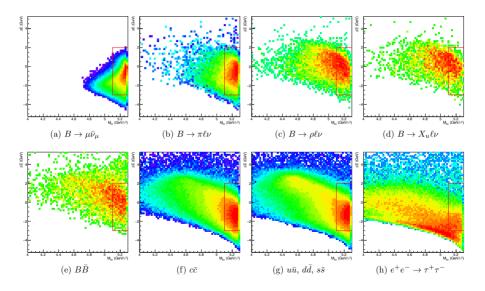
untagged $B^+ o \mu^+ \nu$ Belle (2018)

- ▶ all particles except for the μ^+ are to come from the other B, but its decay chain is not explicitly reconstructed (hence, untagged)
- require $M_{\rm bc} > 5.1$ and $-3.0 < \Delta E < +2.0$





untagged $B^+ o \mu^+ \nu$ Belle (2018)



untagged $B^+ \to \mu^+ \nu^-$ Belle (2018)

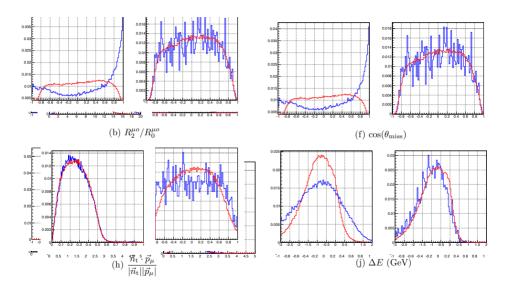
- ▶ all particles except for the μ^+ are to come from the other B, but its decay chain is not explicitly reconstructed (hence, untagged)
- require $M_{\rm bc} > 5.1$ and $-3.0 < \Delta E < +2.0$
- ► In the B^+ rest frame, $p_{\mu} = 2.64$ GeV (*sharp!*), but in the CM frame, $2.45 < p_{\mu}^* < 2.85$ GeV
- Use p_{μ}^* and neural net (NN) for signal extraction (2D fit)

NN variables

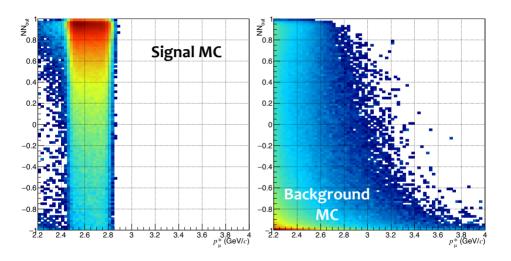
- $R_1^{\mu o}/R_0^{\mu o}$, $R_2^{\mu o}/R_0^{\mu o}$, $R_3^{\mu o}/R_0^{\mu o}$ where $R_i^{\mu o}=\sum_j$ $\vec{n_t} \cdot \vec{p_\mu} = \vec{n_t} \cdot \vec{p_\mu}$ $|\vec{n_t}||\vec{p_u}|$ angle between thrust and muon the muon, and $P_i(x)$ is the i^{th} Legendre polynomial 17b. 17c.
- R_1^{oo}/R_0^{oo} where $R_i^{oo} = \sum_k \sum_j |\vec{p}_k| |\vec{p}_j| P_i(\cos \theta_{kj})$,
- $R_1^{\text{KFW}} = \sum_{k} \sum_{j>k} |\vec{p}_k| |\vec{p}_j| P_i(\cos\theta_{kj})$, the first Kakı cm frame, see Fig. 17e.
- $\cos(\theta_{\text{miss}})$ angle of missing momentum in the
- $\sqrt{\sqrt{\Delta Z^2}}$ distance between reconstructed ztransformation tries to make the strongly peak the neural net catch the small difference bety shown in Fig. 17g. The square root function se discriminating variable away from zero.

- $s = 1 \vec{n}_t^2$ sphericity, see Fig. 17i
- ΔE difference between the sum of energy signal muon and expected energy of B mes
- $\frac{\vec{n}_{t}^{ECL} \cdot \vec{p}_{\mu}}{|\vec{n}^{ECL}||\vec{n}||}$ where the thrust vector \vec{n}_{t}^{ECL} i and calculated in the lab frame, \vec{p}_{μ} is in th
- $(q_{\mu} + q_{\text{tag}}) \times q_{\mu}$ charge balance, see Fig. 1
- $\frac{\vec{p}_{\mu} \cdot \vec{p}_{B_{\text{tag}}}}{|\vec{p}_{\mu}||\vec{p}_{B_{\text{tag}}}|}$ angle between muon and tag i
- $\cos \theta_{\mu}$ muon angle in the cm frame, see F

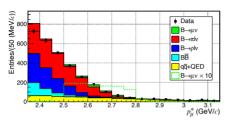
NN variables

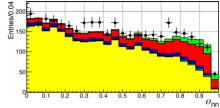


2D distributions (MC) for signal fit



untagged $B^+ \to \mu^+ \nu^-$ Belle (2018) Result





- $B \to \pi \ell \nu$, $\rho \ell \nu$, studied in detail by FF variation
- measure $R \equiv N_{B \to \mu\nu}/N_{B \to \pi \ell \nu}$ for (partial) cancellation of syst. error
- most significant (2.4 σ), and consistent with SM

$$\mathcal{B}(B^+ \to \mu^+ \nu) = (6.46 \pm 2.22 \pm 1.60) \times 10^{-7}$$

 $\in [2.9, 10.7] \times 10^{-7}$ @ 90% C.L.

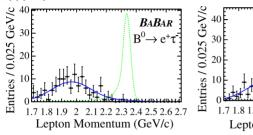
Belle, PRL 121, 031801 (2018)

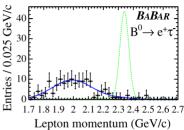
$B^0 o \ell^\pm au^\mp$ (BABAR)



27

PRD 77, 091104(R) (2008)





- ▶ In a hadronic *B*-tagging analysis very similar to $B^+ \to \ell^+ \nu$, *BABAR* also searched for $B^0 \to \ell^{\pm} \tau^{\mp}$.
- Background suppression using m_{ES} and E_{extra}
- Signal extraction by unbinned max. likelihood fit to p_{ℓ}^{B}

$${\cal B}(B^0 o e^\pm au^\mp) < 2.8 imes 10^{-5} \ {\cal B}(B^0 o \mu^\pm au^\mp) < 2.2 imes 10^{-5}$$

$$B^+ \to \ell^+ \nu \gamma$$

► Helicity suppression (of $B^+ \to \ell^+ \nu$) is avoided by γ .

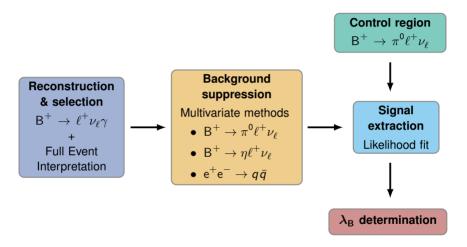
$$\frac{d\Gamma(B^+ \to \ell^+ \nu \gamma)}{dE_{\gamma}} = \frac{\alpha_{\rm em} G_{\rm F}^2 |V_{ub}|^2}{6\pi^2} m_B E_{\gamma}^3 \left(1 - \frac{2E_{\gamma}}{m_B}\right) \left(\left|F_{\rm V}\right|^2 + \left|F_{\rm A} + \frac{e_{\ell} f_B}{E_{\gamma}}\right|^2\right)$$

$$F_{\rm V}(E_{\gamma}),\; F_{\rm V}(E_{\gamma}) \sim rac{e_{
m u}f_Bm_B}{2E_{\gamma}\lambda_B} + \cdots$$

- \triangleright λ_B is needed for QCDF to calculate, e.g., charmless hadronic B decays
- ► SM expectation: $\mathcal{B}(B^+ \to \ell^+ \nu \gamma) \sim \mathcal{O}(10^{-6})$
 - * Calculation is reliable only for $E_{\gamma} > 1$ GeV
- ▶ Previous Belle (2015): $\Delta \mathcal{B}(B^+ \to \ell^+ \nu \gamma) < 3.5 \times 10^{-6}$
- ▶ Updated results from Belle (2018) with 'FEI' algorithm
 - a new B-tagging algorithm developed for Belle II

B⁺ leptonic decays: review and prospects

$B^+ \to \ell^+ \nu \gamma$ Belle (2018) analysis strategy

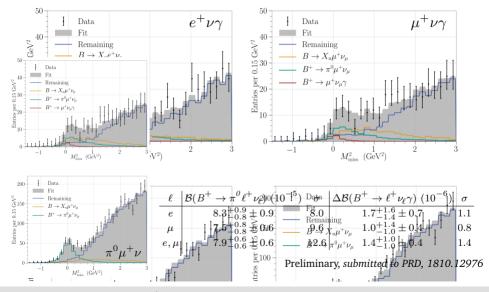


from M. Gelb talk at CKM2018

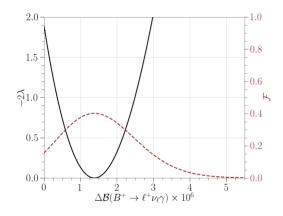
$$B^+ \to \ell^+ \nu \gamma$$
 Belle (2018) features

- ► Measure $B^+ \to \pi^0 \ell^+ \nu$ separately ("control sample"), to constrain the peaking background
- ► Two parameters
 - * $\Delta \mathcal{B}(B^+ \to \ell^+ \nu \gamma)_{E_{\gamma} > 1.0 \text{GeV}}$
 - * $R_{\pi} = \Delta \mathcal{B}(B^+ \to \ell^+ \nu \gamma)_{E_{\gamma} > 1.0 \text{GeV}} / \mathcal{B}(B^+ \to \pi^0 \ell^+ \nu)$ \$\Rightarrow\$ This allows to extract \$\lambda_B\$ independent of $|V_{ub}|$, and some systematics cancel in the ratio R_{π} .

$B^+ \to \ell^+ \nu \gamma$ Belle (2018) results



$B^+ \to \ell^+ \nu \gamma$ Belle (2018) upper limits



Bayesian limit

$$0.9 = \frac{\int_0^{\text{UL}} \mathcal{F}(\Delta \mathcal{B}) d\Delta \mathcal{B}}{\int_0^{\infty} \mathcal{F}(\Delta \mathcal{B}) d\Delta \mathcal{B}}$$

ℓ	BaBar	Belle (2015)	Belle (2018)
\overline{e}	-	< 6.1	< 4.3
μ	-	< 3.4	< 3.4
e, μ	< 14	< 3.5	< 3.0

Preliminary, submitted to PRD, 1810.12976

$$B^+ \to \ell^+ \nu \gamma$$
 Belle (2018) for λ_B

$$R_{\pi}^{\text{meas}} = (1.7 \pm 1.4) \times 10^{-2}$$

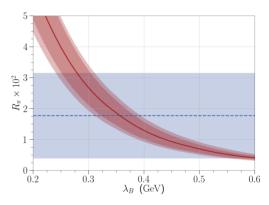
$$R_{\pi} = \frac{\Delta\Gamma(\lambda_B)}{\Gamma(B^+ \to \pi^0 \ell^+ \nu)}$$

Use theory to determine interval for λ_B

- Beneke, Braun, Ji, Wei, JHEP 1807, 154 (2018)
- HFLAV, EPJC 77, 895 (2017)

Two one-sided limits

 $\lambda_B > 0.24 \text{ GeV}$ and $\lambda_B < 0.68 \text{ GeV}$



Preliminary, submitted to PRD, 1810.12976

$B^+ o au^+ u$ Prospects for Belle II

- ► E_{ECL} is crucial for $B^+ \to \tau^+ \nu$ study
 - * In Belle II, beam background is much higher
 - * But such backgrounds can be rejected by tighter selection based on ECL cluster's energy, timing, shape, etc.
- Expected precision at $1 \text{ ab}^{-1} \sim 29\%$ (stat.)
- ▶ Major systematic sources (bkg. PDF, K_L^0 veto eff., B_{tag} eff., etc.) can be improved with more data

	Integrated Luminosity (ab^{-1})	1	5	50
	statistical uncertainty (%)	29	13	4
hadronic tag	systematic uncertainty (%)	13	7	5
	total uncertainty (%)	32	15	6
semileptonic tag	statistical uncertainty (%)	19	8	3
	systematic uncertainty (%)	18	9	5
	total uncertainty (%)	26	12	5

$B^+ \to \tau^+ \nu$ Prospects for Belle II

Two useful (for NP) ratios

$$R_{\rm ps} = \frac{\tau_{B^0}}{\tau_{B^+}} \frac{\mathcal{B}(B^+ \to \tau^+ \nu)}{\mathcal{B}(B^0 \to \pi^- \ell^+ \nu)}$$
$$R_{\rm pl} = \frac{\mathcal{B}(B^+ \to \tau^+ \nu)}{\mathcal{B}(B^+ \to \mu^+ \nu)}$$

$$\begin{split} R_{\mathrm{ps}}^{\mathrm{NP}} &= \left(0.539 \pm 0.043\right) \! \left| 1 + r_{\mathrm{NP}}^{\tau} \right|^2, \\ R_{\mathrm{pl}}^{\mathrm{NP}} &= \frac{m_{\tau}^2}{m_{\mu}^2} \frac{(1 - m_{\tau}^2 / m_B^2)^2}{(1 - m_{\mu}^2 / m_B^2)^2} \! \left| 1 + r_{\mathrm{NP}}^{\tau} \right|^2 \simeq 222.37 \left| 1 + r_{\mathrm{NP}}^{\tau} \right|^2 \end{split}$$

Luminosity	$R_{ m ps}$	$R_{ m pl}$
$5\mathrm{ab^{-1}}$	[-0.22, 0.20]	[-0.42, 0.29]
$50\mathrm{ab^{-1}}$	[-0.11, 0.12]	[-0.12, 0.11]

Expected sensitivity @ 95% CL.

Assumed: NP contribution is real and $|r_{\mathrm{NP}}^{ au}| < 1$

NP contributions to $B^+ \to \tau^+ \nu$ with $|r_{\rm NP}| > \mathcal{O}(0.1)$ can be tested at 95% CL.

$B^+ \to \mu^+ \nu$ Prospects for Belle II

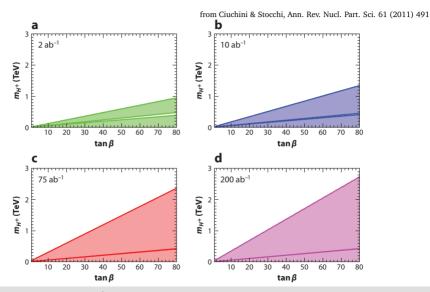
▶ By scaling the FoM of Belle new untagged analysis (PRL 2018),

$$\mathcal{F}_{B2} = \mathcal{F}_{B1} \times \sqrt{50~ab^{-1}/0.711~ab^{-1}} \sim 14.5\%$$

corresponding to $\sim 7\%$ statistical precision

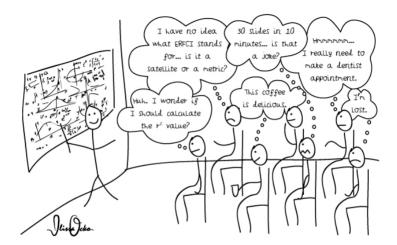
- ▶ naive expectation (Ref. B2TiP draft)
 - * $B^+ \rightarrow \mu^+ \nu$ can reach 5σ with $\sim 6 \text{ ab}^{-1}$
 - * 5% statistical precision, with full 50 ab⁻¹

$B^+ \to \ell^+ \nu$ Prospects beyond 50 ab⁻¹



Concluding Remarks

- ▶ Leptonic *B* decays, in particular $B^+ \to \ell^+ \nu$ ($\ell = e, \mu, \tau$), provide powerful probe for new physics beyond the SM.
- ▶ $B^+ \to \tau^+ \nu$ decays have been measured at nearly 5σ significance, and new physics models such as 2HDM (II) have been tested.
- ▶ With hadronic *B*-tagging, Belle has searched for *invisible*, *massive*, *lepton-like neutral* particle X^0 in $B^+ \to \ell^+ X^0$ for the first time.
- ▶ Belle II with $\int \mathcal{L} dt = 50 \text{ ab}^{-1}$ branching fractions for both $B^+ \to \tau^+ \nu$ and $B^+ \to \mu^+ \nu$ are expected to be measured with precision of $\sim 5\%$.

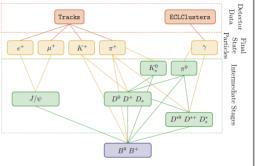


Thank you!

A new B-tagging: Full Event Interpretation

- Hierarchical reconstruction of B_{tag} with a network of classifiers
- Successor of the Belle Full Reconstruction (FR)
- Training and application
- Hadronic and semi-leptonic tag modes
- Generic FFI:
 - 1) FEI trained and applied on full event
 - 2) Signal selection
- Signal-specific FEI (new):
 - 1) Signal selection
 - FEI trained and applied on rest-of-event → trained on specific event topology
- Each B_{tag} candidate has an assigned probability P_{FEI}

from M. Gelb talk at CKM2018



Tagging	officionov	on MC

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Tag	FR ¹	gen. FEI Belle	gen. FEI Belle II			
Hadronic B ⁺	0.28%	0.76%	0.66%			
SL B ⁺	0.67%	1.80%	1.45%			
Hadronic B ⁰	0.18%	0.46%	0.38%			
SL B ⁰	0.63%	2.04%	1.94%			

¹Belle Full Reconstruction algorithm.